### **Electronic Supplementary Information**

#### Magnetoimpedance spectroscopy of phase-separated La<sub>0.5</sub>Ca<sub>0.5</sub>MnO<sub>3</sub> polycrystalline manganites

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ESI 1. Neutron diffraction study of La<sub>0.5</sub>Ca<sub>0.5</sub>MnO<sub>3</sub> at room temperature

ESI 2. X-ray thermodiffraction study of La<sub>0.5</sub>Ca<sub>0.5</sub>MnO<sub>3</sub>

ESI 3. *M* vs *H* curves for La<sub>0.5</sub>Ca<sub>0.5</sub>MnO<sub>3</sub> powder between 5 K and 300 K

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# ESI 1. Neutron diffraction study of $La_{0.5}Ca_{0.5}MnO_3$ at room temperature

### Table ESI 1

T - 44°	5 4159(2)							
Lattice parameter, a / A	5.4158(2)							
Lattice parameter, b / A	7.6430(2)							
Lattice parameter, c / A	5.4247(2)							
La / Ca (position 4c)								
X	0.0189(7)							
Ζ	-0.0013(9)							
$\beta_{11} \times 10^4 / Å^2$	75(10)							
$\beta_{22} \times 10^4 / \text{\AA}^2$	29(5)							
$\beta_{33} \times 10^4 / \text{\AA}^2$	69(10)							
Occ. La / Ca (0.25/0.25)	0.241(8) / 0.247(8)							
Mn (position 4b)								
$\beta_{11} \times 10^4 / Å^2$	44(10)							
$\beta_{22} \times 10^4 / \text{\AA}^2$	10(5)							
$\beta_{33} \times 10^4 / \text{\AA}^2$	65(15)							
Occ. Mn (0.5)	0.492(6)							
O1 (position 4c)								
X	-0.008 (5)							
Ζ	0.4400(8)							
$\beta_{11}  x  10^4 / {\rm \AA}^2$	135(20)							
$\beta_{22} \times 10^4 / Å^2$	19(5)							
$\beta_{33} \times 10^4 / \text{\AA}^2$	90(10)							
Occ. O (0.5)	0.5							
O2 (position 8d)								
x	0.7236(6)							
y	-0.0306(3)							
Z	0.2762(7)							
$\beta_{11}  x  10^4  /  \text{\AA}^2$	73(9)							
$\beta_{22} \times 10^4 / \text{\AA}^2$	44(4)							
$\beta_{33} \times 10^4 / \text{\AA}^2$	117(10)							
Occ. O2 (1)	1							
$R_{p}(\%) / R_{wp}(\%)$	1.66 / 2.21							
R <sub>exp</sub> (%)	4.71							
Bragg R-factor	1.54							
<i>Pnma</i> (#62): 4c(x <sup>1</sup> / <sub>4</sub> z), 4b(0 0 <sup>1</sup> / <sub>2</sub> ), 8d(xyz)								

# ESI 2. X-ray thermodiffraction study of La<sub>0.5</sub>Ca<sub>0.5</sub>MnO<sub>3</sub>

## Table ESI 2

	Pnma				P2 <sub>1</sub> /m				Refinement
T (K)	a	b	c	b/√2	a	b	c	β	agreement factors
298	5.425(1)	7.642(1)	5.4189(9)	5.404					$\begin{array}{c} R_{p}: 5.48; \\ R_{wp}: 7.40; \\ R_{exp}: 6.46 \\ \chi^{2}: 1.31 \end{array}$
270	5.422(1)	7.6408(9)	5.4192(9)	5.403					$\begin{array}{c} R_{p}: 6.06; \\ R_{wp}: 8.52; \\ R_{exp}: 6.02 \\ \chi^{2}: 2.01 \end{array}$
250	5.4215(8)	7.632(1)	5.4131(9)	5.397					$\begin{array}{c} R_{p}: 5.96; \\ R_{wp}: 8.12; \\ R_{exp}: 6.10 \\ \chi^{2}: 1.77 \end{array}$
230	5.420(1)	7.6233(9)	5.411(1)	5.390					$R_{p}: 5.80;R_{wp}:8.02;R_{exp}:6.62\chi^{2}: 1.89$
210	5.421(1)	7.6188(9)	5.414(1)	5.387					$\begin{array}{c} R_{p}: 5.40; \\ R_{wp}: 7.47; \\ R_{exp}: 6.14 \\ \chi^{2}: 1.48 \end{array}$
190	5.420(1)	7.618(1)	5.4147(9)	5.387					$\begin{array}{c} R_{p}: \ 7.13; \\ R_{wp}: 9.68; \\ R_{exp}: 6.66 \\ \chi^{2}: \ 2.11 \end{array}$
170	5.418(1)	7.592(1)	5.400(1)	5.369					$\begin{array}{c} R_{p}: 5.98; \\ R_{wp}: 8.19; \\ R_{exp}: 6.17 \\ \chi^{2}: 1.76 \end{array}$
150	5.4179(9)	7.596(2)	5.401(1)	5.371	5.4546(3)	10.8651(8)	7.5292(6)	90.097(2)	$\begin{array}{c} R_{p}: 4.64; \\ R_{wp}: 6.27; \\ R_{exp}: 6.18 \\ \chi^{2}: 1.03 \end{array}$
130	5.416(1)	7.609(2)	5.402(1)	5.381	5.462(1)	10.8730(9)	7.5209(6)	90.039(2)	$\begin{array}{c} R_{p}: 4.69; \\ R_{wp}: 6.40; \\ R_{exp}: 6.17 \\ \chi^{2}: 1.08 \end{array}$
110	5.416(1)	7.617(2)	5.401(1)	5.387	5.4635(8)	10.8749(7)	7.5191(6)	90.042(1)	$\begin{array}{c} R_{p}: 4.73; \\ R_{wp}: 6.40; \\ R_{exp}: 6.18 \\ \chi^{2}: 1.07 \end{array}$
90	5.416(1)	7.6191(8)	5.401(9)	5.387	5.4641(3)	10.8751(4)	7.5189(4)	90.038(4)	$ \begin{array}{c} R_{p}: 4.75; \\ R_{wp}: 6.43; \\ R_{exp}: 6.13 \\ \chi^{2}: 1.10 \end{array} $
70	5.4158(9)	7.619(2)	5.4002(1)	5.387	5.4636(4)	10.8735(9)	7.5181(4)	90.038(5)	$\begin{array}{c} R_{p}: 4.84; \\ R_{wp}: 6.60; \\ R_{exp}: 6.12 \\ \chi^{2}: 1.16 \end{array}$
15	5.418(1)	7.623(2)	5.394(1)	5.391	5.4614(3)	10.8701(6)	7.5163(4)	90.100(5)	$R_{p}: 4.98;R_{wp}:6.63;R_{exp}:6.18\chi^{2}: 1.15$

**Figures ESI 2** X-ray diffraction patterns for  $La_{0.5}Ca_{0.5}MnO_3$  powder at different temperatures. Observed intensities are printed in red dots, Rietveld refined curves are printed as full black lines, and the difference between the two by blue lines underneath. The expected Bragg reflections are marked by blue and red dashes for *Pnma* and *P2*<sub>1</sub>/*m* phases respectively.





















#### ESI 4. Basic principles of impedance spectroscopy

Impedance spectroscopy is a well-established tool for the separate analysis of several dielectric contributions to the macroscopic dielectric properties of solid matter [1-3]. Impedance spectroscopy experiments consist in the application of an electric stimulus in terms of a time (*t*)-dependent alternating voltage signal U of variable angular frequency  $\omega(\omega = 2\pi f)$  and fixed amplitude  $U_0$  applied to the sample under test:  $U(\omega, t) = U_0 \sin(\omega t)$  (see Figure ESI 3, blue arrow).



**Figure ESI 3** Phasor diagram for an applied voltage signal *U* (blue arrow) and the current response (red arrows) of an (1) ideal resistor ( $I_R$ ), (2) ideal capacitor ( $I_C$ ), (3) ideal RC element ( $I_{RC}$ ), (4) non-ideal constant phase element ( $I_{CPE}$ ), and (5) ideal inductor ( $I_L$ ). The complex impedance for the equivalent circuit components R, C, RC, CPE and L are given. The current response of different components is characterised by the amplitude (length of the red arrows) and a characteristic phase shift.

Effectively, the impedance spectroscopy apparatus measures the amplitude  $I_0$  and phase shift  $\delta$  of the alternating current response signal *I* (Figure ESI 3, red arrows for different possible current response signals) over a wide frequency range:  $I(\omega, t) = I_0 \sin(\omega t + \delta)$ . One period of the applied voltage signal corresponds to a  $2\pi$  (360°) rotation of the  $U(\omega,t)$  arrow on the phasor diagram shown in Figure ESI 3. For the most common equivalent circuit elements the current response is (1) in phase ( $\delta = 0$ ) with the applied voltage for an ideal resistor,  $I_R$ , (2) out of phase by  $\delta = -\pi/2$  for an ideal capacitor,  $I_C$ , (3) out of phase by an *f*-dependent phase shift  $0 < \delta < -\pi/2$  for an ideal parallel resistor-capacitor (RC) element, (4) out of phase by  $\delta = -\pi/2 + \gamma$  for a constant-phase element (CPE) that can be regarded a "non-ideal" capacitor, and (5) out of phase by  $\delta = +\pi/2$  for an ideal inductor,  $I_L$ . Note that CPEs displays an *f*-independent phase shift  $\gamma$ , thus the name constant phase element, where the angle  $\gamma$  is accounted for with respect to an ideal capacitor. The angle  $\gamma$  is correlated to the critical CPE exponent *n* (see Figure ESI 3), which in turn is correlated with the degree of suppression of the typical semicircle arc below the real x-axis in complex impedance plane plots of *Z*" vs *Z*" (Figures 5a and 6a in the main text).

All phase angles are *t*- independent, which implies that all arrows in Figure ESI 3 would rotate at the same  $\omega$  with *t*-independent phase angles. This illustrates that the impedance can be defined as a *t*-independent complex number,  $Z^* = Z' + iZ''$ . The impedance fraction in phase with  $U(\omega,t)$  is defined as the real, and the part out of phase by +/-  $\pi/2$  as the positive/negative imaginary part. By definition, a negative phase shift corresponds to a current preceding the applied voltage *U* as is the case for an ideal capacitor ( $I_C$ ), which leads to a negative imaginary part of impedance *Z*'', whereas a positive phase shift corresponds to a current lagging behind *U* as is the case for an ideal inductor ( $I_L$ ) and leads to positive *Z*''.

In practice, impedance spectra of insulating or semiconducting samples show predominantly negative Z'', but a crossover from negative (capacitive) to positive (inductive) Z'' at high  $\omega$  can sometimes

occur in more conducting samples. Note that the impedance for an inductor  $(Z*_L)$  increases linearly with  $\omega$  or f, whereas the impedance of a capacitor  $(Z*_C)$  decreases with the proportionality  $\sim 1/\omega$  or  $\sim$ 1/f (see Figure ESI 3). This implies that inductive contributions would appear in the impedance spectra at sufficiently high-f. In the case that an inductive contribution may be dominant over a certain high-f regime, it is not possible to extract the capacitance of the relevant dielectric contribution. Z" is positive and the capacitance is thus not accessible. The opposite applies in capacitive lower f regimes, where only the capacitance of the respective dielectric contribution may be determined but not the inductance. The equivalent circuits applied to impedance spectroscopy data usually reflect that fact and are, thus made up of conventional RC elements to represent the f-ranges where capacitive contributions dominate, whereas for inductive dominated high-f ranges, they may contain inductive L-elements.

For the more common capacitive impedance response (Z'' < 0), each dielectric contribution in an insulating or semiconducting sample can usually be represented by an RC element. One RC element is visible in Z'' vs Z' complex impedance plane plots in form of one impedance semicircle arc. The arc diameter corresponds to R and the frequency where the arc has its maximum is given by  $\omega_{max} = 1/RC = \tau^{-1}$ . Each dielectric contribution (RC element) in the impedance response may be associated with a certain area in the sample like grain interior bulk or GBs, and the overall equivalent circuit model is then made up of a series connection of the various RC elements.

The RC element model can be understood intuitively. The application of a voltage stimulus to a sample leads to two instant and concurrent responses from that sample, i.e. dielectric polarization of localized charges and the concomitant displacement currents, as represented by the capacitor, and charge transport of mobile charge carriers and concomitant conduction currents, as represented by the resistor. Considering the example of an RC element representing an insulating material, the capacitor describes the polarizability of the material and its ability to store charge, and the parallel

resistor describes the sample's leakage resistance that arises from the charge transport of mobile charge carriers that bypass or "escape" the ideal charge storage element.

Almost all dielectric contributions encountered in experimental impedance spectra are usually not ideal as represented by ideal RC circuit elements, in which case ideal capacitors are commonly replaced by non-ideal CPEs. The true capacitance C of a dielectric contribution, as represented by a possible R-CPE element, and the critical CPE exponent *n* can be extracted using equivalent circuit fitting software [4], where *n* = 1 for an ideal capacitor, *n* < 1 measures the non-ideality of the dielectric contribution and *n* > 1 is unphysical. As mentioned above, the critical exponent *n* from the CPE is correlated with the degree of suppression of the semicircle arc in plots of *Z*'' vs *Z*', which indicates the width of the distribution of dielectric relaxation times  $\tau$ . The CPE is usually interpreted as a sign for a certain distribution of  $\tau$  across the respective sample area, where  $\tau = R \ge \rho \varepsilon_0 \varepsilon_r'$  in an ideal RC element and decreasing *n* indicates an increasing width of the distribution of  $\tau$ . The exact shape of the distribution function of  $\tau$  is not accessible from impedance spectroscopy data, since the CPE only yields one more additional parameter *n*, which is related to  $\gamma$ , and any attempt to determine the precise distribution function from impedance spectroscopy data should be regarded unphysical [5].

#### References

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[4] C. H. Hsu and F. Mansfeld, Corrosion **57**, 747 (2001).

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### ESI 4. Equivalent circuit fits

A data file (ESI 4. Equivalent circuit fits.xlsx) is attached to this publication, which contains all fitted parameters as extracted from the fitting software. The resistance and capacitance values are not normalized, i.e. given in  $\Omega$  and Farad.