Supporting information

Field enhancement of MoS$_2$: visualization of enhancement and effect of the number of layers

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**Table S1** Reported thicknesses of MoS$_2$ monolayer using four preparation methods

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<tr>
<th>Thickness of monolayer (nm)</th>
<th>Procedures</th>
<th>References</th>
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<td>1.9</td>
<td>sonication</td>
<td><em>Nat Commun.</em>, 2014, <strong>5</strong>, 4576.</td>
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<td>1.5</td>
<td></td>
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<td>0.9-1.2</td>
<td></td>
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<td>1.2</td>
<td></td>
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<td>1.1</td>
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<td>0.7</td>
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<td>0.65-0.7</td>
<td>mechanical</td>
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<td>0.53</td>
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<td>0.72</td>
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*a* Exfoliation by a sonication method in a solvent. *b* Exfoliation by the scotch tape. *c* Exfoliation by a lithium intercalation method, e.g. butyllithium. This method may cause phase transition to a metastable phase where the coordination of Mo atoms becomes octahedral, and the structure has metallic characteristic. *d* Synthesis from chemical vapor deposition (CVD) method.
Fig. S1 EF mappings of electric fields obtained from FDTD calculations of a MoS$_2$ flake, whose x-y-z-positions are determined from AFM measurement. The calculations are conducted as a function of the depth of the MoS$_2$ flake, i.e. from the surface at $d = 0$ to the depth of $d = 100$ nm at the interval of 12.5 nm. The summation of the above nine calculated results gives Fig.4a. The scale bars denote 1µm. $E$ and $E_0$ represent the magnitudes of electric fields with MoS$_2$ and without MoS$_2$, respectively.
Fig. S2 (a) Laser microscope image of MoS₂ flakes. (b)(c) AFM images of MoS₂ flakes. (d)–(f) EF mappings are obtained from enhanced fluorescence intensities of CV using MoS₂ flakes at the same surfaces of the MoS₂ flakes in Fig. 2a–2c respectively. The bars denote the length of 5 µm. As for the grid sizes, the data of laser microscope, AFM and fluorescence micro spectroscopy images are set as 250 nm, 100 nm, and 500 nm, respectively.
Fig. S3 Schematic diagram of a charge transfer (CT) band observed at around 720 nm in Fig. 8c. Black lines and navy rectangular represent electronic states of a crystal violet (CV) molecule and the band structure of bulk MoS$_2$, respectively. The former and latter energy levels are used as the data in ref. 1 and refs. 2 and 3, respectively; the highest occupied molecular orbital (HOMO), the lowest occupied molecular orbital (LUMO), conduction band (CB), and valence band (VB). The black arrows denote an optical transition of an electron of CV and recombination process between electron and hole via a CT process.

Fig. S4 Luminescence spectra measured with a MoS$_2$ (ca.800 layers) flake at an excitation wavelength of 632.8 nm. Red and black curves denote the fluorescence spectrum of CV solution without MoS$_2$ and photoluminescence & Raman spectra of MoS$_2$, respectively.
**Fig. S5** Fluorescence spectra of CV solution with MoS$_2$ ranging from the number of layers ($#$) =10 to 50 (color curves). Fluorescence spectrum of CV solution without MoS$_2$ (black curve).

**Fig. S6** Fluorescence spectra of CV solution with MoS$_2$ ranging from the number of layers ($#$) = 600 to 700 (color curves). Fluorescence spectrum of CV solution without MoS$_2$ (black curve).
Calculation of the extinction efficiency ($Q_{\text{ext}}$)

We obtained the extinction efficiency ($Q_{\text{ext}}$) of MoS$_2$ using the Mie theory$^5$, based on the following equations. $Q_{\text{ext}}$ is expressed by $Q_{\text{ext}} = Q_{\text{abs}} + Q_{\text{sca}}$, where $Q_{\text{abs}}$ is absorption efficiency and $Q_{\text{sca}}$ the scattering efficiency. The $Q_{\text{ext}}$ and $Q_{\text{sca}}$ components are formulated as follows:

\begin{align*}
Q_{\text{ext}} &= \frac{2}{x^2} \sum_{N=1}^{\infty} (2N + 1) \Re[a_N + b_N], \quad (S1) \\
Q_{\text{sca}} &= \frac{2}{x^2} \sum_{N=1}^{\infty} (2N + 1) [|a_N|^2 + |b_N|^2], \quad (S2)
\end{align*}

where $N$ is an integer, i.e. $N=1, 2, \text{and } 3$ correspond to electro and magnetic dipole, quadrupole, and hexapole modes, respectively. $a_N$ (electric component) and $b_N$ (magnetic component) are expressed as follows:

\begin{align*}
a_N &= \frac{m \Psi_N(mx) \Psi'_N(x) - \Psi_N(x) \Psi'_N(mx)}{m \Psi_N(mx) \xi_N(x) - \xi_N(x) \Psi'_N(mx)}, \quad (S3) \\
b_N &= \frac{\Psi_N(mx) \Psi'_N(x) - m \Psi_N(x) \Psi'_N(mx)}{\Psi_N(mx) \xi_N(x) - m \xi_N(x) \Psi'_N(mx)}, \quad (S4)
\end{align*}

where $\Psi_N$ and $\xi_N$ are the Riccati-Bessel functions, $\Psi'_N$ and $\xi'_N$ are the first differentiations, and $x = 2\pi n m R/\lambda$ and $m = (n + ik)/n_m$, where $n$ is the real part of refractive index of a particle, $R$ the radius of the particle, $k$ the imaginary part of the refractive index of the particle (MoS$_2$), $\lambda$ excitation wavelength, and $n_m$ the refractive index of the medium around the particle. We used refractive indices of MoS$_2$ at 632.8 nm ($n = 5.22$, $k = 1.08$) from ref. 4 and $n_m = 1.33$ of methanol. In the main manuscript, the calculations use not only the dipole term but also the quadrupole term for bulk MoS$_2$ (i.e., $N = 1 \text{ and } 2$).

References: