

## Electronic Supplementary Information

### Charge carrier dynamics of methylammonium lead iodide: From $\text{PbI}_2$ -rich to low-dimensional broadly emitting perovskites

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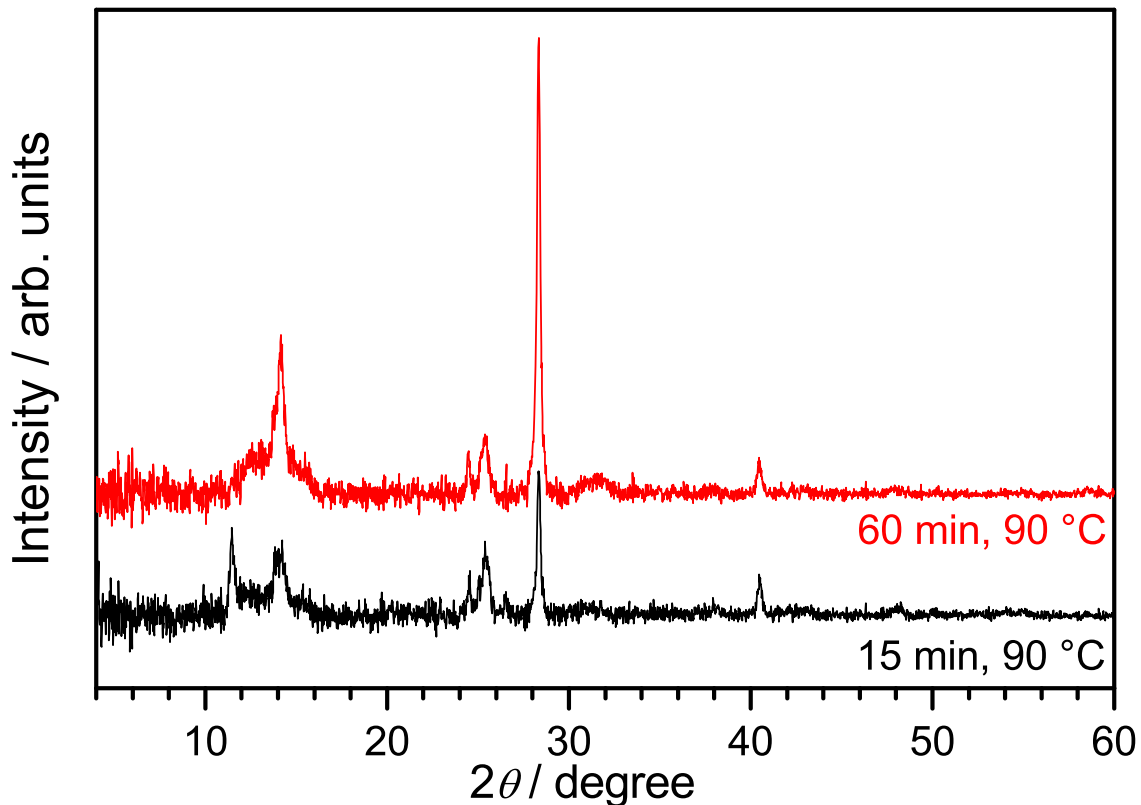
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## S1. Changes of XRD patterns for $x = 0.71$ upon annealing

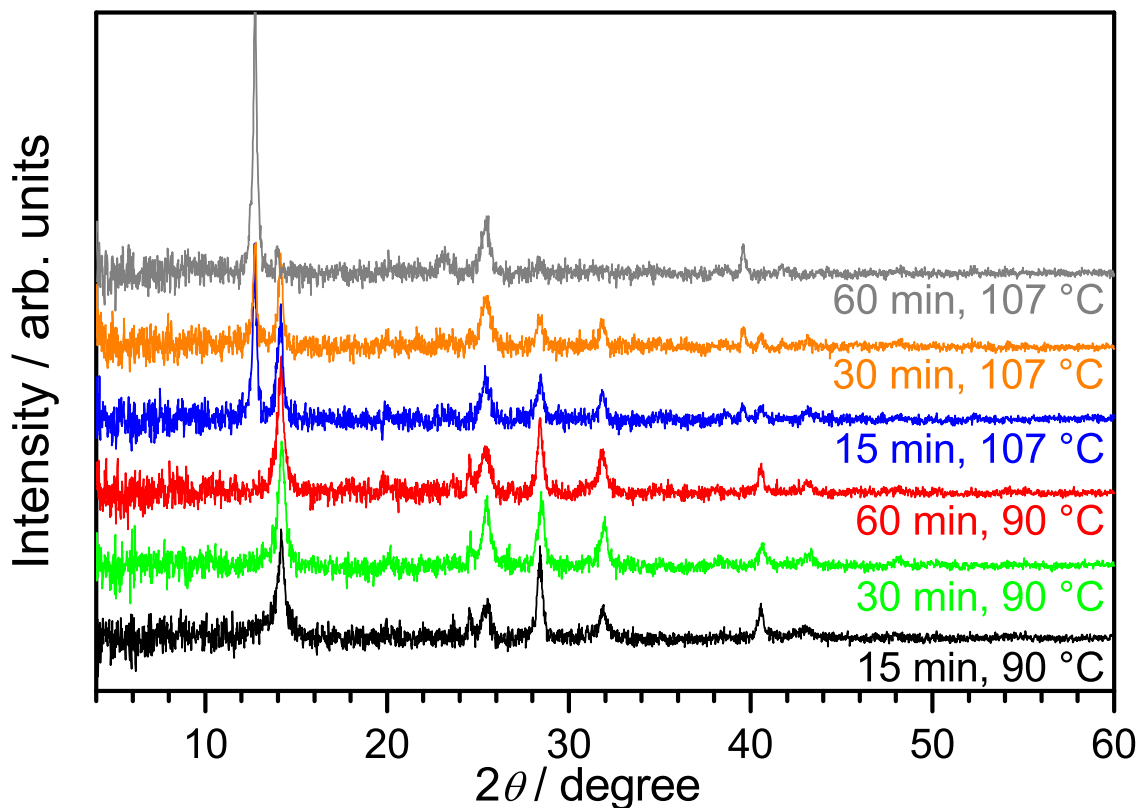
In Fig. S1, we present two XRD patterns of a  $(\text{MAI})_x(\text{PbI}_2)_{1-x}$  thin film with  $x = 0.71$  during heating at  $90^\circ\text{C}$ . According to the “pseudobinary phase-composition processing diagram” of Song *et al.*,<sup>1</sup> a mixture of stacked perovskite sheet (SPS) and low-dimensional perovskite (LDP) structures should be formed initially. Indeed, after 15 min (black line) we see perovskite XRD peaks at  $14.1^\circ$  and  $28.4^\circ$  and an LDP peak at  $11.5^\circ$ . Heating for 60 min (red line) leads to a decay of the LDP peak and an increase of the peak at  $28.4^\circ$ . The behavior is quite similar to that observed by Song *et al.* (see Fig. S3b in their Supporting Information).<sup>1</sup> It appears as if a thermally induced structural rearrangement or transformation of the LDP phase occurs, possibly resulting in smaller isolated 3D perovskite structures. In addition, an increased disorder of the LDPs might contribute to the reduction of the peak at  $11.5^\circ$ .



**Fig. S1** XRD patterns of a  $(\text{MAI})_x(\text{PbI}_2)_{1-x}/\text{TiO}_2/\text{glass}$  sample with  $x = 0.71$ . The black and red XRD patterns are obtained during annealing at  $90^\circ\text{C}$  for 15 and 60 min, respectively.

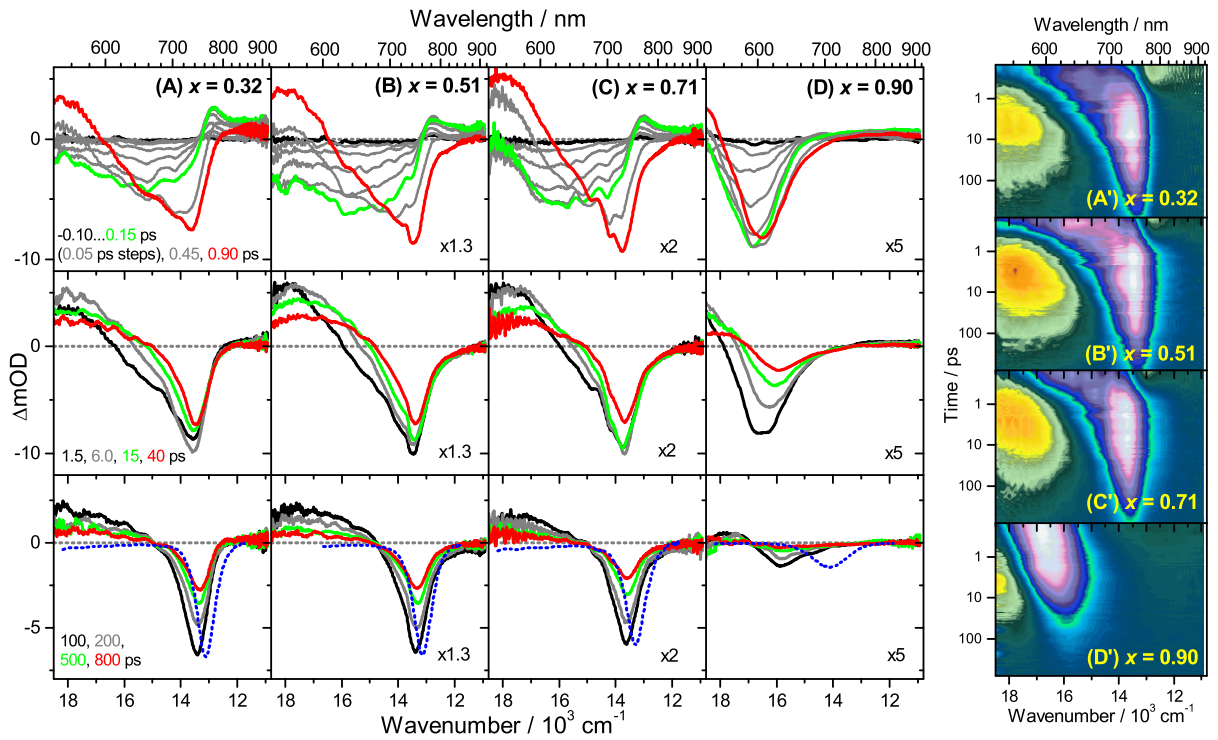
## S2. Optimized annealing condition monitored by XRD

The optimized annealing condition was found to be heating the  $(\text{MAI})_x(\text{PbI}_2)_{1-x}$  thin films at 90 °C for 60 min. No perovskite degradation was taking place under these conditions, as demonstrated by the black, green and red XRD patterns, taken at three different times, shown in Fig. S2 for  $x = 0.50$ . Increasing the hotplate temperature to 107 °C leads to  $\text{PbI}_2$  formation already after 15 min, with almost complete conversion of the perovskite into  $\text{PbI}_2$  reached after 60 min, as also shown in Fig. S2. The decomposition process was further confirmed by recording complementary transient absorption spectra of these films. After 60 min, only a weak residual absorption at 520 nm with a long tail extending into the NIR is visible, consistent with absorption features of  $\text{PbI}_2$ .<sup>2</sup>



**Fig. S2** XRD patterns of  $(\text{MAI})_x(\text{PbI}_2)_{1-x}/\text{TiO}_2/\text{glass}$  samples with  $x = 0.50$ . The black, green and red XRD patterns are obtained during annealing at 90 °C for 15, 30 and 60 min, respectively. The blue, orange and grey XRD patterns are for annealing at 107 °C for 15, 30 and 60 min, respectively.

### S3. Transient Vis-NIR spectra at the pump fluence $14 \mu\text{J cm}^{-2}$



**Fig. S3** Transient Vis-NIR PSCP broadband absorption spectra of  $(\text{MAI})_x(\text{PbI}_2)_{1-x}/\text{TiO}_2/\text{glass}$  at a pump pulse fluence of  $14 \mu\text{J cm}^{-2}$  for different initial mole fractions  $x$ . Columns (A) to (D) are the results for  $x = 0.32, 0.51, 0.71$  and  $0.90$ , respectively. (Top panels)  $-0.10 \dots 0.15$  ps with 50 fs steps; 0.45 ps, 0.90 ps; (middle panels) 1.5, 6.0, 15 and 40 ps; (bottom panels) 100, 200, 500 and 800 ps. Selected transient spectra are shown as thick colored lines for guidance. The blue-dotted lines in the bottom panels are the corresponding scaled steady-state stimulated emission spectra. (A'-D') show the dynamics of the systems in (A-D) as contour plots using a logarithmic time scale. White, pink and bluish colors denote bleach or stimulated emission, whereas green, yellow and red colors correspond to transient absorption.

## S4. Analysis of recombination kinetics

As described in the main manuscript, the second-order recombination rate constant  $k_2$  was obtained *via*:

$$\frac{1}{n(t)} - \frac{1}{n_0} = k_2 t \quad \text{with} \quad n_0 = \alpha F / E_{\text{ph}} \quad (\text{S1})$$

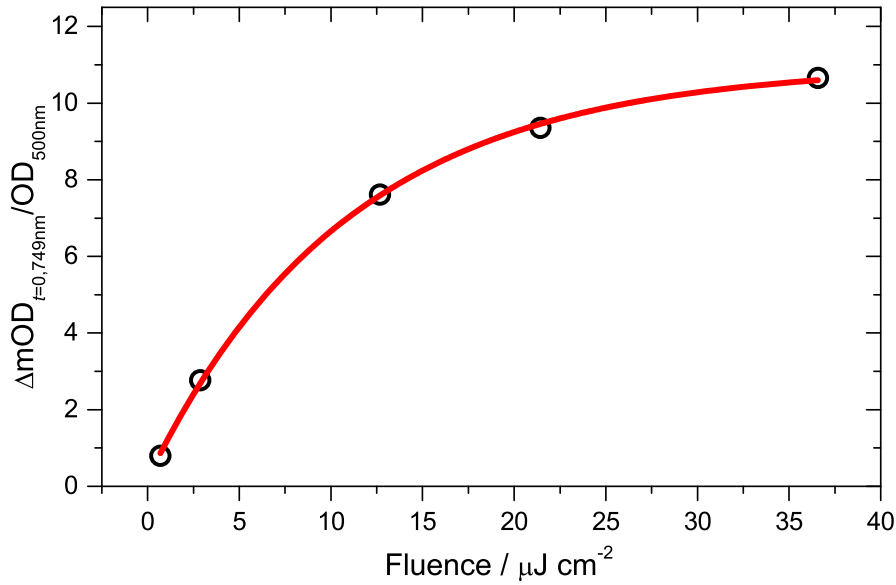
Here,  $n_0$  and  $n(t)$  are the number density of carriers in  $\text{cm}^{-3}$  at  $t = 0$  and at a given time  $t$ , respectively,  $k_2$  is the bimolecular rate constant in  $\text{cm}^3 \text{s}^{-1}$ ,  $\alpha$  is the absorption coefficient in  $\text{cm}^{-1}$  at the pump wavelength (500 nm),  $F$  is the fluence of the pump laser beam in  $\mu\text{J cm}^{-2}$  and  $E_{\text{ph}}$  is the pump photon energy in  $\mu\text{J}$  at 500 nm.

As it turned out, the maximum  $\Delta\text{OD}$  around 750 nm does not linearly depend on the fluence. A similar saturation behavior was previously found by Kamat and co-workers for two semiconductor systems.<sup>3,4</sup> Therefore, an empirical calibration curve for  $\Delta\text{OD}/\text{OD}$  as a function of the number density of carriers  $n(t)$  had to be established. This was done as follows:

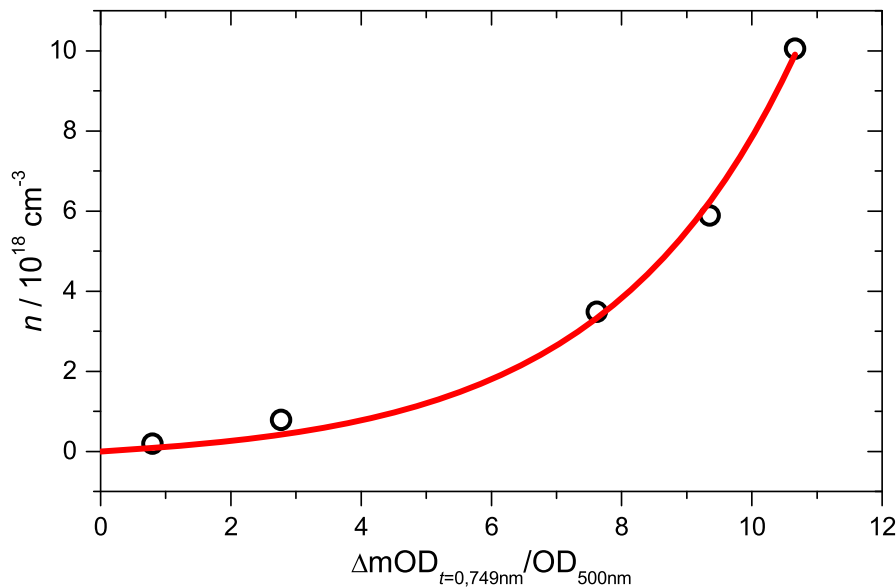
The  $1/\Delta\text{OD}_{t=0}$  value around 750 nm was determined at a certain pump fluence, by extrapolating the  $1/\Delta\text{OD}(t)$  data for  $t > 10$  ps to  $t = 0$ . This way, the influence of carrier cooling effects on the recombination kinetics was minimized. We normalized  $\Delta\text{OD}_{t=0}$  with respect to the OD at the pump wavelength 500 nm. The normalized  $\Delta\text{OD}_{t=0}$  values were then plotted as a function of  $F$ , resulting in typical saturation curves, as shown for one example ( $x = 0.51$ ) in Fig. S4.

Next, the experimental  $\Delta\text{OD}(t)$  values were converted into  $n(t)$  with the corresponding calibration curve. Here we took  $n(t)$  to be  $F \cdot \alpha / E_{\text{ph}}$  in units of  $\text{cm}^{-3}$ . One result based on the data in Fig. S4 is shown in Fig. S5. The expression  $y = A \cdot [\exp(B \cdot x^C) - 1]$  was used for fitting the data points.

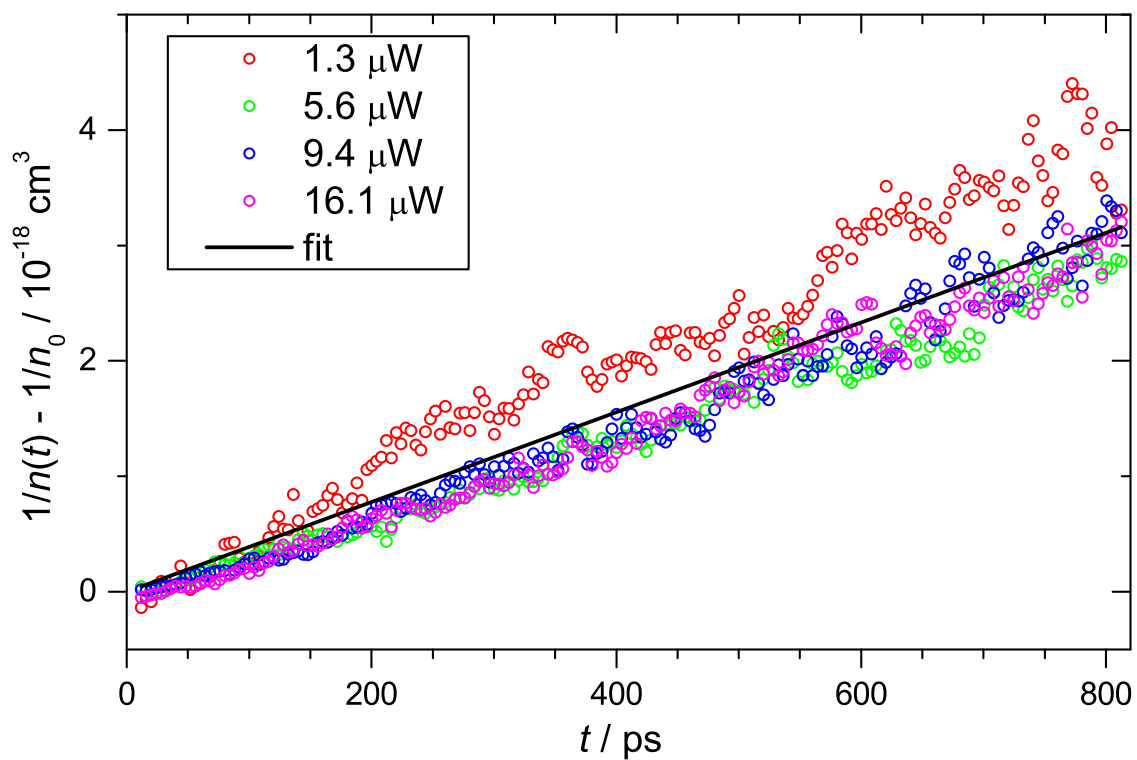
The final second-order kinetics plot was then obtained from eq. (S1). An example of such a representation is shown in Fig. S6 for  $x = 0.51$  at different pump fluences in the range 0.7-50  $\mu\text{J cm}^{-2}$ . All lines are largely coincident in the fluence range 2-50  $\mu\text{J cm}^{-2}$ . At high fluence (e.g.  $> 50 \mu\text{J cm}^{-2}$ ), there is a deviation from linearity at early times ( $< 100$  ps). This is possibly due to some contribution of additional carrier recombination processes, which are not following simple second-order kinetics. From the global fit of the data at different fluences, we finally obtain the recombination rate constant  $k_2 = (3.5 \pm 0.2) \times 10^{-9} \text{ cm}^3 \text{ s}^{-1}$  for  $x = 0.51$ . We performed the same analysis for  $x = 0.32, 0.71$  and  $0.90$ . The results are shown in Fig. 8 of the main manuscript.



**Fig. S4** Normalized  $\Delta\text{mOD}_{t=0}$  values as a function of the fluence of the pump laser pulse for an  $(\text{MAI})_x(\text{PbI}_2)_{1-x}/\text{TiO}_2/\text{glass}$  sample with mole fraction  $x = 0.51$ . The fluences of the pump laser beam are:  $F [\mu\text{J cm}^{-2}] = 0.69, 2.9, 12.7, 21.4$  and  $36.6$ .



**Fig. S5** An example of a calibration curve for converting experimental  $\Delta\text{mOD}(t)$  into  $n(t)$  values. The same experimental data as in Fig. S4 were used. Here we take  $n_0$  to be  $F\alpha$ , with the fluence in units of  $\text{cm}^{-2}$ . The corresponding  $n_0$  values are: (304 nW)  $1.9 \times 10^{17} \text{ cm}^{-3}$ , (1.3  $\mu\text{W}$ )  $7.9 \times 10^{17} \text{ cm}^{-3}$ , (5.6  $\mu\text{W}$ )  $3.5 \times 10^{18} \text{ cm}^{-3}$ , (9.4  $\mu\text{W}$ )  $5.9 \times 10^{18} \text{ cm}^{-3}$  and (16.1  $\mu\text{W}$ )  $1.0 \times 10^{19} \text{ cm}^{-3}$ . The red solid line is the fit result.



**Fig. S6** Plot of the difference of the reciprocal transient ( $n$ ) and initial ( $n_0$ ) carrier density as a function of time for an  $(\text{MAI})_x(\text{PbI}_2)_{1-x}/\text{TiO}_2/\text{glass}$  sample employing the mole fraction  $x = 0.51$ . The black line represents the global fit to all data points.

## References

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