Ultrahigh electron mobility induced by strain engineering in directly semiconducting monolayer Bi₂TeSe₂

Supplementary

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I. 2D STRUCURE PREDICTION USING THE PARTICLE SWARM OPTIMIZATION METHOD

We use the particle swarm optimization (PSO) method within the evolutionary algorithm implemented in the Crystal structure AnaLYsis by Particle Swarm Optimization (CALYPSO) code\textsuperscript{1,2}, to conduct the global structure search for the global energy-minimum structure of the stoichiometric monolayer Bi\textsubscript{2}TeSe\textsubscript{2}. The unit cell containing two Bi atoms, one Te atom and two Se atoms is considered, and three sublayers with a vacuum gap of 30Å are set as the parameters for 2D structure prediction. In the first step, 20 random structures with certain crystal symmetry are constructed, in which the atomic coordinates are generated by the crystallographic symmetry. Then, we use VASP code to carefully optimize the random structures until the Gibbs free energy changes are smaller than $1 \times 10^{-5}$ eV. After finishing the structural optimization of all 20 random structures, 60% of them with lower Gibbs free energies are selected to partially construct the next-generation of PSO. The rest 40% new structures of the next-generation PSO are generated based on the so-called structure fingerprinting technique of bond characterization matrix to avoid identical structures\textsuperscript{2}.

In the case of monolayer Bi\textsubscript{2}TeSe\textsubscript{2}, the structure searching simulations stop at 10 generations and 200 structures are generated and under investigations. The information about the ten structures with lowest formation enthalpies is listed as Table S1.

<table>
<thead>
<tr>
<th>No.</th>
<th>1</th>
<th>2</th>
<th>3</th>
<th>4</th>
<th>5</th>
<th>6</th>
<th>7</th>
<th>8</th>
<th>9</th>
<th>10</th>
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<tbody>
<tr>
<td>Symmetry(Group)</td>
<td>Pmm2(25)</td>
<td>P-3m1(164)</td>
<td>R3m(160)</td>
<td>R3m(160)</td>
<td>R3m(160)</td>
<td>R3m(160)</td>
<td>Fm2m(42)</td>
<td>P3m1(156)</td>
<td>P1(1)</td>
<td>P1(1)</td>
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</tbody>
</table>

The ten structures with lowest formation enthalpies are shown in Table S1. The structure with the symmetry of Pmm2 as shown in Fig. S1(a) and the most lowest formation enthalpy predicted by CALYPSO as shown in Table S1, composes of three isolated monolayers with large distances between each other, so it is invalid and we neglect this structure. The second structure possesses the crystal symmetry of P-3m1, and it is shown in Fig. S1(b). By comparison, we find that, the second structure shown in Fig. S1(b) is the monolayer Bi\textsubscript{2}TeSe\textsubscript{2} we predicted by the atomic transmutation method in the present work.
FIG. S1: Two predicted structures with lowest formation enthalpies, corresponding to (a) No. 1 and (b) No. 2 in Table S1.

II. VDW CORRECTIONS TO THE BAND STRUCTURES, OPTICAL PROPERTIES AND CARRIER MOBILITIES

FIG. S2: Electronic band structures of monolayer Bi$_2$TeSe$_2$ calculated at the (a) PBE (b) PBE+SOC levels with/without vdW interaction.
FIG. S3: Calculations of (a) dielectric functions for in-plane polarization (E⊥c) and out-of-plane polarization (E//c), (b) absorption coefficients, (c) reflectivities of monolayer Bi₂TeSe₂ with HSE+vdW method, respectively.

TABLE S2: Calculated elastic modulus (C²D), deformation potential constant (D₁), effective mass (m₀), and mobility (μ) in x and y directions of monolayer Bi₂TeSe₂ with/without vdW at 300 K.

<table>
<thead>
<tr>
<th>Bi₂TeSe₂</th>
<th>Strain Direction</th>
<th>C²D(N/m)</th>
<th>D₁(e)</th>
<th>D₁(h)</th>
<th>m⁺(m₀)</th>
<th>m⁻(m₀)</th>
<th>μₑ(cm²V⁻¹s⁻¹)</th>
<th>μₕ(cm²V⁻¹s⁻¹)</th>
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<tr>
<td></td>
<td>- x</td>
<td>59.94</td>
<td>7.17</td>
<td>7.75</td>
<td>0.0812</td>
<td>0.1037</td>
<td>2510</td>
<td>1317</td>
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<tr>
<td>w- vdW</td>
<td>- x</td>
<td>61.22</td>
<td>7.32</td>
<td>8.28</td>
<td>0.0871</td>
<td>0.1029</td>
<td>2133</td>
<td>1198</td>
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<tr>
<td></td>
<td>- y</td>
<td>54.65</td>
<td>7.32</td>
<td>8.33</td>
<td>0.0812</td>
<td>0.1037</td>
<td>2195</td>
<td>1039</td>
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<tr>
<td>w- vdW</td>
<td>- y</td>
<td>56.03</td>
<td>7.73</td>
<td>8.27</td>
<td>0.0871</td>
<td>0.1029</td>
<td>1930</td>
<td>1099</td>
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<tr>
<td></td>
<td>-6% x</td>
<td>87.23</td>
<td>7.13</td>
<td>11.61</td>
<td>0.0338</td>
<td>0.0332</td>
<td>21353</td>
<td>8316</td>
</tr>
<tr>
<td>w- vdW</td>
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<td>83.54</td>
<td>7.15</td>
<td>11.54</td>
<td>0.0335</td>
<td>0.0330</td>
<td>20678</td>
<td>8181</td>
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</table>
III. ANALYSIS OF THE BAND STRUCTURES FOR MONOLAYER Bi$_2$TeSe$_2$

We employ the QUANTUM EXPRESSO code with both fully relativistic and scalar relativistic pseudopotentials to analyze the band structures for monolayer Bi$_2$TeSe$_2$ with and without SOC effects$^{3,4}$, as shown in Fig. S4. The irreducible representations (irreps) for energy levels near Fermi level at $\Gamma$ point are obtained by the QUANTUM EXPRESSO code as well and shown in Fig. S4. The small group for $\Gamma$ point is $D_{3d}$. The irreps for maximum of valence band and minimum of conduction band without SOC effects at $\Gamma$ point are two-dimensional (2D) $\Gamma_3^-$ with odd parity and one-dimensional (1D) $\Gamma_{1}^+$ with even parity, as shown in Fig. S4(a). The SOC effect splits the 2D $\Gamma_3^-$ into degenerate 1D $\Gamma_6^-$ and 1D $\Gamma_5^-$, and 2D $\Gamma_4^-$ in the double-group representations, as shown in Fig. S4(b). The SOC effect also transfers the 1D $\Gamma_{1}^+$ into 2D $\Gamma_4^+$. Therefore, the irreps for VBM and CBM in the double-group representations are degenerate $\Gamma_6^-$ & $\Gamma_5^-$ and $\Gamma_4^+$ respectively, as shown in Fig. S4(b). Since the order of parities does not change when considering SOC effects, no band inversion takes place at $\Gamma$ point for monolayer Bi$_2$TeSe$_2$, thus it is topologically trivial. We have further checked the topological properties for monolayer Bi$_2$TeSe$_2$ by using the state-of-art software SymTopo$^5$, and the calculated symmetry-based indicator is (0,0,0,0), which confirms that monolayer Bi$_2$TeSe$_2$ is a trivial insulator.

![FIG. S4: Band structures for monolayer Bi$_2$TeSe$_2$ (a) without and (b) with SOC effects.](image-url)
IV. SOC EFFECTS ON BANDGAPS AND MOBILITIES

Since Te is below Se in the periodic table, the SOC effect of Te atom is stronger than Se atoms. As shown in Table 1, the calculated SOC strength for Te atom is 140 meV smaller than Se atom, which means that the SOC effect of heavier Te atoms is stronger than Se atoms. To further investigate influences of the SOC effects on band structures and carrier mobilities, we change the ratios of Te and Se in monolayer Bi$_2$TeSe$_2$ with fixed crystal structures by replacing Se atoms with Te atoms, and monolayers of Bi$_2$Se$_3$, Bi$_2$TeSe$_2$, Bi$_2$Te$_2$Se and Bi$_2$Se$_3$ with identical crystal structures are generated. By using first-principle method, the calculated bandgaps for the four monolayers are shown in Fig. S5, which shows that, when the Te concentration increases, the value of bandgap decreases from 1.08 eV to 0.29 eV. The decreased bandgap is resulted from the stronger SOC strength of Te atoms, considering similar chemical properties of Te and Se atoms and identical crystal structures. According to the simple two-band $k\cdot p$ model shown as Eq. (12), the value of $m^*$ increases when the value of bandgap increases. Thus, when the value of bandgap decreases with increasing concentration of Te atoms, the value of $m^*$ increases, which finally decreases the value of carrier mobilities, according to the deformation-potential theory shown as Eq. (7).
V. CLEAVAGE ENERGY FOR Bi$_2$Te$_2$Se

Fig. S6 shows the resulting cleavage energy by determining the total energy difference (per unit cell) with respect to the ground state of the fractured structure with four quintuple layers (QLs) slab as a function of distance. And the calculated cleavage energy of Bi$_2$TeSe$_2$ is 0.37 J/m$^2$, which confirms that monolayer Bi$_2$TeSe$_2$ could be feasible to be experimentally exfoliated from its bulk single crystals considering the weak vdW interactions between adjacent QLs.

FIG. S6: Cleavage energy as a function of the distance between a top QL and the bottom three QLS of Bi$_2$TeSe$_2$. Inserted is the schematic view of the cleavage process of monolayer from the four-QLs slab of Bi$_2$TeSe$_2$.

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