

# Mechanics of Twisted Hippuric Acid Crystals Untwisting as they Grow

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## Electronic Supplementary Information

**Fig. S1** Relationship between twist period ( $2P$ , mm) and the smallest size in a cross section ( $h$ ,  $\mu\text{m}$ ) fitted with by the power law,  $P = kh^n$ , and the linear function,  $P = k_0 + k_1h$ .

**Fig. S2** Correlation between twist intensity ( $r/P$ ) and temperature at the bench,  $T_{\text{bench}}$  (a) and the substrate,  $T$  (b).

**Supplementary information text.** Reformulating of discrete untwisting models in terms of continuous growth.

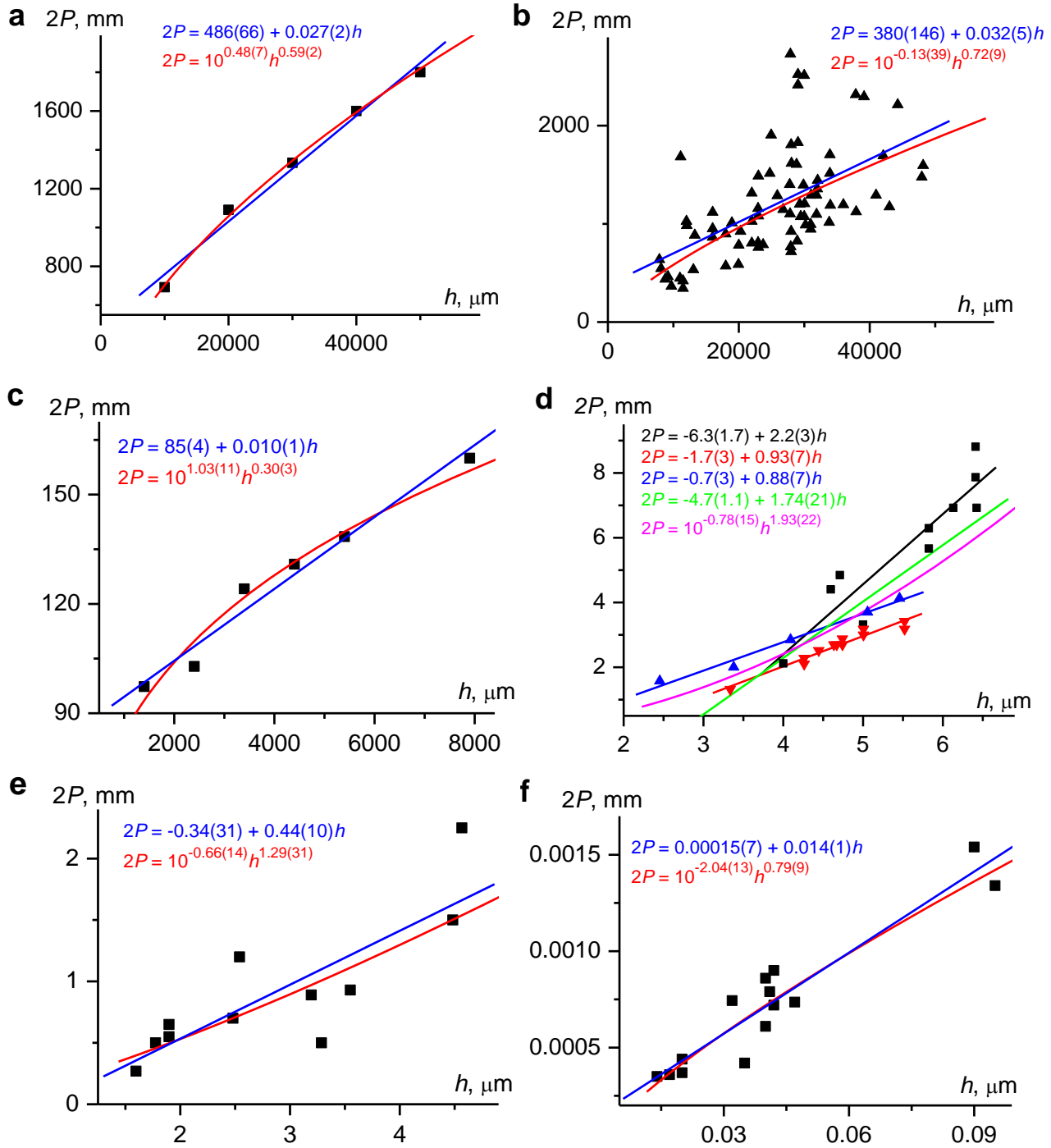


Fig. S1 Relationship between twist period ( $2P$ , mm) and the smallest size in a cross section ( $h$ ,  $\mu\text{m}$ ) fitted with by the power law,  $P = kh^n$ , and the linear function,  $P = k_0 + k_1h$ . (a) Natural quartz, data from ref. 7; (b) Natural quartz, data from ref. 6; (c) Oxalic acid dihydrate, data from ref. 7; (d) Melt grown hippuric acid; separate fittings for three crystals and simultaneous fitting, data from ref. 12; (e) (1,4-bis[2-(pyrene-1-yl) vinyl]-2,5-dimethylbenzene)IBr<sub>2</sub>, data from ref. 11; (f) Poly(*m*-phenylene isophthalamide), data from ref. 14.

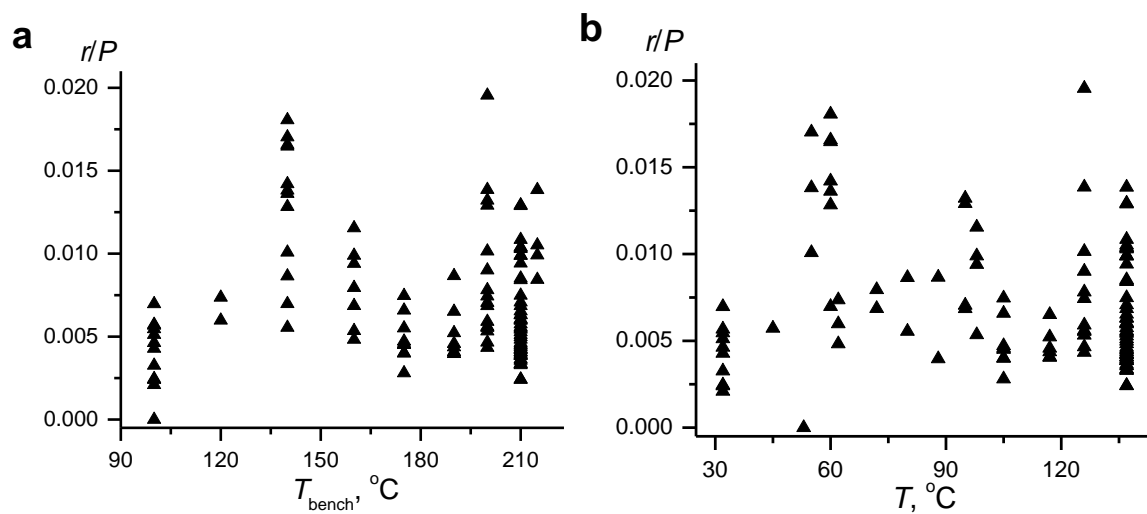


Fig. S2 Correlation between twist intensity ( $r/P$ ) and temperature at the bench,  $T_{\text{bench}}$  (a) and the substrate,  $T$  (b). In order to minimize the effect of crystal size the twist intensity,  $1/P$ , is multiplied by the crystal radius,  $r$ .

## Supplementary information text

### Reformulating of discrete untwisting models in terms of continuous growth

It should not be a big error to replace growth by discrete layers of thickness  $\Delta r = 0.7$  nm with continuous growth by infinitesimal layers, since experimentally observed  $\Delta r$  is much less than the minimal observed crystal size  $r_0 > 50$  nm. The equation (8) can be rewritten as

$$\theta_{T(N)} r_N^4 = \theta_{p,0} r_0^4 + \sum_{i=1}^N \theta_{p,i} (r_i^4 - r_{i-1}^4), \quad (\text{S1})$$

where summation can be replaced by integration

$$\theta_{T(N)} r_N^4 = \theta_{p,0} r_0^4 + \int_{r_0}^{r_N} \theta_p(r) dr^4 = \theta_{p,0} r_0^4 + 4 \int_{r_0}^{r_N} \theta_p(r) r^3 dr. \quad (\text{S2})$$

The condition (7) can be rewritten as  $\gamma_{p,i} = \theta_{p,i} r_i = \theta_{p,i-1} r_{i-1} \alpha = \gamma_{p,i-1} \alpha$ . Then the difference in shear strain  $\Delta \gamma_{p,i} = \gamma_{p,i} - \gamma_{p,i-1} = \gamma_{p,i-1} \alpha$  in a continuous form is

$$\frac{d\gamma}{dr} \Delta r = \gamma(\alpha - 1) \quad (\text{S3})$$

Solution to eq. (S3) gives  $\gamma_p = \gamma_{p,0} \exp\left([\alpha - 1] \frac{r - r_0}{\Delta r}\right)$  or in terms of twist intensity

$$\theta_p = \theta_{p,0} \frac{r_0}{r} \exp\left([\alpha - 1] \frac{r - r_0}{\Delta r}\right). \quad (\text{S4})$$

The final expression is obtained by combining eqs. (S2) and (S4)

$$\frac{\theta_{T(N)} r_N^4}{\theta_{p,0} r_0^4} = 1 + \frac{4}{r_0^3} \int_{r_0}^{r_N} r^2 \exp\left([\alpha - 1] \frac{r - r_0}{\Delta r}\right) dr. \quad (\text{S5})$$

This integral can be solved straightforwardly to give

$$\frac{P_0 r_N^4}{P_N r_0^4} = 1 + \frac{4}{r_0^3} \left[ \left( \frac{r_N^2}{k} - \frac{2r_N}{k^2} + \frac{2}{k^3} \right) \exp(k[r_N - r_0]) - \left( \frac{r_0^2}{k} - \frac{2r_0}{k^2} + \frac{2}{k^3} \right) \right], \quad (\text{S6})$$

where  $k = (\alpha - 1)/\Delta r$ . Eq. (9) cannot be directly obtained from eq. (S6) because of the zero denominator,  $k = 0$ , but it can be obtained easily by solving eq. (S5) with  $\alpha = 1$ .