**ARTICLE TYPE** 

## **Appendix 2**

## 2.1 Symbolic dynamics and analysis blocks

Complex nonlinear dynamical systems can be modelled and studied in the context of coarse-graining, *i.e.*, they can viewed as information capacitors containing discrete series of symbolic messages.<sup>1-5</sup> Coarse-graining can be outlined through symbolisation, a process which allows partitioning of the full continuous phase space into finite number of cells.<sup>1</sup> In this manner, symbolisation provides a rigorous way of studying actual complex dynamics under finite precision. This operational procedure is often referred as symbolic dynamics.<sup>1,4–7</sup> One of the merits of symbolic dynamics is that it provides a strong link between dynamical systems and information theory.<sup>1,4,5,8</sup> According to symbolic dynamics, time-series are re-organised into new symbolic sequences, in which every symbol stands for a partition of the initial time-series. Every different symbol is represented by an alphabet letter and the whole procedure is called lettering.<sup>1–3,6,7,9</sup> The new subsequences of symbols are called words.<sup>2,3,7</sup> Reading of symbolic sub-sequences can be derived through the processes of lumping or gliding.<sup>2,3,7</sup> Lumping is the interpretation of symbolic words through independent sequential discrete portions of certain number of words, called blocks, opposed to gliding, where the portions are not independent.<sup>2,3,7</sup> Note that gliding is the standard convention in literature and is often referred also under the term sliding or moving-frame.<sup>2,3,7</sup> Block timeseries symbolisation prerequisites selection of  $\lambda$  different letters from an alphabet and choice of the size, n, of blocks or words, *i.e.*, the number of sequential letters that will be treated as a whole. Depending on  $\lambda$  and *n* the maximum number, N of different words is determined in the selected alphabet. For example in a  $\lambda = 2$  lettering, a threshold C may be considered. Each value above this threshold may be symbolised as 1 and each below, as 0.<sup>2,3,7</sup> Initial timeseries, for instance, of length L = 20, may be transformed through symbolic dynamics to e.g. 11001010111000101010 in the  $\lambda = 2$  lettering. Through lumping, the  $\lambda = 2$ -letter symbolic time-series may be organised in sets of n = 2blocks as (11|00|10|10|11|10| 00|10|10|10|), in which each block is one of the  $N = 2^2 = 4$  different words in this lettering, *i.e.*, (00, 01, 10,11). The same symbolic sequence through gliding-sliding will be treated as (11|10|00|01|10|01| 10|01|11|11|10|00| 00|01|10|01|10|01| 10|01|10), *i.e.*, the sequence will be of greater length. The  $\lambda = 2$  letter sequence may be also organised in blocks of n = 3 letters with maximum of  $N = N_{3,2} = n^{\lambda} = 3^2 = 9$  different words,*i.e.*, (000, 001, 010, 100, 110, 011, 010, 001, 111). Other sequences of words may be generated as well. In general, 1-3,6,7,9through symbolic dynamics a L-length time-series is transformed to a symbolic time-series sequence,  $[A_1, A_2...A_n...A_L]$ , composed by  $\lambda$  different letters,  $[A^1, A^2 \dots A^{\lambda}]$ , from a  $\lambda$ -length alphabet. Symbolic time-series sequences are re-organised in *n*-sized words-blocks composed by letters of the alphabet  $[A^1, A^2 \dots A^{\lambda}]$ . In linguistics the word size is unconfined and, hence, linguistic words contain some or, potentially, all letters. On the other hand, in symbolic dynamics, the words are of fixed lengths  $n, n \geq \lambda$  and are chosen from  $N_{max} = N_{n,\lambda} = n^{\lambda}$  different fixed-sized words in the  $[A^1, A^2 \dots A^{\lambda}]$  alphabet. In this manner, the symbolic time-series are reorganised as  $\dots A_1 \dots A_n A_{n+1} \dots A_{2n} \dots A_{in+1} \dots A_{(i+1)n} \dots$  blocks,

where *i* is the consecutive number of the block, *i.e.*, i = 1...Total number of blocks. In lumping the *n* words-blocks are sequentially independent, on the contrary to the dependent sequential blocks in the gliding-sliding process. The total number of blocks of the symbolic time-series is greater for gliding-sliding and hence more computation is needed. The probability of occurrence of a block,  $[A_1, A_2...A_n]$ , of size *n* is calculated by

$$p^{(n)}(A_1, A_2...A_n) = \frac{Number of occurrences of block [A_1, A_2...A_n]}{Total number of blocks}$$
(2.1)

## 2.2 Block entropy analysis

In the framework of complex signal analysis, specific entropy methodologies based on symbolic dynamics have been developed in the previous decade.<sup>5,10–13</sup> All these methodologies are referred as block entropies. Most common techniques rely on the extension of Shannon entropy<sup>14</sup>

$$H_S = -\sum p_i ln p_i \tag{2.2}$$

where  $p_i$  is the number of possible microscopic configurations. Note that equation (2.2) represents the classical Boltzman's entropy for the Gibbs canonical ensemble (B-Gentropy).<sup>9</sup> Combining equations (2.1) and (2.2), the Shannon block entropy, H(n) of *n*-sized blocks is derived by (2.3):

$$H(n) = -\sum_{(A_1, A_2, \dots, A_n)} p^{(n)}(A_1, A_2, \dots, A_n) ln p^{(n)}(A_1, A_2, \dots, A_n)$$
(2.3)

Equation (2.3) calculates the entropy due to all possible words. It is a measure of uncertainty or disorder, *i.e.*, it measures organisation deficiency of a complex system. It also gives the average amount of information necessary to predict a sub-sequence of words or blocks of length n

From equation (2.3), the Shannon block entropy per letter may be derived by:

$$h^{(n)} = \frac{H(n)}{n} = \frac{-\sum_{(A_1, A_2, \dots, A_n)} p^{(n)}(A_1, A_2, \dots, A_n) ln p^{(n)}(A_1, A_2, \dots, A_n)}{n}$$
(2.4)

This journal is © The Royal Society of Chemistry [year]

This entropy may be interpreted as the average uncertainty of a block of size n per letter.<sup>6</sup>

From the Shannon block entropy the conditional entropy may be derived by equation (2.5):

$$h_{(n)} = H(n+1) - H(n)$$
(2.5)

The conditional entropy  $h_{(n)}$  measures the uncertainty of predicting a state one step into the future, provided a history of the preceding n states.<sup>6</sup>

For physical phenomena with long-range interactions or long-range memory effects, an important property observed is the violation of Boltzmann-Gibbs (B-G) statistics.<sup>15</sup> A generalised expression of the B-G statistics has been proposed based on multifractal concepts by Tsallis<sup>16,17</sup>

$$S_q = \frac{1}{q-1} \left( 1 - \sum_{i=1}^{W} p_i^q \right)$$
(2.6)

where  $p_i$  denotes, in these references, the probabilities of a sequence and W their total number. q is a real nuber which is the measure of non-extensivity of the system.<sup>16,17</sup> Using  $p_i^{(q-1)} = e^{(q-1)\ln(p_i)} \sim 1 + (q-1)\ln(p_i)$  in the limit  $q \rightarrow 1$  the B-G entropy is derived.<sup>6,9</sup> The generalization of B-G expression, suggests the non-extensive statistical mechanics. The entropic index q characterises the degree of nonadditivity in the following pseudo-additivity rule<sup>2,3,6,9,16</sup>

$$S_q(A,B) = S_q(A) + S_q(B) + (1-q)S_q(A)S_q(B)$$
(2.7)

with q > 1 referring to sub-additivity and q < 1 to superadditivity. Systems that called non-extensive, have special probability correlations and extensivity may occur for  $S_q$  for specific value of index q.<sup>18</sup> Tsallis entropy has been used in terms of symbolic dynamics for electromagnetic time series prior to earthquakes.<sup>2,3,6,9,16,19</sup> By properly adjusting symbolisation in accordance to (2.3) and (2.4), the Tsallis entropy of a block  $[A_1, A_2...A_n]$  of length n in a  $\lambda$ -letter alphabet can be calculated by <sup>2,3,6,9,19</sup>

$$S_q(n) = \frac{1}{q-1} \left(1 - \sum_{(A_1, A_2, \dots, A_n)} \left[p^{(n)}(A_1, A_2 \dots A_n)\right]^q\right)$$
(2.8)

where  $p^{(n)}(A_1, A_2...A_n)$  is the probability of occurrence of block  $[A_1, A_2...A_n]$ . As already mentioned, high level of organisation is indicated when low values of Tsallis entropy are produced. Tsallis entropy has been explored in the field of earthquake time series analysis.<sup>6,9,15,20–22</sup> Recent work has also been published in biomedical imaging, with suggestions in replacing Shannon's theorem<sup>23</sup> and bioinformatics.<sup>23</sup> An interconnection between fractals and Tsallis entropy that has been introduced in previous decade could provide natural frame for studying fractally structured systems.<sup>24</sup> Moreover, a possible interconnection could exist between generalized Tsallis statistics and quantum groups.<sup>17</sup>

From (2.8) the normalised Tsallis entropy may be derived. The formula for the computation is  $^{16}$ 

$$\hat{\mathbf{S}}_{q}(n) = \frac{\frac{1}{q-1} (1 - \sum_{(A_{1}, A_{2}, \dots, A_{n})} [p^{(n)}(A_{1}, A_{2} \dots, A_{n})]^{q})}{\sum_{(A_{1}, A_{2}, \dots, A_{n})} [p^{(n)}(A_{1}, A_{2} \dots, A_{n})]^{q}}$$
(2.9)

where the symbolisation was adjusted in accordance to equations (2.3) and (2.4) following the approach proposed by other investigators. <sup>6,9,19,20</sup> In equation (2.9),  $p^{(n)}(A_1,A_2...A_n)$  is the probability of occurrence of the block  $[A_1,A_2...A_n]$  and q is the corresponding real number of equations (2.6),(2.8) and (2.9).

The appropriate choice of the entropic index q has crucial meaning for the Tsallis and the normalised Tsallis entropy computation and requires further exploration for its proper use.<sup>9</sup> For every specific use of Tsallis entropy the ranges of the q values will result in significant discrimination.<sup>17</sup> Nonadditive Tsallis entropy combined with Gutenberg-Richter law provided excellent fit to seismicities with q-values range from 1.4 to 1.85. The q-values are rooted in a rather solid physical background and describe the non-additivity of a seismic emmision in a correct manner.<sup>9</sup> Moreover, index q can be consider as bias parameter with q < 1 refer to rare events and q > 1 refer to prominent events.<sup>22</sup> For pre-eqrthquake electromagnetic disturbances, the q-values are restricted in the region 1 < q < 2, and are consistent with several studies that suggest the upper limit to be equal to  $2.^{9}$  It is noteworthy that entropic index q is not a measure of complexity but measures the nonextensivity of the system.<sup>6</sup>

## References

- K. Karamanos and G. Nicolis, *Chaos Soliton. Frac.*, 1999, 10, 1135–1150.
- 2 K. Karamanos, D. Dakopoulos, K. Aloupis, A. Peratzakis, L. Athanasopoulou, S. Nikolopoulos, P. K. P and K. Eftaxias, *Phys. Rev. E.*, 2006, **74**, 21–36.
- 3 K. Karamanos, J Phys. A Math. Gen., 2001, 34, 9231– 9241.
- 4 J. Nicolis, *Chaos and Information Processing*, World Scientific, Singapore, 1991.
- 5 G. Nicolis and P. Gaspard, *Chaos Soliton. Fract.*, 1994, 4, 41–57.
- K. Eftaxias, G. Balasis, Y. Contoyiannis, C. Papadimitriou, M. Kalimeri, L. Athanasopoulou, S. Nikolopoulos, J. Kopanas, G. Antonopoulos and C. Nomicos, *Nat. Hazard Earth Sys.*, 2009, 9, 1953–1971.

- 7 K. Karamanos, D. Dakopoulos, K. Aloupis, A. Peratzakis, L. Athanasopoulou, S. Nikolopoulos, P. K. P and K. Eftaxias, *Phys. Rev. E.*, 2006, **74**, 21–36.
- 8 A. Voss, J. Kurths, H. Kleiner, A. Witt, N. Wessel, N. S. ad K. Osterziel, R. Schurath and R. Dietz, *Cardiovasc. Res.*, 1996, **31**, 419–433.
- 9 M. Kalimeri, C. Papadimitriou, G. Balasis and K. Eftaxias, *Physica A*, 2008, **387**, 1161–1172.
- 10 W. Ebeling, Physica D, 1997, 109, 42-52.
- 11 W. Ebeling and G. Nicolis, *Chaos Soliton Fract.*, 1992, **2**, 635–650.
- 12 W. Ebeling and G. Nicolis, *Europhys Lett.*, 1991, **14**, 191–196.
- 13 K. Karamanos, Lect. Notes Phys., 2000, 550, 357-371.
- 14 C. Shannon, Tech. J., 1948, 27, 379-423.
- 15 Y. Contoyiannis and K. Eftaxias, Nonlin. Processes Geophys., 2008, 15, 379–388.
- 16 H. Suyari, 2002.
- 17 C. Tsallis, J. Stat. Phys., 1988, 52, 479-487.
- 18 J. Boon and K. Tsallis, *Europhys. News*, 2005, **36**, 185– 186.
- 19 K. Eftaxias, G. Balasis, Y. Contoyiannis, C. Papadimitriou, M. Kalimeri, L. Athanasopoulou, S. Nikolopoulos, J. Kopanas, G. Antonopoulos and C. Nomicos, *Nat. Hazard Earth Sys.*, 2010, **10**, 275–294.
- 20 K. Eftaxias, Physica A, 2010, 389, 133-140.
- 21 V. Surkov, S. Uyeda, H. Tanaka and M. Hayakawa, J. Geodyn., 2002, 33, 477–487.
- 22 L. Zunino, D. Perez, A. Kowalski, M. Martin, M. Garavaglia, A. Plastino and O. Rosso, *Physica A*, 2008, **387**, 60057–6068.
- 23 J. Mohanalin, S. Beenamol, P. Kalra and N. Kumar, *Comput. Math. Appl.*, 2010, **60**, 2426–2432.
- 24 F. Lopes, E. de Oliveira and R. Cesar, *BMC Syst. Biol.*, 2011, **5**, 61–69.