

# Supporting Information:

## Investigations of Metal Phosphides as Potential Thermoelectric Materials

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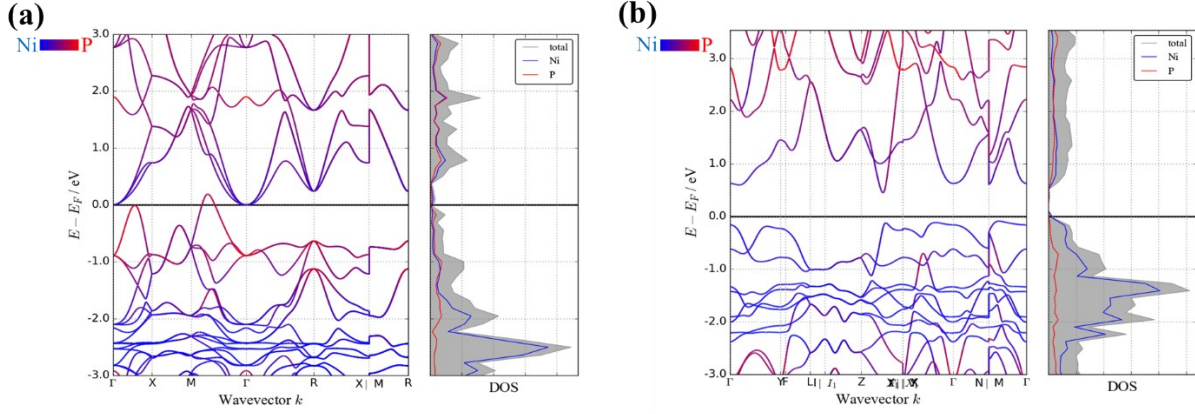
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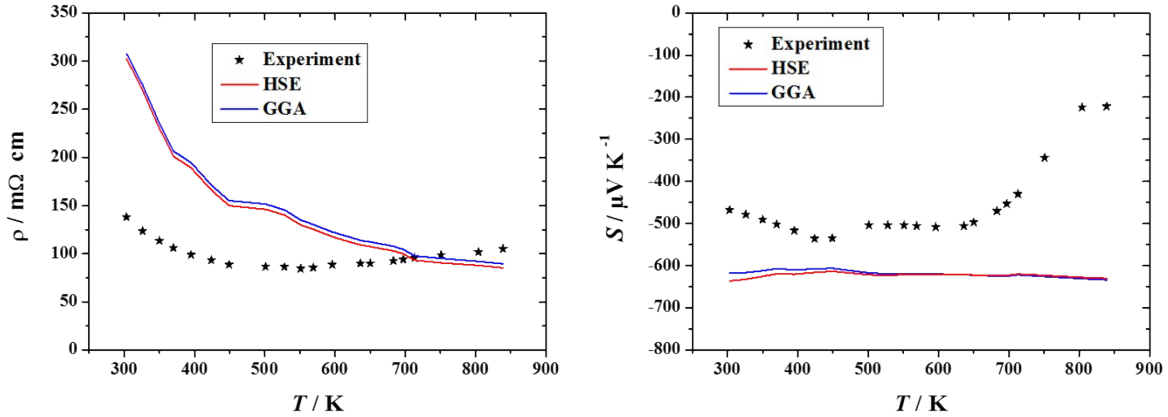
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**Figure S1:** Electronic band structure of (a) cubic NiP<sub>2</sub> (mp-22619) and (b) monoclinic NiP<sub>2</sub> (mp-486) computed with DFT-PBE-GGA.



**Figure S2:** Comparison between computed electrical resistivity and Seebeck coefficient of BP to experimental values reported by Kumashiro *et al.*<sup>[1]</sup> While the computed resistivity is similar to experiments, the magnitude of the computed Seebeck coefficient is overestimated.

## 1. Minimum thermal conductivity models

The most established model to determine the amorphous limit is the Cahill-Pohl model<sup>[2,3]</sup> with

$$\kappa_{min,CP} = \frac{1}{2.48} k_B \left( \frac{N}{V} \right)^{2/3} (v_L + 2v_T) \quad \text{SI-1}$$

where  $k_B$  is the Boltzmann constant,  $N$  is the total number of atoms in the unit cell volume  $V$ , and  $v_L$  and  $v_T$  are the longitudinal and transverse speeds of sound, respectively. The speeds of sound were calculated from the computed elastic constants using first-principle calculations and a Voigt-Reuss-Hill average; the stress-strain method used for these calculations is described elsewhere.<sup>[4]</sup> The longitudinal and transverse speeds of sound, therefore, are given by

$$v_L = \left( \frac{K + 4/3G}{d} \right)^{1/2} \quad \text{SI-2}$$

and

$$v_T = \left( \frac{G}{d} \right)^{1/2}, \quad \text{SI-3}$$

where  $K$  is the bulk modulus and  $G$  is the shear modulus.

In a recent report, we demonstrated that the experimental phononic contribution of the thermal conductivity has a strong correlation with the amorphous thermal conductivity.<sup>[5]</sup> However, several materials showed thermal conductivity lower than the predicted amorphous limit; the thermal conductivity in these materials is defined as ‘ultralow’. In a recent study, the origins of the ultralow thermal conductivity of [6,6]-phenyl-C<sub>61</sub>-butyric acid methyl ester (PCBM) was delineated.<sup>[6]</sup> As shown in Equation SI-1 the amorphous limit is calculated from the speed of sound which is taken from the linear slope of the acoustic phonons near the center of the

Brillouin zone. However, a small phonon mean free path ( $\lambda$ ) has a large wavevector ( $k = \frac{2\pi}{\lambda}$ ) and using the linear slope at the center of the Brillouin zone overestimates the phonon mean speed,  $\bar{v}$ . Therefore, the mean phonon speed over the entire Brillouin zone was calculated for the longitudinal and transverse acoustic phonons. The three phonon mean speeds were averaged by

$$\langle \bar{v} \rangle = \left( \frac{1}{3} \left[ \frac{1}{\bar{v}_L^3} + \frac{1}{\bar{v}_{T1}^3} + \frac{1}{\bar{v}_{T2}^3} \right] \right)^{-1/3} \quad \text{SI-4}$$

and an average Debye temperature,  $\langle \theta_D \rangle$ , was computed from

$$\langle \theta_D \rangle = \frac{\langle \bar{v} \rangle \hbar (6\pi^2 N)}{k_B V} \quad \text{SI-5}$$

where  $\hbar$  is the reduced Planck constant. The thermal conductivity in a material can be limited by

a constant phonon mean free path  $\left( \lambda = \left[ \frac{\pi V}{6N} \right]^{1/3} \right)$ <sup>[6]</sup> and a frequency-dependent phonon mean free

path  $\left( \lambda = \frac{\pi}{\omega} \langle \bar{v} \rangle \right)$  as assumed by Einstein.<sup>[7]</sup> While the latter does not allow for the grain boundary limit for low frequencies, it can represent the limit where dynamical disorder is responsible for the low thermal conductivity (referred as ‘dynamic’). The constant mean free path, on the other hand, can delineate the limit for static disordering (referred as ‘static’) which was most likely the

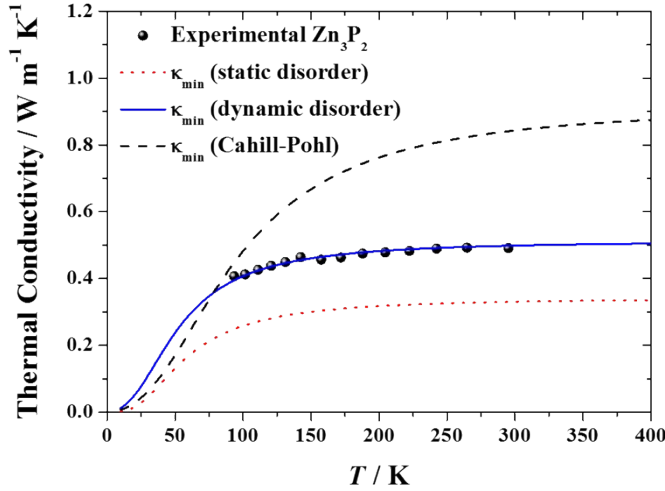
case for the ultralow thermal conductivity in PCBM and  $C_{60}/C_{70}$ .<sup>[6]</sup> The thermal conductivities for the static and dynamic disordering are calculated by  $\langle\theta_D\rangle$  (Equation SI-5) and the atomic density ( $N/V$ )

$$\kappa_{min,st}(T) = \frac{3}{6^{2/3}\pi^{1/3}} \frac{k_B^2}{\hbar} \left(\frac{N}{V}\right)^{1/3} \frac{\langle\theta_D\rangle}{x_D^3} \int_0^{x_D} \frac{x^4 e^x}{(e^x - 1)^2} dx \quad \text{SI-6}$$

and

$$\kappa_{min,dy}(T) = \frac{3}{6^{2/3}\pi^{1/3}} \frac{k_B^2}{\hbar} \left(\frac{N}{V}\right)^{1/3} \frac{\langle\theta_D\rangle}{x_D^2} \int_0^{x_D} \frac{x^3 e^x}{(e^x - 1)^2} dx \quad \text{SI-7}$$

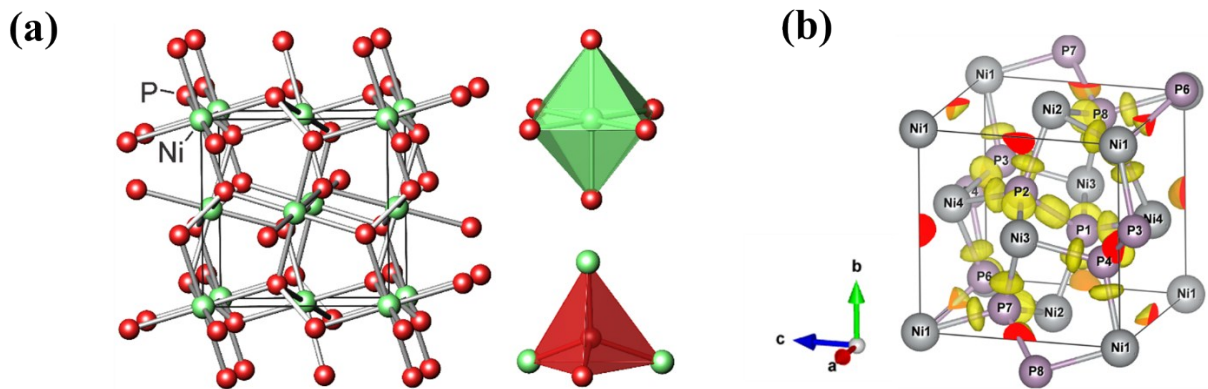
respectively, where  $x_D = \frac{\langle\theta_D\rangle}{T}$ . More information of the derivation of these models can be found in reference [6].



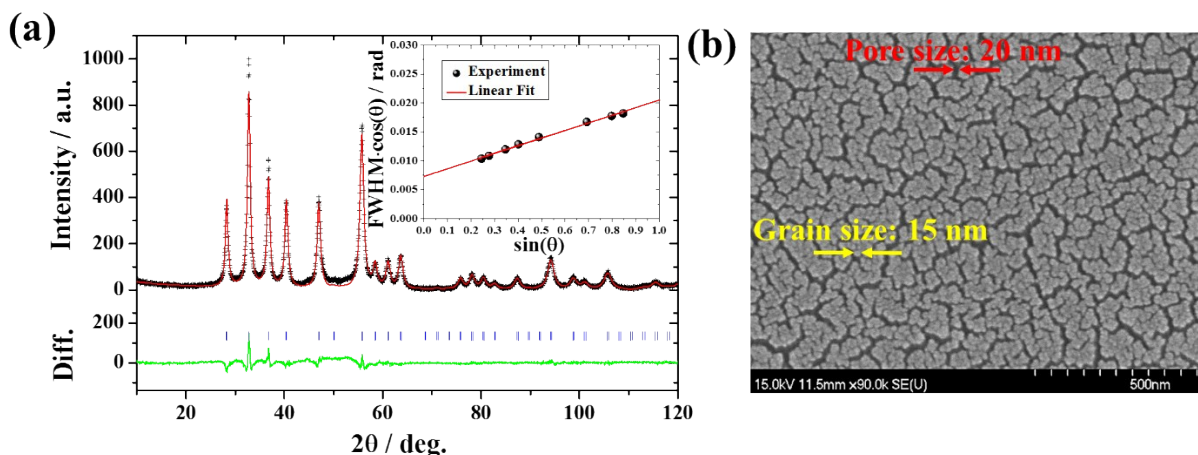
**Figure S3:** Ultralow thermal conductivity in  $Zn_3P_2$  can be explained using an average phonon mean speed from phonon band calculations. While static disordering (frequency-independent phonon mean free path) predicts a lower thermal conductivity, dynamic disordering (frequency-dependent mean free path) agrees well with the experiments.

## 2. Characterization of $c$ -NiP<sub>2</sub>

While the monoclinic phase of  $\text{NiP}_2$  indicates better predicted TE performance than the cubic phase, the synthesis of the former using a solid-state approach as reported by Gillot *et al.*<sup>[8]</sup> was not successful and toxic phosphine gas was produced. Using high-energy ball milling, pure *c*- $\text{NiP}_2$  was synthesized. The crystal structure of *c*- $\text{NiP}_2$  is pyritic where Ni and P occupy octahedral and tetrahedral environments, respectively (Figure S4). The crystal structure of *c*- $\text{NiP}_2$  was confirmed by powder X-ray diffraction (PXRD) analysis (Figure S5 (a)). The refined lattice parameter of the cubic structure ( $a = 5.469(2) \text{ \AA}$ ) is close to the reported value of  $5.4706(2) \text{ \AA}$ <sup>[9]</sup> and to that in our computed structure ( $5.453 \text{ \AA}$ ). In the crystal structure, each P is bonded to one P and three Ni, while each Ni is surrounded by six P atoms (Figure S4 (a)). The tetrahedral (for P) and octahedral (for Ni) environments are slightly distorted. Electron localization function (ELF) calculations show shared electrons for Ni-P indicating polar covalent bonding (Figure S4 (b)). A Bader charge analysis demonstrated that the charge transfer between Ni (+0.11) and P (-0.04; -0.07) is subtle, due to the similar electronegativity of Ni (1.91) and P (2.19). The PXRD spectrum showed broad peaks indicating small grains (Figure S5 (a)). From the Rietveld refinement it was found that the reflections can be fitted only with Lorentzian functions and therefore, the grain size can be determined using a Hall-Williamson plot.<sup>[10]</sup> With the Hall-Williamson plot the peak broadening due to the grain size as well as inhomogeneous strain can be determined (see inset of Figure S5 (a)). From the intercept of the linear fit a grain size of  $19 \pm 1 \text{ nm}$  was determined. This is consistent with the SEM image, which indicates grains in the range of 10 to 20 nm (see Figure S5 (b)). In addition to the grain size, SEM reveals that the *c*- $\text{NiP}_2$  sample contains pores in the range of 10 to 30 nm which results in a reduced density ( $4.04 \text{ g cm}^{-3}$ ) compared to the computed theoretical density ( $4.94 \text{ g cm}^{-3}$ ). The charge transfers of the other MPs are summarized in Table S9.

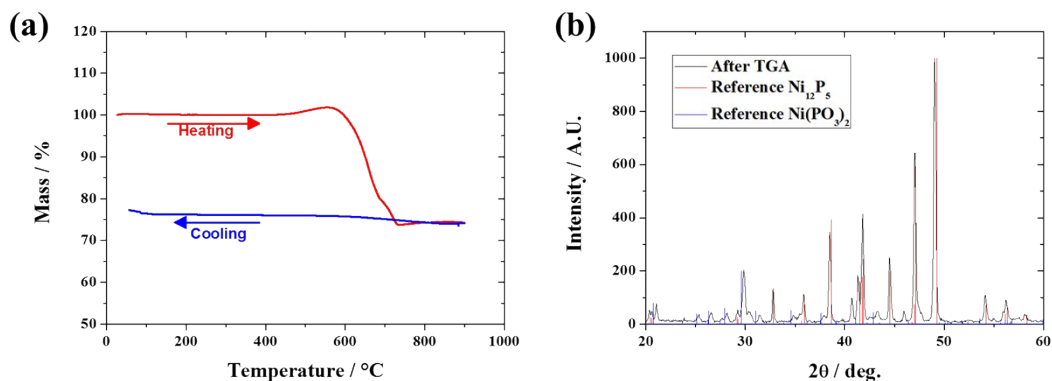


**Figure S4:** (a) Crystal structure and (b) electron localization function (ELF) calculation of cubic  $\text{NiP}_2$ . Each P atom is tetrahedrally bonded to three Ni atoms (P  $-0.07$ ;  $-0.04$ ; Ni  $+0.11$ ) and one P atom; each Ni atom fills the octahedral interstices. The electron densities are found closer to the P atom for Ni-P and at the bond middle point for P-P indicating polar covalent and pure covalent interactions, respectively.

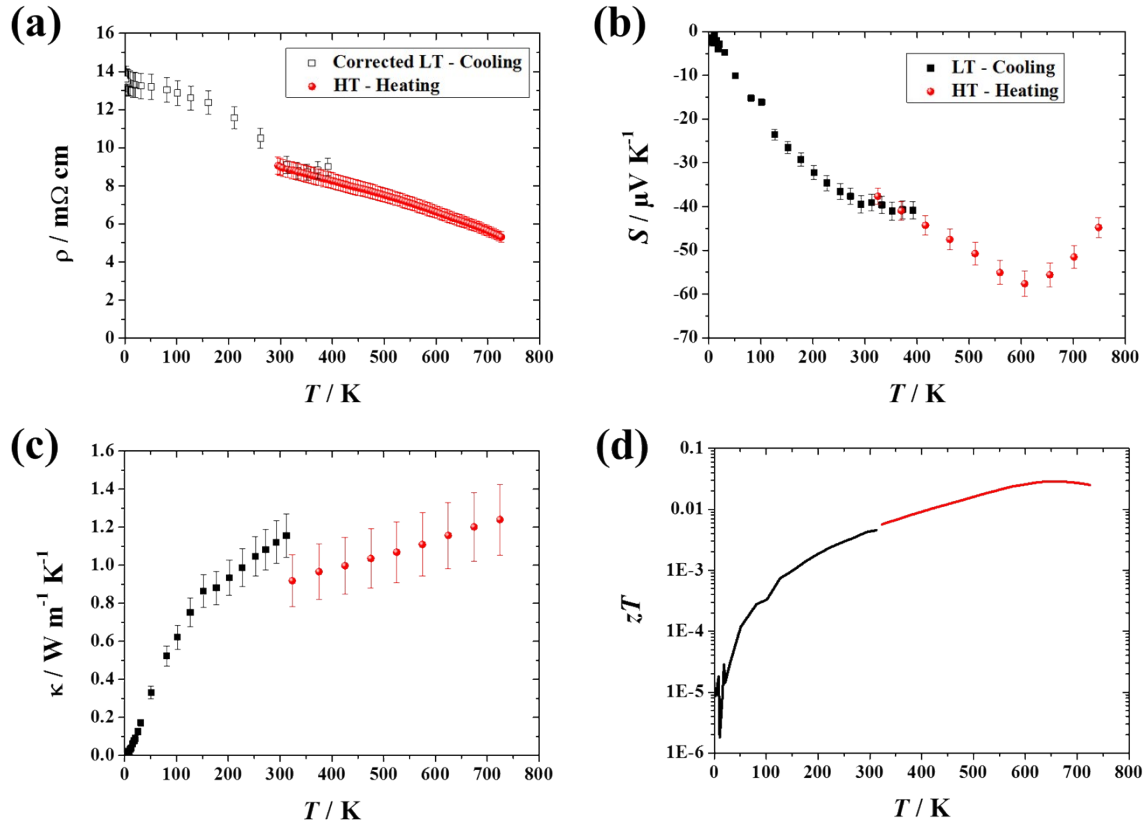


**Figure S5:** (a) Rietveld refinement pattern of cubic  $\text{NiP}_2$  (space group:  $Pa\bar{3}$ ). Blue tick marks show the calculated reflections and the green baseline corresponds to the residuals of the Rietveld refinement. Inset: Hall-Williamson plot gives a grain size of  $19 \pm 1$  nm. (b) SEM image of the surface of a consolidated  $\text{NiP}_2$  pellet, showing grains in the range of 10 to 20 nm.

The temperature stability of  $c\text{-NiP}_2$  is shown in Figure S6 (a). A mass increase was observed at approximately 470 °C, following a decrease in mass from 570 °C to 730 °C. Although the measurement was conducted under inert atmosphere, the increase in mass can be assigned to oxidation of  $\text{NiP}_2$  to  $\text{Ni}(\text{PO}_3)_2$  as observed by the minor phase in PXRD after the thermogravimetric measurement (Figure S6 (b)). At higher temperatures ( $T > 570$  °C), P is removed from the non-oxidized  $\text{NiP}_2$  until it is completely converted to  $\text{Ni}_{12}\text{P}_5$ .



**Figure S6:** (a) Thermogravimetric analysis of  $\text{NiP}_2$  indicates major decomposition sets in at 570 °C. (b) PXRD pattern of sample after TGA shows the major phase is  $\text{Ni}_{12}\text{P}_5$  with a minor phase of  $\text{Ni}(\text{PO}_3)_2$ .



**Figure S7.** The experimental (a) electrical resistivity,  $\rho$ , (b) Seebeck coefficient,  $S$ , (c) thermal conductivity,  $\kappa$ , and (d) figure of merit of  $c$ -NiP<sub>2</sub> as functions of temperature,  $zT$ .

**Table S1:** Computed electronic transport properties for different *p*-type metal phosphorous compounds at  $n = 10^{20} \text{ cm}^{-3}$  and the temperature where the predicted  $zT_{calc}$  is maximum. The electronic transport properties were calculated from the PBE band structures with the Boltztrap code where the constant relaxation time was set to  $10^{-14} \text{ s}$ .

Material	mp-id	Space group	$E_g / \text{eV}$ HSE (PBE)	$T / \text{K}$	$S / \mu\text{V K}^{-1}$	$\rho / \text{m}\Omega \text{ cm}$	$\kappa_{el} / \text{W m}^{-1} \text{ K}^{-1}$	$PF / \text{mW m}^{-1} \text{ K}^{-2}$
BP	1479	216	1.98 (1.24)	1300	318	1.39	1.67	7.28
GaP	2490	216	(1.59)	1300	339	1.67	1.43	6.88
InP	20351	216	(0.47)	1100	324	1.89	1.26	5.54
FeP	1005	62	0.00 (0.00)	1300	39	0.09	38.73	1.69
ZnNaP	4824	129	(0.91)	1300	356	2.26	1.43	5.61
SrLiP	10614	187	(0.80)	1000	378	3.13	0.87	4.56
BaLiP	10615	187	(0.70)	1100	341	2.96	0.03	3.93
SnNaP	29529	186	(0.58)	800	291	1.92	0.10	4.41
NiP <sub>2</sub>	22619	205	0.13 (0.00)	700	80	0.17	16.35	3.76
NiP <sub>2</sub>	486	15	1.15 (0.53)	500	357	2.55	0.00	5.00
ZnP <sub>2</sub>	11025	96	2.20 (1.46)	1300	366	7.93	0.17	1.69
ZnP <sub>2</sub>	2782	92	2.20 (1.46)	1300	365	7.93	0.17	1.68
AgP <sub>2</sub>	8200	14	1.30 (0.66)	1100	278	2.27	0.99	3.40
TiP <sub>2</sub>	20327	62	0.00 (0.00)	1100	31	0.1	37.01	0.96
CuP <sub>2</sub>	927	14	(0.86)	1100	350	2.81	0.74	4.36
RuP <sub>2</sub>	1413	58	(0.48)	800	307	2.08	0.98	4.53
CoP <sub>2</sub>	14285	14	(0.44)	700	290	2.72	0.53	3.09
IrP <sub>2</sub>	10155	14	(0.78)	1100	329	2.07	1.04	5.23
BeP <sub>2</sub>	27148	15	(0.83)	1300	315	2.13	1.55	4.66
RhP <sub>2</sub>	15953	14	(0.37)	800	265	1.62	1.30	4.33
PdP <sub>2</sub>	28266	15	(0.46)	600	298	2.3	0.58	3.86
FeP <sub>2</sub>	20027	58	(0.43)	500	442	12.8	0.09	1.53
Zn <sub>3</sub> P <sub>2</sub>	2071	137	1.17 (0.31)	900	278	1.76	1.04	4.39
Mg <sub>3</sub> P <sub>2</sub>	8085	224	(1.11)	1200	451	6.56	0.24	3.10
Mg <sub>3</sub> P <sub>2</sub>	2514	206	2.20 (1.61)	1300	408	8.4	0.16	1.98
Cd <sub>3</sub> P <sub>2</sub>	2441	137	0.65 (0.002)	600	224	1.97	0.86	2.55
MnP <sub>4</sub>	569522	2	(0.48)	700	315	9.83	0.26	1.01
MnP <sub>4</sub>	487	2	1.58 (0.44)	700	343	5.78	0.38	2.04
PbTe	19717	225	(0.81)	1100	419	1.08	1.97	16.26
Bi <sub>2</sub> Te <sub>3</sub>	568390	166	(0.25)	1000	315	1.46	2.02	6.80

**Table S2:** Computed electronic transport properties for different  $n$ -type metal phosphorous compounds at  $n = 10^{20} \text{ cm}^{-3}$  and the temperature where the predicted  $zT_{calc}$  is maximum. The electronic transport properties were calculated from the PBE band structures with the Boltztrap code where the constant relaxation time was set to  $10^{-14} \text{ s}$ .

Material	mp-id	Space group	$E_g / \text{eV}$ HSE (PBE)	$T / \text{K}$	$S / \mu\text{V K}^{-1}$	$\rho / \text{m}\Omega \text{ cm}$	$\kappa_{el} / \text{W m}^{-1} \text{ K}^{-1}$	$PF / \text{mW m}^{-1} \text{ K}^{-2}$
BP	1479	216	1.98 (1.24)	1300	-339	1.46	1.79	7.87
GaP	2490	216	(1.59)	1300	-368	1.14	2.22	11.88
InP	20351	216	(0.47)	1300	-113	0.73	6.67	1.76
FeP	1005	62	0.00 (0.00)	1300	-39	0.09	38.77	1.69
ZnNaP	4824	129	(0.91)	1300	-288	1.29	0.86	6.43
SrLiP	10614	187	(0.8)	1000	-345	4.04	0.61	2.95
BaLiP	10615	187	(0.7)	1100	-338	2.6	0.42	4.39
SnNaP	29529	186	(0.58)	900	-257	1.24	1.36	5.33
NiP <sub>2</sub>	22619	205	0.13 (0.00)	800	-69	0.17	21.50	2.8
NiP <sub>2</sub>	486	15	1.15 (0.53)	800	-242	1.03	1.76	5.69
ZnP <sub>2</sub>	11025	96	2.20 (1.46)	1300	-333	5.56	0.73	1.99
ZnP <sub>2</sub>	2782	92	2.20 (1.46)	1300	-333	5.57	0.73	1.99
AgP <sub>2</sub>	8200	14	1.30 (0.66)	1000	-325	2.56	0.81	4.13
TiP <sub>2</sub>	20327	62	0.00 (0.00)	1100	-34	0.1	37.03	1.16
CuP <sub>2</sub>	927	14	(0.86)	1100	-350	3.29	0.98	3.72
RuP <sub>2</sub>	1413	58	(0.48)	800	-274	2.86	0.70	2.63
CoP <sub>2</sub>	14285	14	(0.44)	600	-390	4.4	0.27	3.46
IrP <sub>2</sub>	10155	14	(0.78)	1000	-264	4.28	0.68	1.63
BeP <sub>2</sub>	27148	15	(0.83)	1300	-334	2.46	1.52	4.53
RhP <sub>2</sub>	15953	14	(0.37)	600	-322	6.58	0.32	1.58
PdP <sub>2</sub>	28266	15	(0.46)	700	-248	1.15	1.02	5.35
FeP <sub>2</sub>	20027	58	(0.43)	700	-293	4.73	0.67	1.81
Zn <sub>3</sub> P <sub>2</sub>	2071	137	1.17 (0.31)	1100	-201	1.39	3.03	2.91
Mg <sub>3</sub> P <sub>2</sub>	8085	224	(1.11)	1300	-295	1.61	1.93	5.41
Mg <sub>3</sub> P <sub>2</sub>	2514	206	2.20 (1.61)	1300	-242	1.51	1.70	3.88
Cd <sub>3</sub> P <sub>2</sub>	2441	137	0.65 (0.002)	1200	-142	0.77	5.72	2.62
MnP <sub>4</sub>	569522	2	(0.48)	700	-351	14.66	0.12	0.84
MnP <sub>4</sub>	487	2	1.58 (0.44)	700	-362	4.43	0.39	2.96
PbTe	19717	225	(0.81)	1300	-343	0.95	5.35	12.38
Bi <sub>2</sub> Te <sub>3</sub>	568390	166	(0.25)	1100	-283	1.04	3.07	7.7

**Table S3:** Calculated power factors in  $x$ ,  $y$ , and  $z$ -direction in the units  $\text{mW m}^{-1} \text{K}^{-2}$  at  $n = 10^{20} \text{ cm}^{-3}$  and  $T = 600 \text{ K}$ . Note that a constant relaxation time of  $10^{-14} \text{ s}$  was used for the electronic transport property calculations.

Material	NiP <sub>2</sub>	CoP <sub>2</sub>	MnP <sub>4</sub>	PdP <sub>2</sub>	FeP <sub>2</sub>	TiP <sub>2</sub>	RuP <sub>2</sub>	SrLiP	SnNaP	ZnNaP	BaLiP
mp-id	486	14258	487	28266	20027	20327	1413	10614	29529	4824	10615
$n\text{-}PF_{xx}$	1.36	1.54	2.47	2.71	0.96	0.25	2.04	1.66	2.72	2.28	1.37
$n\text{-}PF_{yy}$	6.67	3.48	4.05	3.99	1.88	2.04	2.83	1.66	4.18	2.28	1.37
$n\text{-}PF_{zz}$	2.48	5.35	2.28	6.06	2.55	0.95	1.15	3.28	4.18	0.33	2.21
$p\text{-}PF_{xx}$	1.37	2.01	1.85	2.73	1.33	0.45	3.05	2.38	0.84	6.39	1.57
$p\text{-}PF_{yy}$	5.73	2.89	1.93	4.87	0.95	1.80	3.51	2.38	0.84	6.39	1.57
$p\text{-}PF_{zz}$	5.08	2.97	1.79	3.25	1.97	0.89	4.86	6.90	7.70	0.36	3.14

**Table S4:** Computed *p*-type (*p*-PF) and *n*-type (*n*-PF) power factors of metal phosphorous compounds at 600 K and  $10^{20} \text{ cm}^{-3}$ . The electronic transport properties were calculated from the HSE and PBE-GGA band structures, respectively, with the Boltztrap code where the constant relaxation time was set to  $10^{-14} \text{ s}$ .

Material	mp-id	Space group	<i>p</i> -PF / $\text{mW m}^{-1} \text{ K}^{-2}$ (HSE)	<i>n</i> -PF / $\text{mW m}^{-1} \text{ K}^{-2}$ (HSE)	<i>p</i> -PF / $\text{mW m}^{-1} \text{ K}^{-2}$ (PBE-GGA)	<i>n</i> -PF / $\text{mW m}^{-1} \text{ K}^{-2}$ (PBE-GGA)
BP	1479	216	3.48	4.32	4.07	4.19
FeP	1005	62	0.0588	0.0517	0.838	0.847
NiP <sub>2</sub>	22619	205	4.41	0.674	3.21	2.01
NiP <sub>2</sub>	486	15	2.60	2.77	5.41	3.37
ZnP <sub>2</sub>	11025	96	2.52	1.22	2.08	0.922
ZnP <sub>2</sub>	2782	92	2.52	1.22	2.08	0.921
AgP <sub>2</sub>	8200	14	1.66	1.79	2.19	2.92
TiP <sub>2</sub>	20327	62	0.366	0.362	1.05	1.06
Zn <sub>3</sub> P <sub>2</sub>	2071	137	1.86	1.40	3.59	1.14
Mg <sub>3</sub> P <sub>2</sub>	2514	206	1.56	1.66	2.09	1.44
Cd <sub>3</sub> P <sub>2</sub>	2441	137	2.05	0.36	2.54	0.70
MnP <sub>4</sub>	487	2	2.25	2.60	2.09	1.44

**Table S5:** Minimum thermal conductivity using different models. Using static disordering and an average phonon mean speed indicates the amorphous limit. The minimum thermal conductivities using the Cahill-Pohl model where the speed of sounds were calculated from bulk and shear moduli predict the highest values and were taken to calculate the theoretical  $zT_{calc}$  in Figure 5.

	Material	$\kappa_{\min}$ (Cahill-Pohl, bulk)	$\kappa_{\min}$ (Cahill-Pohl, phonon)	$\kappa_{\min}$ (static disorder, avg. phonon)	$\kappa_{\min}$ (dynamic disorder, avg. phonon)
mp-22619	NiP <sub>2</sub>	1.6	1.36	0.81	1.22
mp-486	NiP <sub>2</sub>	1.32	1.06	0.53	0.8
mp-1479	BP	2.81	2.43	0.89	1.34
mp-2490	GaP	0.97	0.68	0.27	0.4
mp-8200	AgP <sub>2</sub>	0.95	0.58	0.39	0.58
mp-2071	Zn <sub>3</sub> P <sub>2</sub>	0.92	0.71	0.34	0.51
mp-2782	ZnP <sub>2</sub>	1.14	0.86	0.39	0.58
mp-11025	ZnP <sub>2</sub>	0.97	0.86	0.47	0.71
mg-2514	Mg <sub>3</sub> P <sub>2</sub>	1.07	0.92	0.44	0.67
mp-487	MnP <sub>4</sub>	1.77	1.57	0.91	1.41

**Table S6:** Comparison of computed thermal conductivity using the amorphous limit by Cahill and Pohl ( $\kappa_{\min}$ )<sup>[2,3]</sup> and using a semi-empirical approach described by Miller *et al.* ( $\kappa_{\text{SE}}$ ) at 300 K.<sup>[11]</sup> All studied XYPs have a low estimated thermal conductivity ( $< 2.5 \text{ W m}^{-1} \text{ K}^{-1}$ ) using both approaches pointing to excellent high-performance TE materials.

Material	mp-id	Space group	$\kappa_{\min} / \text{W m}^{-1} \text{ K}^{-1}$	$\kappa_{\text{SE}} / \text{W m}^{-1} \text{ K}^{-1}$
BP	1479	216	2.81	243.50
GaP	2490	216	0.97	47.74
InP	20351	216	0.62	27.10
FeP	1005	62	1.73	4.85
ZnNaP	4824	129	0.80	1.84
SrLiP	10614	187	0.79	2.12
BaLiP	10615	187	0.60	1.55
SnNaP	29529	186	0.52	1.75
<i>c</i> -NiP <sub>2</sub>	22619	205	1.60	12.73
<i>m</i> -NiP <sub>2</sub>	486	15	1.32	14.26
ZnP <sub>2</sub>	11025	96	0.97	6.36
ZnP <sub>2</sub>	2782	92	1.14	6.36
AgP <sub>2</sub>	8200	14	0.95	4.18
TiP <sub>2</sub>	20327	62	1.60	4.88
CuP <sub>2</sub>	927	14	1.12	6.90
RuP <sub>2</sub>	1413	58	1.66	13.2
CoP <sub>2</sub>	14285	14	1.78	9.69
IrP <sub>2</sub>	10155	14	1.24	10.15
BeP <sub>2</sub>	27148	15	9.19	10.90
RhP <sub>2</sub>	15953	14	1.51	9.65
PdP <sub>2</sub>	28266	15	1.03	13.12
FeP <sub>2</sub>	20027	58	2.00	13.36
Zn <sub>3</sub> P <sub>2</sub>	2071	137	0.92	2.60
Mg <sub>3</sub> P <sub>2</sub>	8085	224	1.14	4.18
Mg <sub>3</sub> P <sub>2</sub>	2514	206	1.07	4.21
Cd <sub>3</sub> P <sub>2</sub>	2441	137	0.52	1.84
MnP <sub>4</sub>	569522	2	1.78	8.44
MnP <sub>4</sub>	487	2	1.77	20.65
PbTe	19717	225	0.35	2.26
Bi <sub>2</sub> Te <sub>3</sub>	568390	166	0.22	0.56

**Table S7:**  $zT_{calc}$  for different  $p$ -type metal phosphorous compounds at  $n = 10^{20} \text{ cm}^{-3}$  and the temperature where the predicted  $zT_{calc}$  ( $\kappa_{min}$ ) is maximum. The thermal conductivity is calculated using the amorphous limit by Cahill and Pohl ( $\kappa_{min}$ )<sup>[2,3]</sup> and a semi-empirical approach by Miller *et al* ( $\kappa_{SE}$ ).<sup>[10]</sup> The electronic transport properties were calculated from the PBE band structures with the Boltztrap code where the constant relaxation time was set to  $10^{-14} \text{ s}$ .

Material	mp-id	Space group	$T / \text{K}$	$zT_{calc} (\kappa_{min})$	$zT_{calc} (\kappa_{SE})$
BP	1479	216	1300	2.11	0.16
GaP	2490	216	1300	3.74	0.68
InP	20351	216	1100	3.24	0.66
FeP	1005	62	1300	0.06	0.05
ZnNaP	4824	129	1300	3.27	2.40
SrLiP	10614	187	1000	2.75	2.05
BaLiP	10615	187	1100	6.88	3.87
SnNaP	29529	186	800	5.69	2.42
NiP <sub>2</sub>	22619	205	700	0.15	0.12
NiP <sub>2</sub>	486	15	500	1.89	0.33
ZnP <sub>2</sub>	11025	96	1300	1.94	0.62
ZnP <sub>2</sub>	2782	92	1300	1.67	0.62
AgP <sub>2</sub>	8200	14	1100	1.93	1.04
TiP <sub>2</sub>	20327	62	1100	0.03	0.03
CuP <sub>2</sub>	927	14	1100	2.59	1.10
RuP <sub>2</sub>	1413	58	800	1.38	0.55
CoP <sub>2</sub>	14285	14	700	0.93	0.38
IrP <sub>2</sub>	10155	14	1100	2.52	1.07
BeP <sub>2</sub>	27148	15	1300	0.56	0.90
RhP <sub>2</sub>	15953	14	800	1.23	0.56
PdP <sub>2</sub>	28266	15	600	1.43	0.29
FeP <sub>2</sub>	20027	58	500	0.36	0.10
Zn <sub>3</sub> P <sub>2</sub>	2071	137	900	2.01	1.15
Mg <sub>3</sub> P <sub>2</sub>	8085	224	1200	2.69	1.19
Mg <sub>3</sub> P <sub>2</sub>	2514	206	1300	2.10	0.75
Cd <sub>3</sub> P <sub>2</sub>	2441	137	600	1.10	0.60
MnP <sub>4</sub>	569522	2	700	0.35	0.12
MnP <sub>4</sub>	487	2	700	0.66	0.14

PbTe	19717	225	1100	7.78	6.19
Bi <sub>2</sub> Te <sub>3</sub>	568390	166	1000	3.01	2.69

**Table S8:**  $zT_{calc}$  for different  $n$ -type metal phosphorous compounds at  $n = 10^{20} \text{ cm}^{-3}$  and the temperature where the predicted  $zT_{calc} (\kappa_{min})$  is maximum. The thermal conductivity is calculated using the amorphous limit by Cahill and Pohl ( $\kappa_{min}$ )<sup>[2,3]</sup> and a semi-empirical approach by Miller *et al* ( $\kappa_{SE}$ ).<sup>[10]</sup> The electronic transport properties were calculated from the PBE band structures with the Boltztrap code where the constant relaxation time was set to  $10^{-14} \text{ s}$ .

Material	mp-id	Space group	$T / \text{K}$	$zT_{calc} (\kappa_{min})$	$zT_{calc} (\kappa_{SE})$
BP	1479	216	1300	2.23	0.17
GaP	2490	216	1300	4.84	0.92
InP	20351	216	1300	0.31	0.17
FeP	1005	62	1300	0.06	0.05
ZnNaP	4824	129	1300	5.06	3.39
SrLiP	10614	187	1000	2.11	1.45
BaLiP	10615	187	1100	4.71	3.27
SnNaP	29529	186	800	2.55	1.85
NiP <sub>2</sub>	22619	205	700	0.10	0.08
NiP <sub>2</sub>	486	15	500	1.49	0.59
ZnP <sub>2</sub>	11025	96	1300	1.52	0.63
ZnP <sub>2</sub>	2782	92	1300	1.39	0.63
AgP <sub>2</sub>	8200	14	1100	2.34	1.27
TiP <sub>2</sub>	20327	62	1100	0.03	0.03
CuP <sub>2</sub>	927	14	1100	1.96	0.93
RuP <sub>2</sub>	1413	58	800	0.89	0.33
CoP <sub>2</sub>	14285	14	700	1.01	0.36
IrP <sub>2</sub>	10155	14	1100	0.85	0.31
BeP <sub>2</sub>	27148	15	1300	0.55	0.87
RhP <sub>2</sub>	15953	14	800	0.52	0.16
PdP <sub>2</sub>	28266	15	600	1.82	0.55
FeP <sub>2</sub>	20027	58	500	0.48	0.16
Zn <sub>3</sub> P <sub>2</sub>	2071	137	900	0.81	0.59
Mg <sub>3</sub> P <sub>2</sub>	8085	224	1200	2.28	1.46
Mg <sub>3</sub> P <sub>2</sub>	2514	206	1300	1.82	1.02
Cd <sub>3</sub> P <sub>2</sub>	2441	137	600	0.50	0.43
MnP <sub>4</sub>	569522	2	700	0.31	0.10
MnP <sub>4</sub>	487	2	700	0.96	0.20

PbTe	19717	225	1100	2.83	2.60
Bi <sub>2</sub> Te <sub>3</sub>	568390	166	1000	2.58	2.38

**Table S9:** Computed charge transfer between the metal and phosphorous atoms. While most compounds have only a small charge transfer between the atoms, the *XYP* subclass has a relative large charge transfer.

Material	mp	Space group	Charge Transfer
BP	1479	216	B: +0.40; P: -0.40
GaP	2490	216	Ga: +0.78; P: -0.78
FeP	1005	62	Fe: +0.31; P: -0.31
ZnNaP	4824	129	Na: +0.81; Zn: +0.51; P: -1.32
SrLiP	10614	187	Sr: +1.33; Li: +0.81; P: -2.14
BaLiP	10615	187	Ba: +1.20; Li: +0.82; P: -2.02
SnNaP	29529	186	Na: +0.81; Sn: +0.32; P: -1.13
NiP <sub>2</sub>	22619	205	Ni: +0.11; P <sub>1</sub> : -0.04; P <sub>2</sub> : -0.07
NiP <sub>2</sub>	486	15	Ni: +0.13; P <sub>1</sub> : -0.06; P <sub>2</sub> : -0.07
ZnP <sub>2</sub>	11025	96	Zn: +0.59; P <sub>1</sub> : -0.29; P <sub>2</sub> : -0.30
ZnP <sub>2</sub>	2782	92	Zn: +0.59; P <sub>1</sub> : -0.29; P <sub>2</sub> : -0.30
AgP <sub>2</sub>	8200	14	Ag: +0.19; P <sub>1</sub> : -0.03; P <sub>2</sub> : -0.17
TiP <sub>2</sub>	20327	62	Ti: +1.36; P <sub>1</sub> : -0.84; P <sub>2</sub> : -0.52
CuP <sub>2</sub>	927	14	Cu: +0.27; P <sub>1</sub> : -0.04; P <sub>2</sub> : -0.23
RuP <sub>2</sub>	1413	58	Ru: +0.07; P <sub>1</sub> : -0.03; P <sub>2</sub> : -0.04
CoP <sub>2</sub>	14285	14	Co: +0.09; P <sub>1</sub> : -0.05; P <sub>2</sub> : -0.05
IrP <sub>2</sub>	10155	14	Ir: -0.16; P <sub>1</sub> : +0.09; P <sub>2</sub> : +0.07
BeP <sub>2</sub>	27148	15	Be: +1.51; P <sub>1</sub> : -0.76; P <sub>2</sub> : -0.76
RhP <sub>2</sub>	15953	14	Rh: -0.04; P <sub>1</sub> : +0.01; P <sub>2</sub> : +0.04
PdP <sub>2</sub>	28266	15	Pd: -0.05; P <sub>1</sub> : +0.02; P <sub>2</sub> : +0.03
FeP <sub>2</sub>	20027	58	Fe: +0.20; P <sub>1</sub> : -0.08; P <sub>2</sub> : -0.12
Zn <sub>3</sub> P <sub>2</sub>	2071	137	Zn <sub>1</sub> : +0.61; Zn <sub>2</sub> : +0.62; Zn <sub>3</sub> : +0.63; P <sub>1</sub> : -0.91; P <sub>2</sub> : -0.95
Mg <sub>3</sub> P <sub>2</sub>	8085	224	Mg <sub>1</sub> : +1.53; Mg <sub>2</sub> : +1.53; Mg <sub>3</sub> : +1.53; P <sub>1</sub> : -2.29; P <sub>2</sub> : -2.29
Mg <sub>3</sub> P <sub>2</sub>	2514	206	Mg <sub>1</sub> : +1.56; Mg <sub>2</sub> : +1.56; Mg <sub>3</sub> : +1.56; P <sub>1</sub> : -2.34; P <sub>2</sub> : -2.34
Cd <sub>3</sub> P <sub>2</sub>	2441	137	Cd <sub>1</sub> : +0.57; Cd <sub>2</sub> : +0.57; Cd <sub>3</sub> : +0.57; P <sub>1</sub> : -0.85; P <sub>2</sub> : -0.87

MnP <sub>4</sub>	487	2	Mn: +0.47; P <sub>1</sub> : -0.06; P <sub>2</sub> : -0.07; P <sub>3</sub> : -0.14; P <sub>4</sub> : -0.20
MnP <sub>4</sub>	569522	2	Mn: +0.47; P <sub>1</sub> : -0.06; P <sub>2</sub> : -0.07; P <sub>3</sub> : -0.14; P <sub>4</sub> : -0.20

**Table S10:** Rietveld refined PXRD data for *c*-NiP<sub>2</sub>.

Parameter	<i>c</i> -NiP <sub>2</sub>
Space group	<i>Pa</i> $\bar{3}$ (No. 205)
<i>a</i> / Å	5.469
<i>b</i> / Å	5.469
<i>c</i> / Å	5.469
$\alpha$ / °	90.000
$\beta$ / °	90.000
$\gamma$ / °	90.000
Number of molecules per unit cell, <i>Z</i>	4
Volume / Å <sup>3</sup>	164
Calculated density / g cm <sup>-3</sup>	4.90
<i>R<sub>p</sub></i>	11.21
<i>R<sub>wp</sub></i>	16.96
$\chi^2$	1.65

Atom	Site	<i>x</i>	<i>y</i>	<i>z</i>	Occupation
Ni	4 <i>a</i>	0.000	0.000	0.000	1.000
P	8 <i>c</i>	0.385	0.385	0.385	2.001

**Table S11:** Heat capacity data for *c*-NiP<sub>2</sub> using <sup>4</sup>He cooling: 27.82±0.01 mg; in order of data collection.

<i>T</i> / K	<i>C<sub>p</sub></i> / J K <sup>-1</sup> mol <sup>-1</sup>	<i>T</i> / K	<i>C<sub>p</sub></i> / J K <sup>-1</sup> mol <sup>-1</sup>	<i>T</i> / K	<i>C<sub>p</sub></i> / J K <sup>-1</sup> mol <sup>-1</sup>
303.64	66.96	53.29	7.99	9.96	0.13
283.41	65.83	48.28	6.32	9.02	0.11
263.26	63.06	43.74	4.96	8.17	0.09
243.06	60.55	39.62	3.87	7.40	0.08
222.89	57.95	35.91	2.99	6.70	0.07
202.69	54.63	32.53	2.31	6.08	0.06
182.48	50.84	29.48	1.78	5.49	0.06
162.27	46.19	26.71	1.38	4.98	0.05
142.05	40.57	24.21	1.05	4.51	0.05
121.79	34.15	21.94	0.82	4.08	0.04
101.57	26.67	19.89	0.64	3.70	0.04
96.53	24.58	18.02	0.49	3.35	0.04
87.45	20.71	16.33	0.39	3.03	0.04
79.24	17.70	14.78	0.31	2.74	0.04
71.75	14.92	13.40	0.25	2.49	0.03
65.06	12.09	12.14	0.20	2.25	0.03
58.84	9.97	11.00	0.16	2.04	0.03

**Table S12:** Heat capacity data for *c*-NiP<sub>2</sub> using <sup>3</sup>He cooling: 21.01±0.01 mg; in order of data collection.

<i>T</i> / K	<i>C<sub>p</sub></i> / J K <sup>-1</sup> mol <sup>-1</sup>	<i>T</i> / K	<i>C<sub>p</sub></i> / J K <sup>-1</sup> mol <sup>-1</sup>	<i>T</i> / K	<i>C<sub>p</sub></i> / J K <sup>-1</sup> mol <sup>-1</sup>
301.88	69.855	3.45	0.032	1.10	0.016
10.05	0.125	3.17	0.031	1.02	0.015
9.26	0.106	2.92	0.029	0.94	0.014
8.52	0.091	2.69	0.028	0.87	0.013
7.85	0.079	2.48	0.027	0.80	0.012
7.22	0.069	2.29	0.026	0.74	0.011
6.65	0.060	2.11	0.025	0.69	0.010
6.13	0.054	1.94	0.024	0.63	0.009
5.65	0.049	1.79	0.023	0.59	0.009
5.20	0.045	1.65	0.022	0.54	0.008
4.79	0.041	1.52	0.020	0.50	0.008

4.41	0.039	1.40	0.019	0.46	0.008
4.07	0.036	1.29	0.018	0.43	0.007
3.74	0.034	1.19	0.017		

**Table S13:** Experimental electrical resistivity ( $\rho$ ) for  $c$ -NiP<sub>2</sub>; in order of data collection. Low-temperature (LT) measurements were conducted with a two-probe method in a Quantum Design PPMS while high-temperature (HT)  $\rho$  were measured using the van der Pauw technique. It is important to note that the LT data were empirically adjusted to the HT data at  $T = 312$  K, using a contact resistance (10 m $\Omega$ ;  $\rho_{\text{contact}} = 24$  m $\Omega$  cm). HT measurements have an uncertainty of 5%.

$T / \text{K}$	$\rho$ (LT) /m $\Omega$ cm	$T / \text{K}$	$\rho$ (HT) /m $\Omega$ cm	$T / \text{K}$	$\rho$ (HT) /m $\Omega$ cm	$T / \text{K}$	$\rho$ (HT) /m $\Omega$ cm	$T / \text{K}$	$\rho$ (HT) /m $\Omega$ cm
312.27	9.11	295.25	9.05	406.55	8.21	516.15	7.39	626.15	6.37
292.29	9.63	298.45	9.01	411.45	8.17	521.05	7.34	630.95	6.33
272.33	10.21	299.65	8.99	415.95	8.14	525.85	7.31	635.65	6.28
252.40	10.72	304.75	8.93	420.75	8.1	530.45	7.27	640.45	6.24
226.92	11.29	314.05	8.86	425.85	8.06	535.35	7.23	645.45	6.19
202.04	11.74	321.85	8.82	430.35	8.03	540.35	7.18	649.95	6.15
176.81	12.13	326.15	8.8	435.25	7.99	544.85	7.14	654.55	6.11
151.93	12.43	329.75	8.77	439.85	7.96	549.85	7.1	659.45	6.06
126.76	12.67	334.15	8.73	444.75	7.92	554.35	7.06	664.25	6.01
101.57	12.85	339.45	8.7	449.55	7.89	559.15	7.02	669.05	5.96
81.35	13.03	344.65	8.66	454.25	7.86	563.95	6.97	673.75	5.92
51.09	13.19	349.55	8.62	459.15	7.83	568.55	6.93	678.45	5.87
31.10	13.25	354.25	8.59	463.75	7.8	573.65	6.88	683.35	5.82
20.58	13.28	358.85	8.55	468.65	7.76	578.35	6.84	687.95	5.77
18.55	13.31	363.65	8.52	473.65	7.72	583.15	6.79	692.85	5.71
15.60	13.31	368.65	8.48	478.05	7.69	587.85	6.75	697.65	5.66
10.50	13.36	373.05	8.45	482.85	7.65	592.75	6.7	702.35	5.6
8.29	13.39	377.95	8.42	487.65	7.61	597.55	6.66	707.05	5.55
6.28	13.42	382.75	8.38	492.35	7.58	602.25	6.61	712.05	5.49
4.35	13.48	387.45	8.35	497.15	7.54	607.15	6.56	716.75	5.43
2.41	13.59	392.25	8.32	501.85	7.5	611.65	6.51	721.55	5.37
		397.15	8.28	506.75	7.46	616.75	6.47	726.15	5.31
		401.75	8.25	511.45	7.42	621.35	6.42		

**Table S14:** Experimental Seebeck coefficient ( $S$ ) for  $c$ -NiP<sub>2</sub>; in order of data collection. Low-temperature (LT) measurements were conducted in a Quantum Design PPMS while high-temperature (HT)  $S$  were measured with an oscillating temperature gradient. LT and HT measurements have an uncertainty of 5%.

$T / \text{K}$	$S \text{ (LT)} / \mu\text{V K}^{-1}$	$T / \text{K}$	$S \text{ (LT)} / \mu\text{V K}^{-1}$	$T / \text{K}$	$S \text{ (LT)} / \mu\text{V K}^{-1}$	$T / \text{K}$	$S \text{ (HT)} / \mu\text{V K}^{-1}$
391.86	-40.86	176.81	-29.24	10.50	-0.86	324.55	-37.63
371.97	-40.73	151.93	-26.53	8.29	-2.54	369.63	-40.98
352.10	-41.00	126.76	-23.54	4.35	-1.73	415.88	-44.27
332.22	-39.67	101.57	-16.15	2.41	-1.41	462.83	-47.51
312.28	-39.06	81.35	-15.25			511.20	-50.75
292.29	-39.48	51.09	-10.06			559.07	-55.08
272.33	-37.66	31.10	-4.79			606.79	-57.65
252.40	-36.56	20.58	-2.89			654.87	-55.60
226.92	-34.63	18.55	-3.97			701.40	-51.51
202.04	-32.18	15.60	-2.02			748.50	-44.77

**Table S15:** Experimental thermal conductivity ( $\kappa$ ) for  $c$ -NiP<sub>2</sub>; in order of data collection. Low-temperature (LT) measurements were conducted in a Quantum Design PPMS where the thermal conductivity was corrected for blackbody radiation with an emissivity of 1. High-temperature (HT)  $\kappa$  was calculated from the thermal diffusivity ( $\alpha$ ), the density ( $d$ ), and the computed heat capacity ( $C_p$ ). While the LT measurement has an uncertainty of 5%, the HT measurement has an uncertainty of 15%.

$T / \text{K}$	$\kappa \text{ (LT)} / \text{W m}^{-1} \text{ K}^{-1}$	$T / \text{K}$	$\kappa \text{ (LT)} / \text{W m}^{-1} \text{ K}^{-1}$	$T / \text{K}$	$\alpha / \text{mm}^2 \text{ s}^{-1}$	$d / \text{g cm}^{-3}$	$C_p / \text{J g}^{-1} \text{ K}^{-1}$	$\kappa \text{ (HT)} / \text{W m}^{-1} \text{ K}^{-1}$
312.28	1.16	51.09	0.33	323.45	0.40	4.04	0.57	0.92
292.29	1.12	31.10	0.17	375.15	0.41	4.04	0.59	0.97
272.33	1.08	20.58	0.09	425.45	0.41	4.04	0.61	1.00
252.40	1.05	18.56	0.08	475.15	0.42	4.04	0.62	1.04
226.92	0.99	15.62	0.06	524.95	0.43	4.04	0.62	1.07
202.04	0.94	10.51	0.03	574.55	0.44	4.04	0.63	1.11
176.81	0.88	8.30	0.02	624.35	0.45	4.04	0.63	1.16
151.93	0.86	6.30	0.01	674.25	0.47	4.04	0.63	1.20
126.76	0.75	4.36	0.01	724.15	0.48	4.04	0.64	1.24
101.57	0.62	2.42	0.00					

81.35	0.52							
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