

Supporting Information

Molecular dynamics simulation of the conductivity mechanism of nanorod filled polymer nanocomposites

Yangyang Gao¹, Dapeng Cao⁴, Jun Liu^{1*}, Jianxiang Shen¹, Youping Wu¹, Liquan Zhang^{1, 2, 3, 4*}

¹Key Laboratory of Beijing City on Preparation and Processing of Novel Polymer Materials,
People's Republic of China

²Beijing Engineering Research Center of Advanced Elastomers, People's Republic of China

³Engineering Research Center of Elastomer Materials on Energy Conservation and Resources,
Ministry of Education, PRC

⁴State Key Laboratory of Organic-Inorganic Composites, Beijing University of Chemical
Technology, 100029 Beijing, People's Republic of China

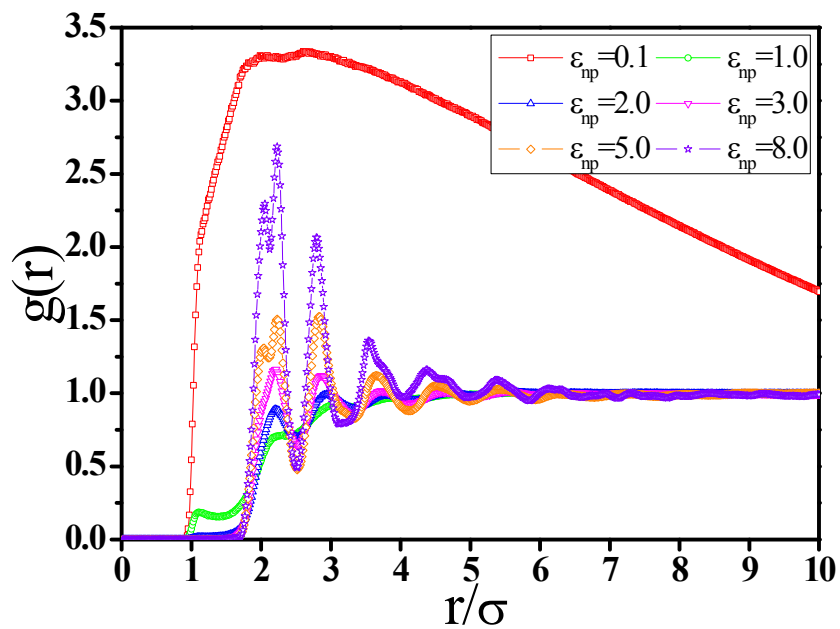
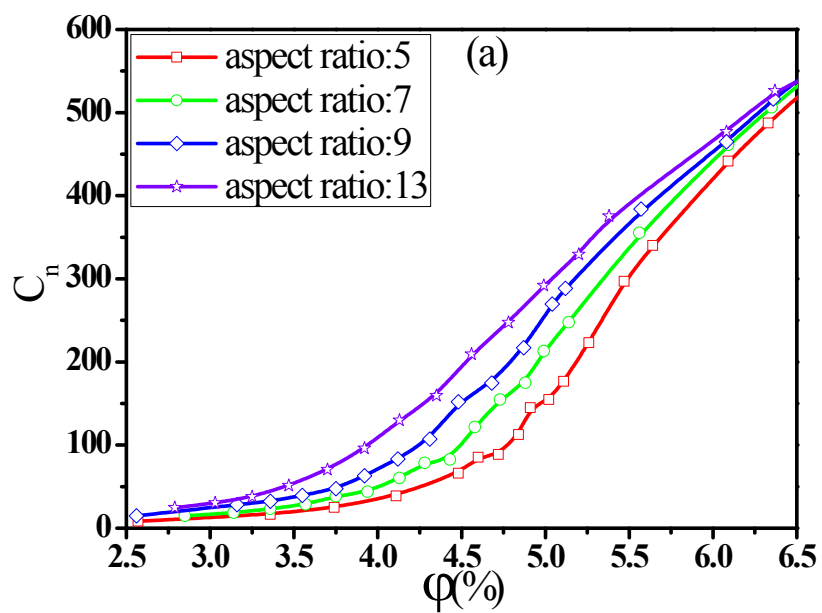


Fig. S1 RDF of nanorods for different polymer-nanorod interactions ε_{np} . ($T=1.0$, $\varphi = 4.49\%$)



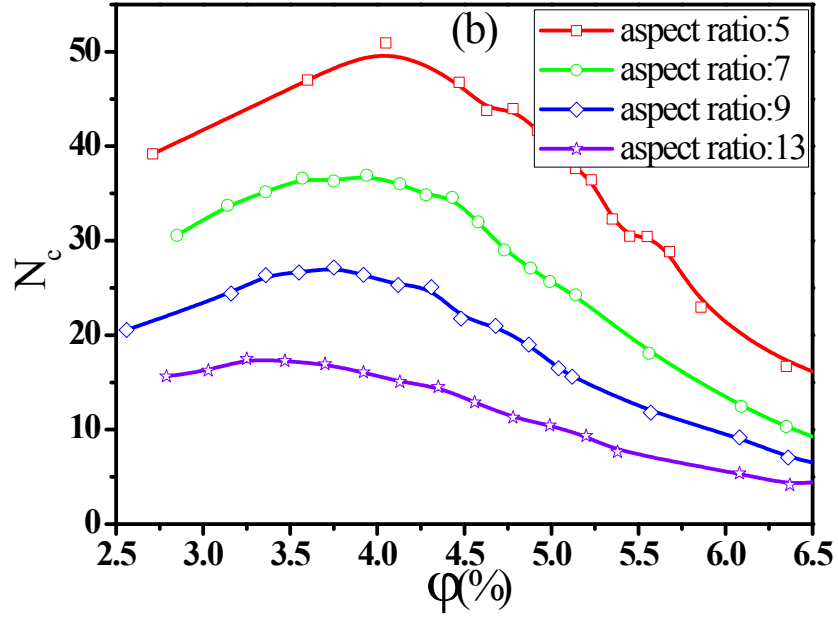
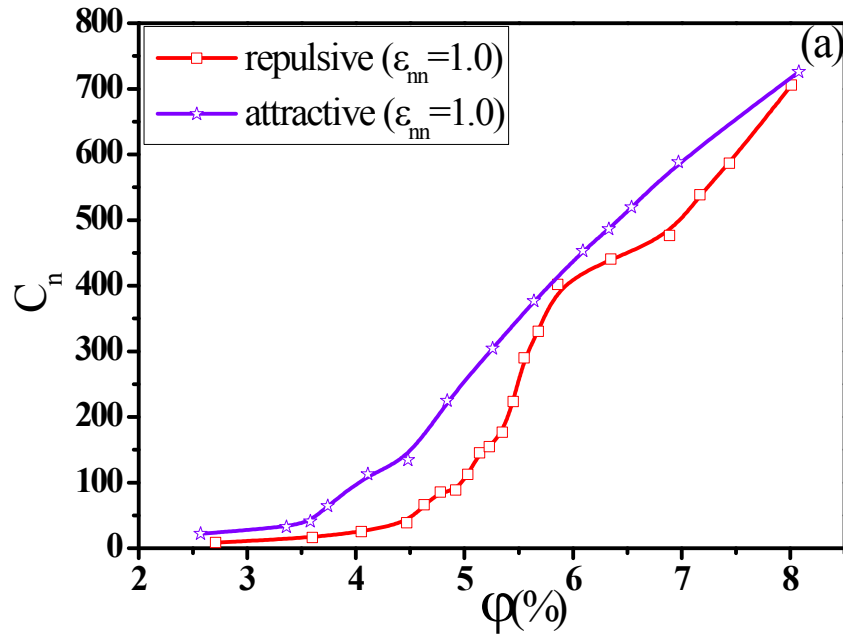


Fig. S2(a) The maximum cluster size C_n , (b) The total number of clusters N_c as a function of the volume fraction ϕ of nanorods for different aspect ratios. ($T=1.0$, $\varepsilon_{np}=1.0$)



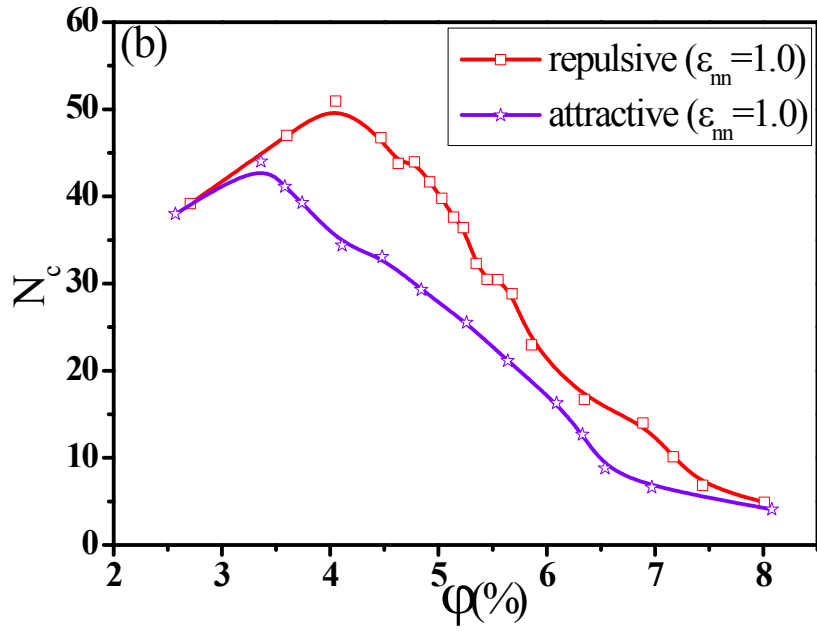


Fig. S3(a) The maximum cluster size C_n , (b) The total number of clusters N_c as a function of the volume fraction ϕ of nanorods for repulsive and attractive filler-filler interaction ε_{nn} . ($T=1.0$, $\varepsilon_{np}=1.0$)

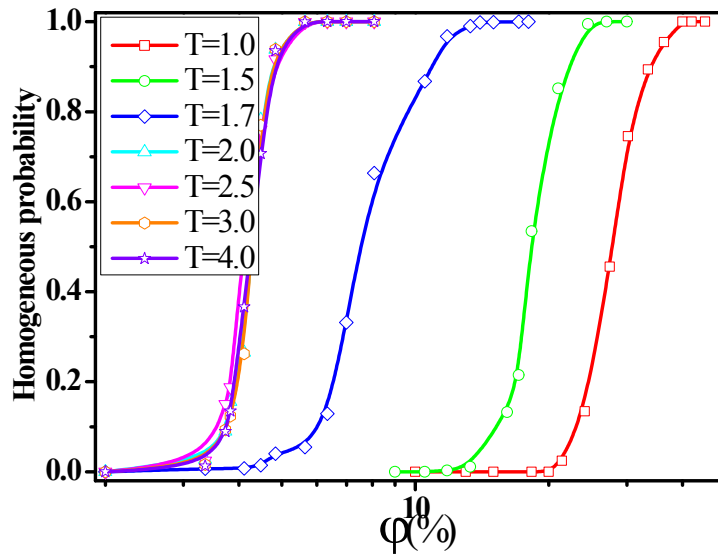


Fig. S4 Homogeneous probability as a function of the volume fraction ϕ of nanorods for different temperatures T . ($\varepsilon_{np}=0.1$)

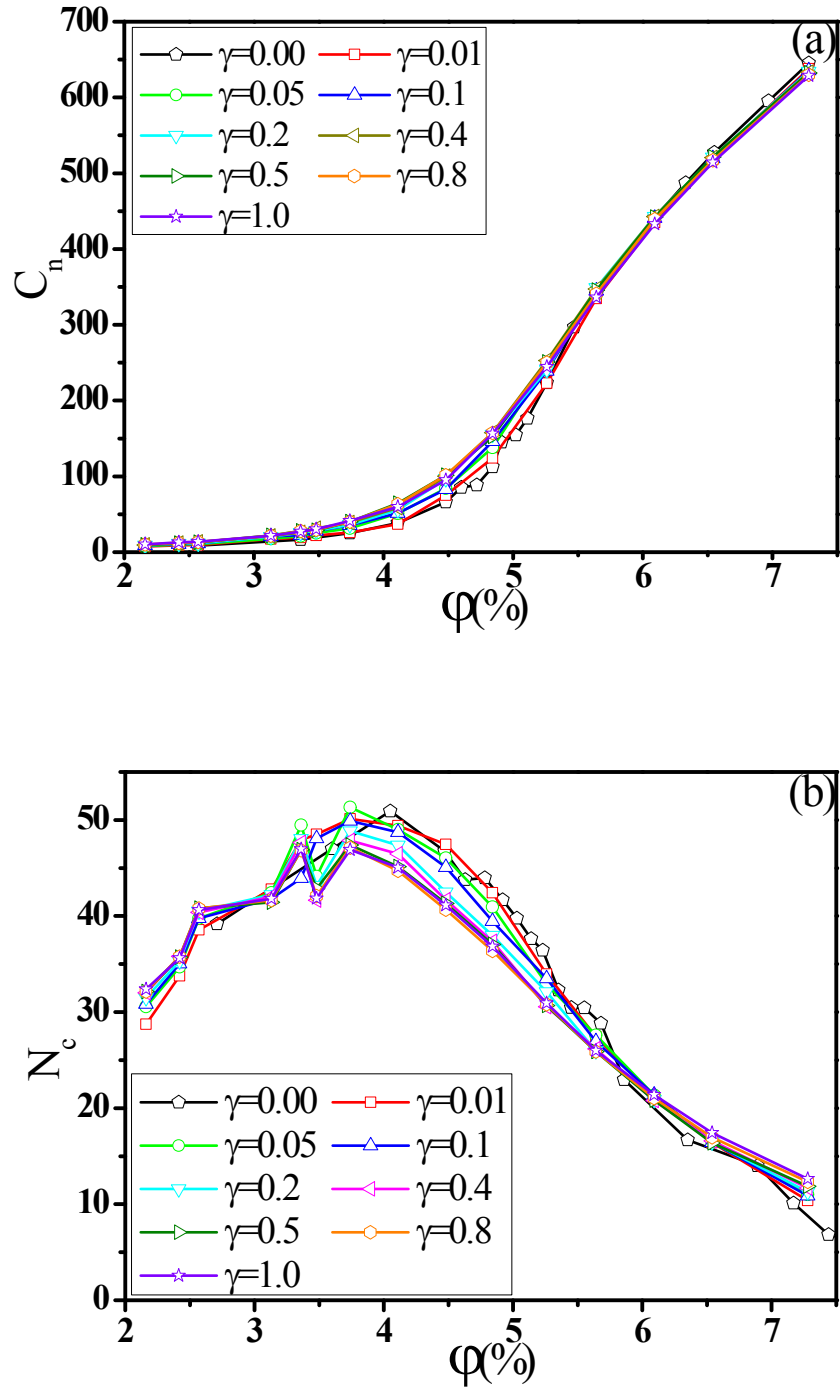


Fig. S5(a) The maximum cluster size C_n , (b) The total number of clusters N_c as a function of the volume fraction ϕ of nanorods for different shear rates γ . ($T=1.0$, $\varepsilon_{np}=1.0$)

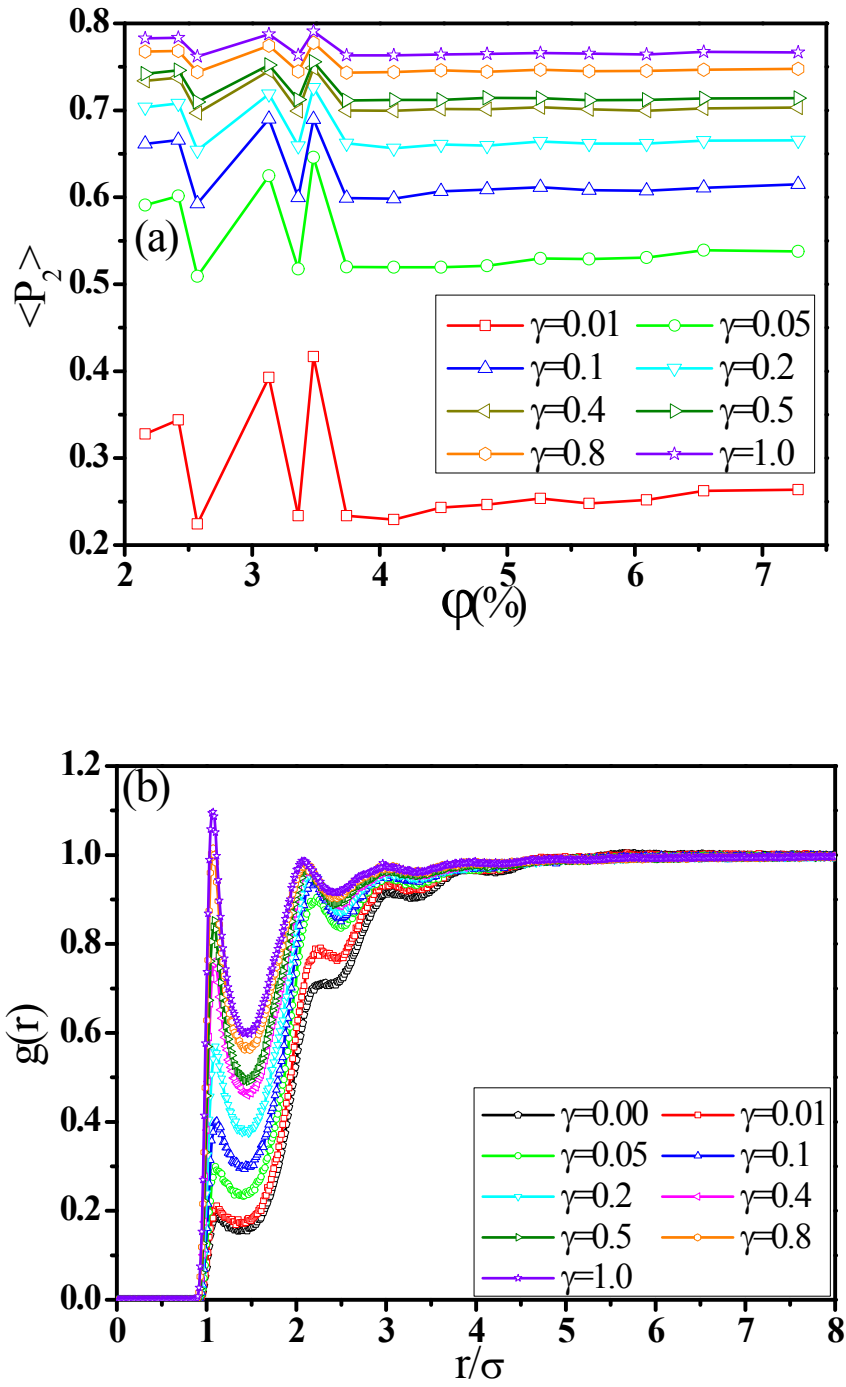


Fig. S6(a) The second-order Legendre polynomials $\langle P_2(\cos \theta(t)) \rangle$ as a function of the volume fraction ϕ of nanorod. (b) RDF of nanorods for different shear rates γ . ($T=1.0$, $\varepsilon_{np}=1.0$)

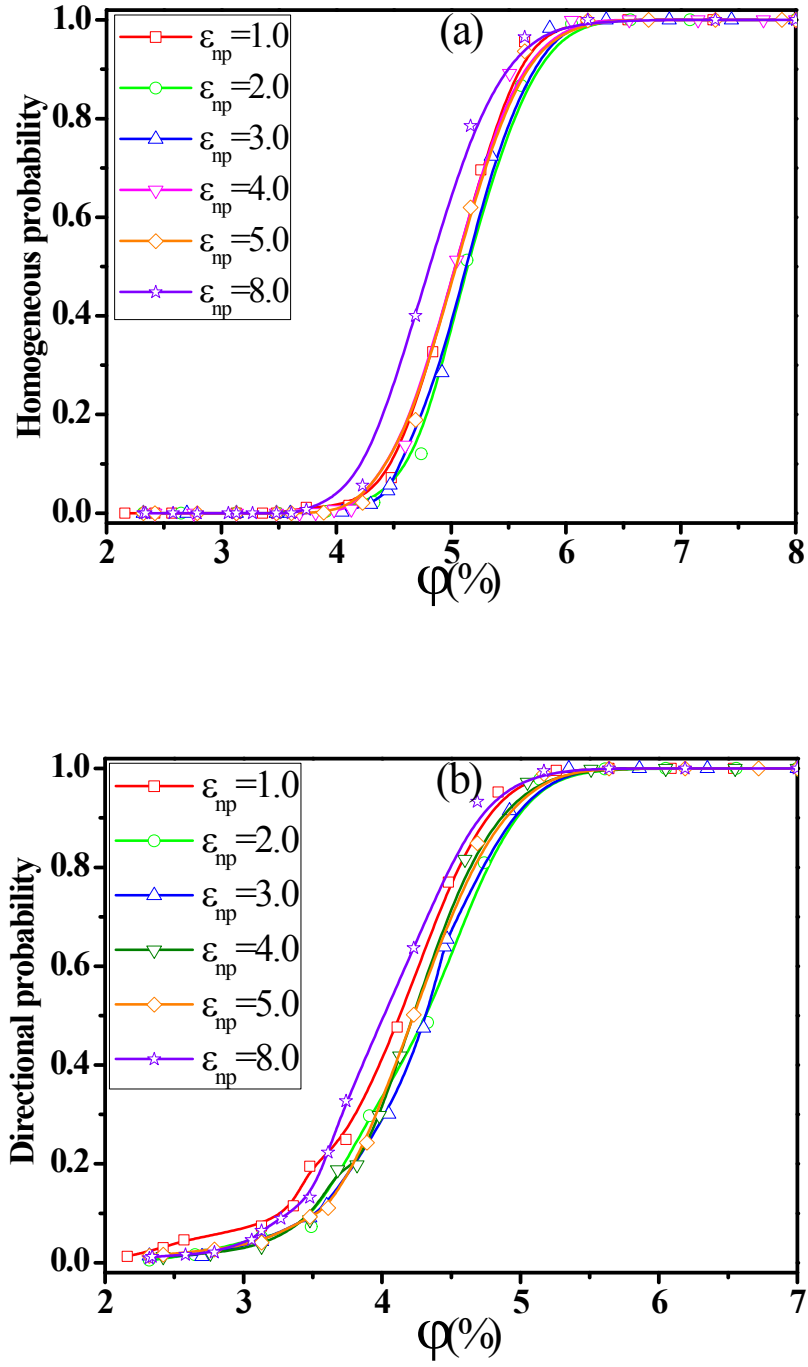


Fig. S7(a) Homogeneous probability and (b) directional probability of PNC as a function of the volume fraction ϕ of nanorods for different interfacial strengths ϵ_{np} . ($T=1.0$, $\gamma = 0.1$)

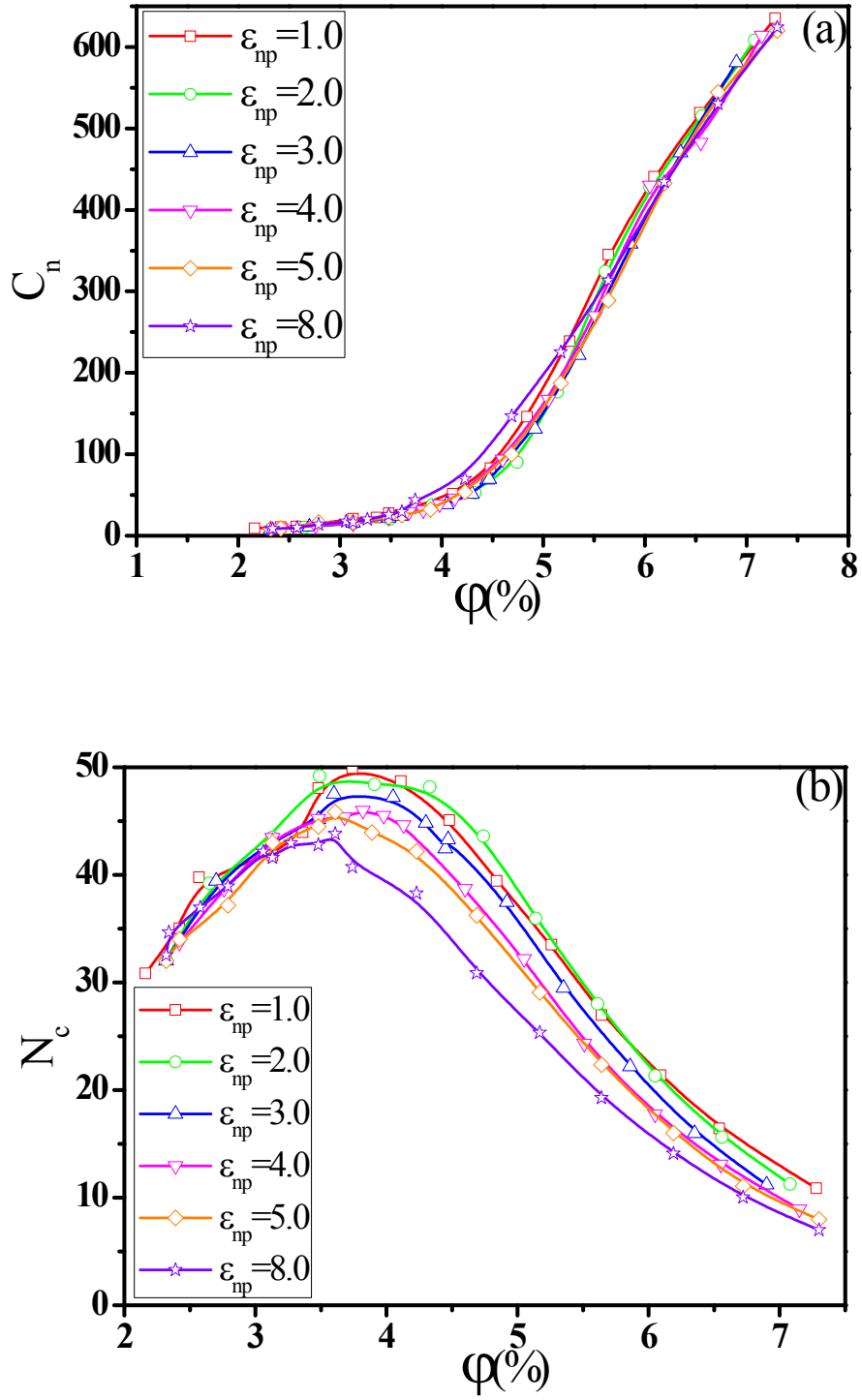


Fig. S8(a) The maximum cluster size C_n , (b) The total number of clusters N_c as a function of the volume fraction ϕ of nanorods for different interfacial strengths ϵ_{np} . ($T=1.0$, $\gamma = 0.1$)

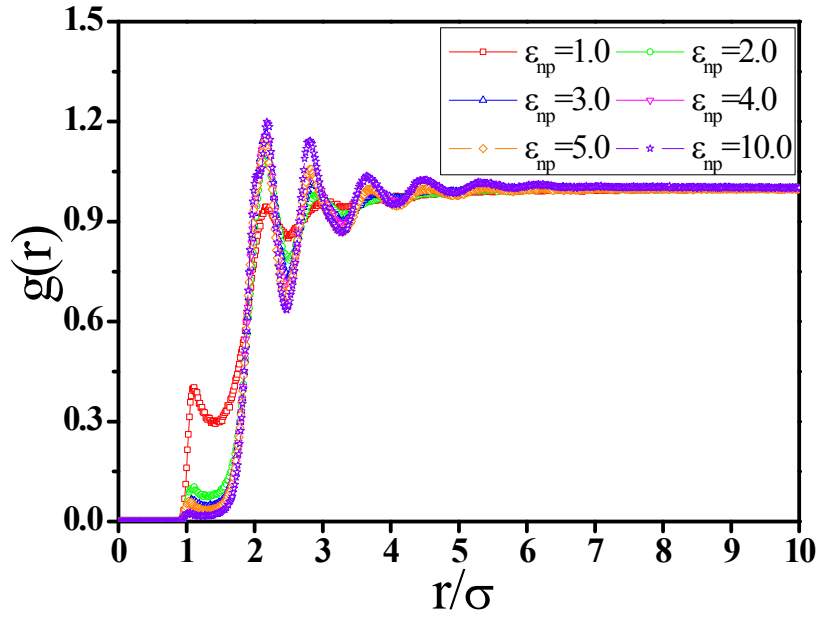
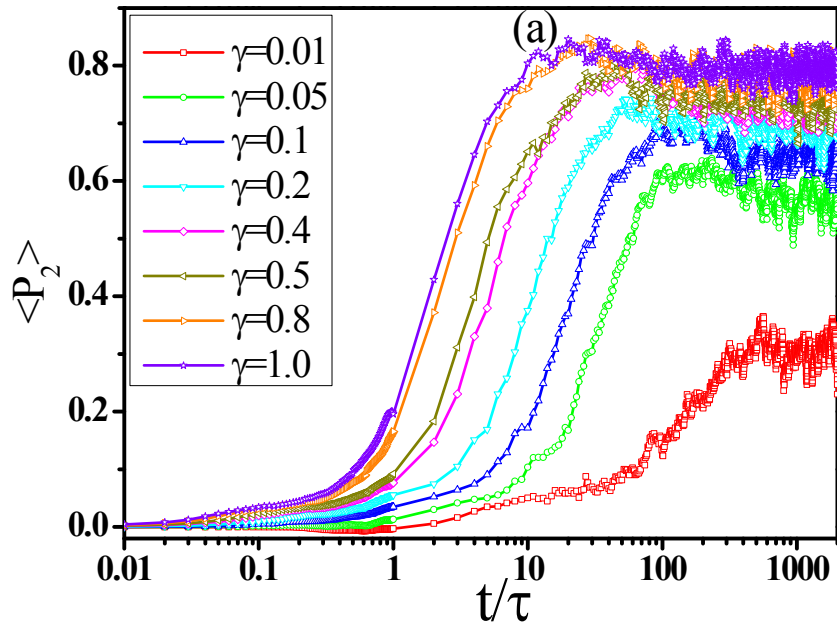


Fig. S9 RDF of nanorods for different interfacial interactions ε_{np} . ($T=1.0$, $\gamma=0.1$)



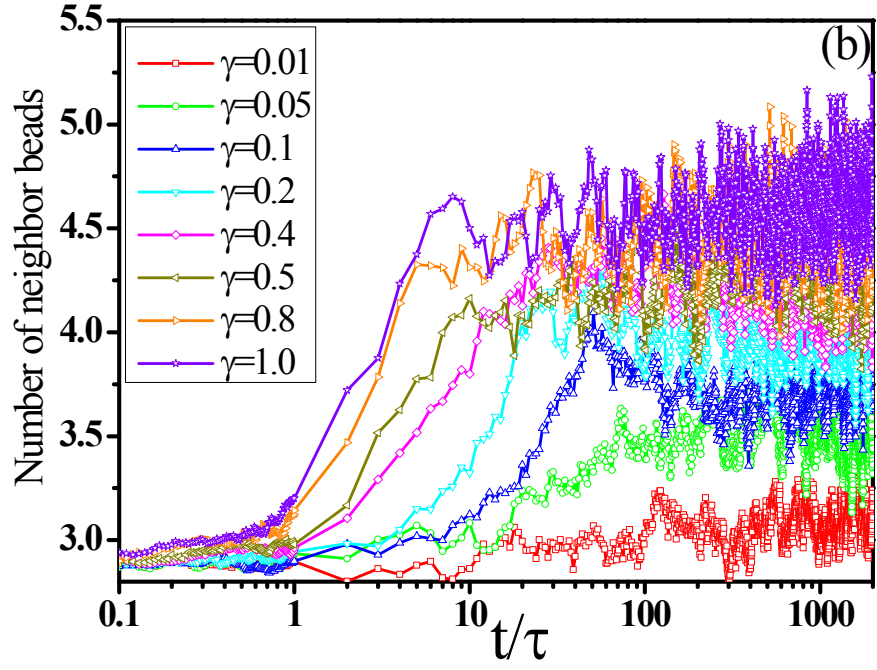


Fig. S10 The change of (a) the coordination number and (b) the second-order Legendre polynomials $\langle P_2 \rangle$ as a function of the shear time t for different shear rates $\dot{\gamma}$. ($T=1.0$, $\varepsilon_{np} = 1.0$, $\varphi = 4.49\%$)

Here we calculated the rotational diffusivity D_{rot} and the rotational Peclet number Pe_{rot}

($\text{Pe}_{\text{rot}} = \frac{\gamma}{D_{\text{rot}}}$). As shown in Fig. S11(a), for the case of $\varepsilon_{np} = 1.0$, the rotational

diffusivity D_{rot} increases with the shear rate, while the effect of the volume fraction is

little because of low interfacial interaction. And D_{rot} decreases significantly with the

increase of the interfacial interaction shown in Fig. S11(b) because stronger

interaction can inhibit the mobility of nanorods. From Fig. S12(a), the values of Pe_{rot}

for different volume fractions are nearly the same, while they increase with the shear

rate. Meanwhile, the percolation threshold of the directional conductivity also

increases with the shear rate, which is similar to the trend of Pe_{rot} . From Fig. S12(b),

the effect of the interfacial interaction on Pe_{rot} is small except the case $\varepsilon_{np} = 1.0$. This

is consistent with the percolation threshold of conductivity for different interactions.

The reason why Pe_{rot} is low at the case $\varepsilon_{np} = 1.0$ results from the fact that some

nanorods aggregate directly (the peak at $r = 2\sigma$ in Fig. S9), which leads to the low

rotational diffusivity D_{rot} . And for other interfacial interactions, nanorods do not

aggregate directly.

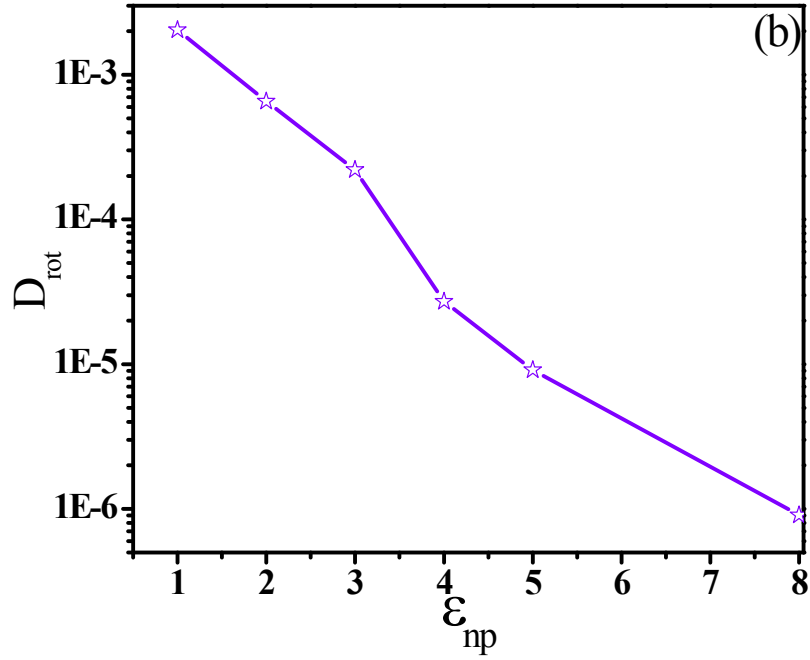
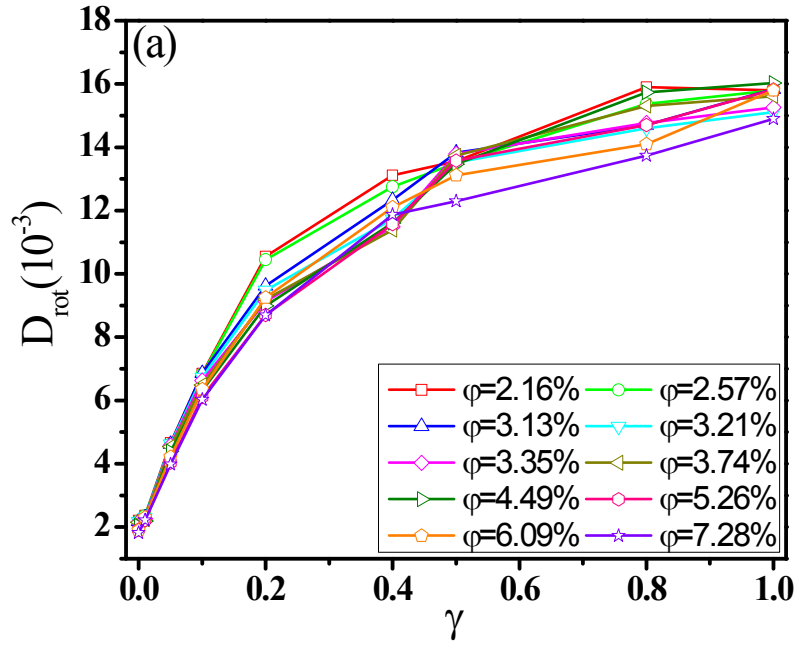


Fig. S11 The rotational diffusivity D_{rot} as a function of (a) shear rate γ for different volume fractions φ ($\varepsilon_{np}=1.0$) and (b) the interfacial interaction ε_{np} ($\gamma=0.0$, $\varphi=4.49\%$). ($T=1.0$)

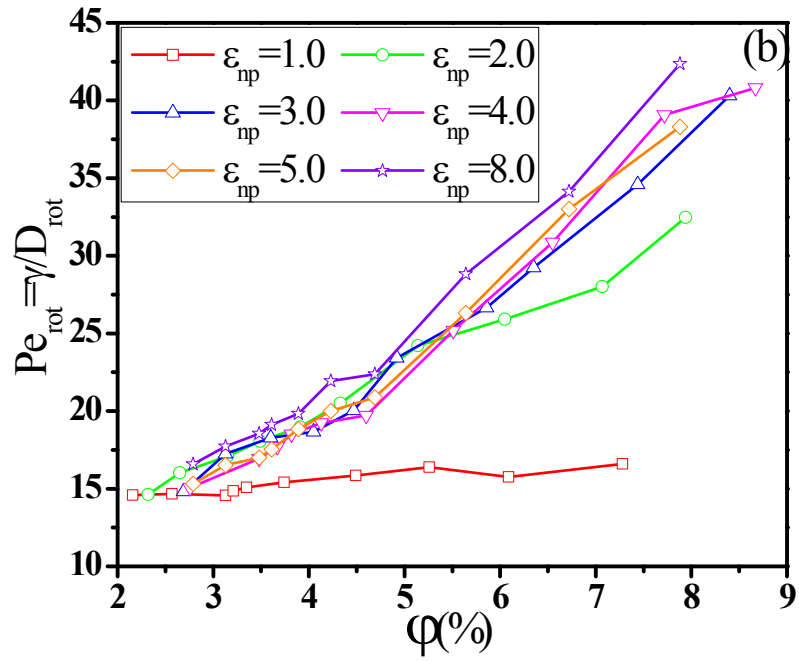
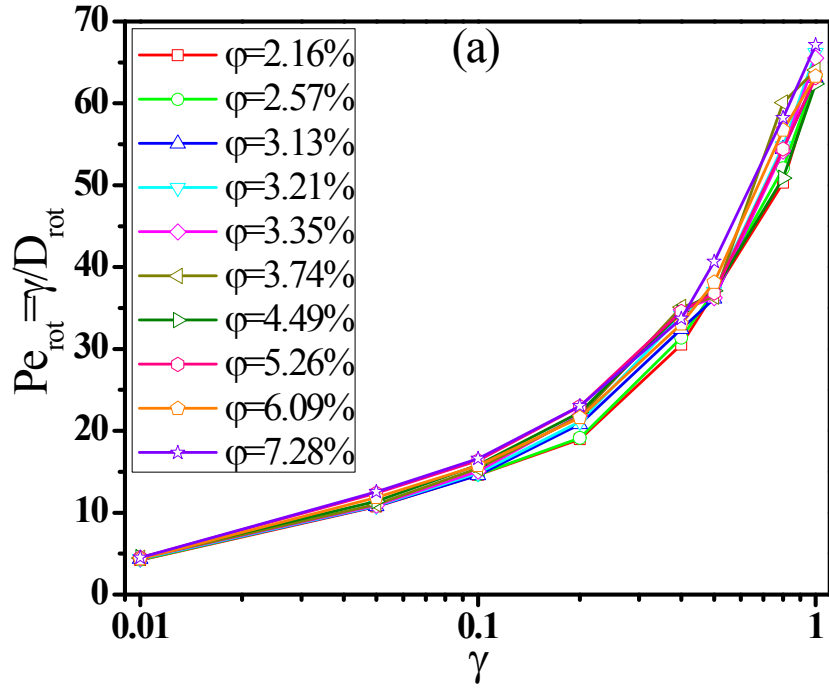


Fig. S12 The rotational Peclet number (a) as a function of shear rate γ for different volume fractions ϕ ($\epsilon_{np}=1.0$). (b) as a function of the volume fraction ϕ for different interfacial interaction ϵ_{np} with $\gamma=0.1$. ($T=1.0$)

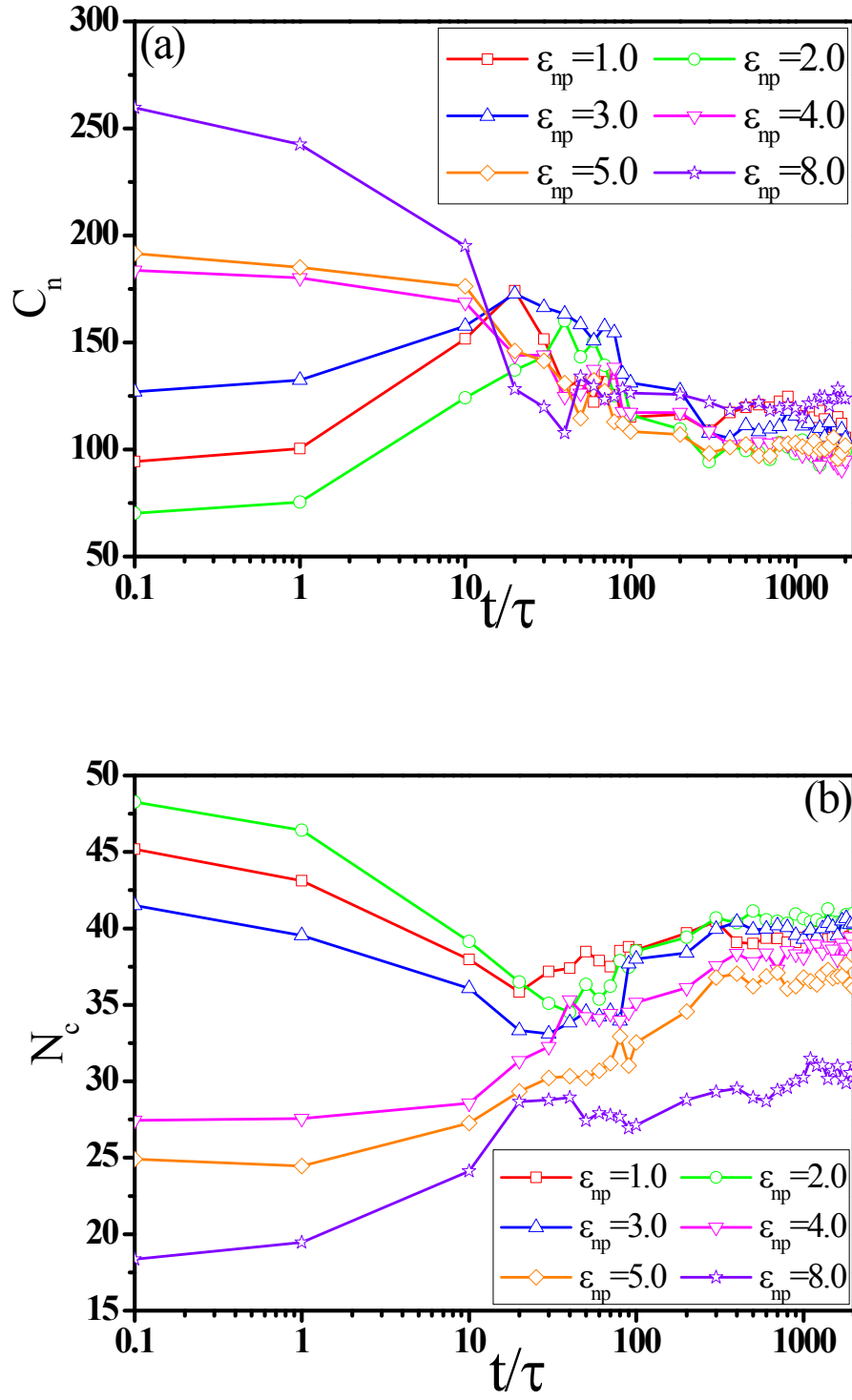


Fig. S13(a) The maximum cluster size C_n , (b) The total number of clusters N_c as a function of the shear time t for different interfacial strengths ε_{np} . ($T=1.0$, $\gamma = 0.1$, $\varphi = 4.49\%$)