

Supplementary Information for *Single-ion 4f element magnetism: an ab-initio look at Ln(COT)₂⁻*

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Optimized Distances of $\text{Ln}(\text{COT})_2^-$

Table S1: Comparison of selected optimized distances (\AA) along the $\text{Ln}(\text{COT})_2^-$ series. SF ZORA/B3LYP/TZ2P. Fractional occupation $(4f_\sigma 4f_\pi 4f_\delta 4f_\phi)^n$.

Complex	Ln-COT	Ln-C	C-C
$\text{Ce}(\text{COT})_2^-$	2.133	2.817	1.408
$\text{Pr}(\text{COT})_2^-$	2.123	2.809	1.408
$\text{Nd}(\text{COT})_2^-$	2.131	2.815	1.408
$\text{Pm}(\text{COT})_2^-$	2.156	2.834	1.408
$\text{Sm}(\text{COT})_2^-$	2.156	2.834	1.408
$\text{Eu}(\text{COT})_2^-$	2.101	2.793	1.408
$\text{Gd}(\text{COT})_2^-$	2.023	2.734	1.408
$\text{Tb}(\text{COT})_2^-$	2.004	2.721	1.408
$\text{Dy}(\text{COT})_2^-$	1.999	2.716	1.408
$\text{Ho}(\text{COT})_2^-$	2.009	2.724	1.408
$\text{Er}(\text{COT})_2^-$	2.032	2.741	1.408
$\text{Tm}(\text{COT})_2^-$	2.030	2.740	1.408
$\text{Yb}(\text{COT})_2^-$	1.981	2.704	1.408

Table S2: $\text{Ce}(\text{COT})_2^-$: Principal optimized distances (\AA) and relative energies (cm^{-1}) of the SF and SO states using three different fractional occupation schemes^a.

	Frac17	Frac13	NoFrac
Ce-COT	2.133	2.134	2.141
Ce-C	2.817	2.817	2.822
C-C	1.408	1.408	1.408
Spin-Free 2F term			
${}^2\Sigma$	0	0	0
${}^2\Pi$	518	518	518
${}^2\Phi$	621	621	612
${}^2\Delta$	1917	1917	1897
Spin-Orbit ${}^2F_{5/2}$ term			
$\pm 1/2$	0	0	0
$\pm 5/2$	479	480	469
$\pm 3/2$	1040	1040	1032

^a Frac17: $(4f_\sigma 4f_\pi 4f_\delta 4f_\phi)^1$ Frac13: $(4f_\sigma 4f_\pi)^1$ NoFrac: $4f_\sigma^1$

Model vs. *ab-initio* for $\text{Ce}(\text{COT})_2^-$ and $\text{Pr}(\text{COT})_2^-$

Table S3: Model wave functionsⁱ $|\psi\rangle$, orbital $m_u^L(\mathbf{r})$ and spin $m_u^S(\mathbf{r})$ magnetizations for $\text{Ce}(\text{COT})_2^-$ and $\text{Pr}(\text{COT})_2^-$.

		$\text{Ce}(\text{COT})_2^-$	
ψ	$= A^2\Sigma_{1/2} - B^2\Pi_{1/2}$ ($u = \parallel$)	ψ	$= \frac{1}{\sqrt{2}}[A^2\Sigma_{1/2} - B^2\Pi_{1/2} + A^2\Sigma_{-1/2} - B^2\Pi_{-1/2}]$ ($u = \perp$)
	$= AY_3^0 - B\bar{Y}_3^1$	ψ	$= \frac{1}{\sqrt{2}}[AY_3^0 - B\bar{Y}_3^1 + A\bar{Y}_3^0 - BY_3^{-1}]$
$m_{\parallel}^S(\mathbf{r})$	$= \frac{1}{2}[A^2Y_3^0Y_3^0 + B^2Y_3^1Y_3^{-1}]$	$m_{\perp}^S(\mathbf{r})$	$= \frac{A^2}{2}[Y_3^0Y_3^0] - \frac{B^2}{4}[Y_3^1Y_3^1 + Y_3^{-1}Y_3^{-1}]$
	$= \frac{A^2}{2}f_0^2 - \frac{B^2}{4}[f_{1+}^2 + f_{1-}^2]$		$= \frac{A^2}{2}f_0^2 - \frac{B^2}{4}[f_{1-}^2 - f_{1+}^2]$
$m_{\parallel}^L(\mathbf{r})$	$= -B^2Y_3^1Y_3^{-1}$	$m_{\perp}^L(\mathbf{r})$	$= -\frac{AB\sqrt{12}}{2}Y_3^0Y_3^0 + \frac{AB\sqrt{3}}{4}[2Y_3^1Y_3^{-1} + Y_3^1Y_3^1 + Y_3^{-1}Y_3^{-1}]$
	$= \frac{B^2}{2}[f_{1+}^2 + f_{1-}^2]$		$= -\frac{AB\sqrt{12}}{2}f_0^2 - AB\sqrt{3}f_{1+}^2$
		$\text{Pr}(\text{COT})_2^-$	
ψ	$= A^3\Gamma_{-3} + B^3\Phi_{-3} + C^3\Delta_{-3}$ ($u = \parallel$)		
	$= A Y_3^{-3}Y_3^{-1} + \frac{B}{\sqrt{2}}[\sqrt{\frac{2}{3}}(Y_3^{-3}\bar{Y}_3^0 - \bar{Y}_3^{-3}Y_3^0) + \sqrt{\frac{1}{3}}(Y_3^{-2}\bar{Y}_3^{-1} - \bar{Y}_3^{-2}Y_3^{-1})]$		
	$+ C[\sqrt{\frac{2}{3}} \bar{Y}_3^{-2}\bar{Y}_3^0 + \sqrt{\frac{1}{3}} \bar{Y}_3^{-3}\bar{Y}_3^1]$		
$m_{\parallel}^S(\mathbf{r})$	$= -\frac{A^2}{2}[Y_3^3Y_3^{-3} + Y_3^1Y_3^{-1}] - \frac{C^2}{3}[Y_3^2Y_3^{-2} + Y_3^0Y_3^0] + \frac{C^2}{6}[Y_3^3Y_3^{-3} + Y_3^1Y_3^{-1}]$		
	$= \frac{A^2}{4}[f_{3+}^2 + f_{3-}^2 + f_{1+}^2 + f_{1-}^2] - \frac{C^2}{6}[f_{2+}^2 + f_{2-}^2] - \frac{C^2}{3}f_0^2 - \frac{C^2}{6}[f_{3+}^2 + f_{3-}^2 + f_{1+}^2 + f_{1-}^2]$		
$m_{\parallel}^L(\mathbf{r})$	$= 2A^2[Y_3^3Y_3^{-3} + Y_3^1Y_3^{-1}] - \frac{3B^2}{2}[\frac{2}{3}[Y_3^0Y_3^0 - Y_3^3Y_3^{-3}] + \frac{1}{3}[Y_3^2Y_3^{-2} - Y_3^1Y_3^{-1}]]$		
	$- C^2[\frac{2}{3}[Y_3^2Y_3^{-2} + Y_3^0Y_3^0] - \frac{1}{3}[Y_3^3Y_3^{-3} + Y_3^1Y_3^{-1}]]$		
	$= -A^2[f_{3+}^2 + f_{3-}^2 + f_{1+}^2 + f_{1-}^2] - \frac{3B^2}{2}[\frac{2}{3}[f_0^2 + \frac{1}{2}(f_{3+}^2 + f_{3-}^2)] + \frac{1}{6}[f_{2+}^2 + f_{2-}^2 + f_{1+}^2 + f_{1-}^2]]$		
	$- C^2[\frac{2}{3}[f_0^2 + \frac{1}{2}(f_{2+}^2 + f_{2-}^2)] + \frac{1}{6}[f_{3+}^2 + f_{3-}^2 + f_{1+}^2 + f_{1-}^2]]$		

ⁱ Notation: Angular behavior only. Spherical harmonics $Y_{\ell}^{m_{\ell}}$, a bar indicates beta spin. Tesseral harmonics $f_{\ell\pm}^{|m|}$. Two-electron Slater determinants $|a, b| = 1/\sqrt{2} \det |a, b|$. Real coefficients A, B, C . Normalization implies $A^2 + B^2 = 1$ for $\text{Ce}(\text{COT})_2^-$ and $A^2 + B^2 + C^2 = 1$ for $\text{Pr}(\text{COT})_2^-$.

The spherical harmonics $Y_{\ell}^{m_{\ell}}$ are related to the real tesseral harmonics $f_{\ell\pm}^{|m|}$ as follows:

$$\begin{aligned}
 Y_3^0 &= f_0 \\
 Y_3^1 &= \frac{f_{1+} - if_{1-}}{i\sqrt{2}} \\
 Y_3^{-1} &= \frac{f_{1+} + if_{1-}}{i\sqrt{2}} \\
 Y_3^2 &= \frac{f_{2+} + if_{2-}}{\sqrt{2}} \\
 Y_3^{-2} &= \frac{f_{2+} - if_{2-}}{\sqrt{2}}
 \end{aligned}$$

$$Y_3^3 = \frac{f_{3+} - if_{3-}}{i\sqrt{2}}$$

$$Y_3^3 = \frac{f_{3+} + if_{3-}}{i\sqrt{2}}$$

Table S4: Comparison of the orbital and spin expectation values obtained with the model and with the *ab-initio* calculations for $\text{Ce}(\text{COT})_2^-$ and $\text{Pr}(\text{COT})_2^-$.

	Model	<i>ab-initio</i>
	$\text{Ce}(\text{COT})_2^-$	
$\langle L_{\parallel} \rangle$	0.459	0.461
$\langle S_{\parallel} \rangle$	0.040	0.038
$\pm g_{\parallel}$	1.078	1.074
$\langle L_{\perp} \rangle$	-1.726	-1.716
$\langle S_{\perp} \rangle$	0.269	0.269
$\pm g_{\perp}$	2.376	2.354
	$\text{Pr}(\text{COT})_2^-$	
$\langle L_{\parallel} \rangle$	-3.578	-3.609
$\langle S_{\parallel} \rangle$	0.609	0.613
$\pm g_{\parallel}$	4.720	4.765

Energies and assignment of the lowest Spin-Free electronic states in $\text{Ln}(\text{COT})_2^-$

Table S5: Relative energies (cm^{-1}) and assignment (per-cent) of the SF states issued from the SF ground multiplet ^{2s+1}L . CAS(n,7)SCF calculations.

^{2s+1}L Multiplet	M_L States ^a	ΔE	Major Configurations ^b
$\text{Ce}(\text{COT})_2^-$			
2F	$^2\Sigma$	0	100 σ
	$^2\Pi$	518	100 π
	$^2\Phi$	621	100 ϕ
	$^2\Delta$	1917	100 δ
$\text{Pr}(\text{COT})_2^-$			
3H	$^3\Phi$	0	76 ($\sigma\phi$), 22 ($\delta\pi$)
	$^3\Gamma$	4	93 ($\phi\pi$)
	$^3\Delta$	421	61 ($\sigma\delta$), 36 ($\phi\pi$)
	$^3\Sigma$	637	72 ($\pi\pi$), 26 ($\delta\delta$)
	$^3\Pi$	644	58 ($\pi\delta$), 28 ($\sigma\pi$)
	3H	1298	100 ($\phi\delta$)
$\text{Nd}(\text{COT})_2^-$			
4I	$^4\Phi$	0	76 ($\pi\pi\phi$), 15 ($\phi\delta\delta$), 8 ($\sigma\pi\delta$)
	$^4\Gamma$	28	51 ($\pi\phi\delta$), 48 ($\sigma\pi\phi$)
	$^4\Delta$	288	48 ($\pi\phi\delta$), 27 ($\sigma\pi\phi$), 26 ($\pi\pi\delta$)
	4H	413	100 ($\sigma\phi\delta$)
	$^4\Pi$	606	39 ($\sigma\phi\delta$), 25 ($\pi\delta\delta$), 23 ($\sigma\pi\delta$), 12 ($\pi\phi\phi$)
	$^4\Sigma$	732	53 ($\sigma\delta\delta$), 22 ($\pi\phi\delta$), 19 ($\sigma\phi\phi$), 6 ($\sigma\pi\pi$)
	4I	758	100 ($\pi\phi\delta$)
$\text{Pm}(\text{COT})_2^-$			
5I	5I	0	100 ($\sigma\pi\phi\delta$)
	$^5\Sigma$	26	57 ($\pi\pi\phi\phi$), 20 ($\sigma\pi\phi\delta$), 16 ($\pi\pi\delta\delta$), 5 ($\phi\phi\delta\delta$)
	$^5\Pi$	129	37 ($\pi\pi\phi\delta$), 28 ($\sigma\pi\phi\phi$), 20 ($\pi\phi\phi\delta$), 11 ($\sigma\pi\delta\delta$)
	5H	319	100 ($\pi\pi\phi\delta$)
	$^5\Delta$	363	46 ($\sigma\pi\phi\delta$), 27 ($\sigma\phi\phi\delta$), 22 ($\pi\phi\delta\delta$)
	$^5\Phi$	560	69 ($\sigma\phi\delta\delta$), 21 ($\sigma\pi\pi\phi$), 9 ($\pi\phi\phi\delta$)
	$^5\Gamma$	577	54 ($\sigma\pi\phi\delta$), 42 ($\pi\phi\delta\delta$)
$\text{Sm}(\text{COT})_2^-$			
6H	6H	0	100 ($\sigma\pi\pi\phi\delta$)
	$^6\Sigma$	379	49 ($\sigma\phi\phi\delta\delta$), 47 ($\sigma\pi\pi\phi\phi$)
	$^6\Pi$	473	64 ($\sigma\pi\phi\phi\delta$), 20 ($\pi\phi\phi\delta\delta$), 15 ($\sigma\pi\pi\phi\delta$)
	$^6\Delta$	629	69 ($\pi\pi\phi\phi\delta$), 31 ($\sigma\pi\phi\delta\delta$)
	$^6\Phi$	802	60 ($\pi\pi\phi\delta\delta$), 40 ($\sigma\pi\phi\phi\delta$)
	$^6\Gamma$	861	100 ($\sigma\pi\phi\delta\delta$)
$\text{Eu}(\text{COT})_2^-$			
7F	$^7\Delta$	0	100 ($\sigma\pi\pi\phi\phi\delta$)
	$^7\Phi$	525	100 ($\sigma\pi\pi\phi\delta\delta$)
	$^7\Pi$	871	100 ($\sigma\pi\phi\phi\delta\delta$)
	$^7\Sigma$	1102	100 ($\pi\pi\phi\phi\delta\delta$)

^a: The SF states characterized by the total angular momentum projection $M_L = \sum m_\ell$ are denoted using the $D_\infty h$ parent symmetry. ^b: Numbers indicate the weight of the configurations in %, only contributions larger than 5% are listed.

Table S6: Relative energies (cm⁻¹) and assignment (per-cent) of the SF states issued from the SF ground multiplet ^{2s+1}L. CAS(n,7)SCF calculations.

^{2s+1} L Multiplet	M _L States ^a	ΔE	Major Configurations ^b
⁸ S	⁸ Σ	0	Gd(COT) ₂ ⁻ 100 (σππφφδδ)
			Tb(COT) ₂ ⁻ 100(σ ² π ¹ π ¹ φ ¹ φ ¹ δ ¹ δ ¹) 100(σ ¹ π ² π ¹ φ ¹ φ ¹ δ ¹ δ ¹) 100(σ ¹ π ¹ π ¹ φ ² φ ¹ δ ¹ δ ¹) 100(σ ¹ π ¹ π ¹ φ ² φ ¹ δ ¹ δ ¹)
⁷ F	⁷ Σ	0	
	⁷ Π	268	
	⁷ Φ	675	
	⁷ Δ	1147	
⁶ H	⁶ Φ	0	Dy(COT) ₂ ⁻ 70 (σ ² π ¹ π ¹ φ ² φ ¹ δ ¹ δ ¹), 30 (σ ¹ π ² π ¹ φ ¹ φ ¹ δ ² δ ¹)
	⁶ Γ	35	100 (σ ¹ π ² π ¹ φ ² φ ¹ δ ¹ δ ¹)
	⁶ Δ	152	65 (σ ² π ¹ π ¹ φ ¹ φ ¹ δ ² δ ¹), 34 (σ ¹ π ² π ¹ φ ² φ ¹ δ ¹ δ ¹)
	⁶ Σ	217	66 (σ ¹ π ² π ² φ ¹ φ ¹ δ ¹ δ ¹), 31 (σ ¹ π ¹ π ¹ φ ¹ φ ¹ δ ² δ ²)
	⁶ Π	223	63 (σ ¹ π ² π ¹ φ ¹ φ ¹ δ ² δ ¹), 27 (σ ² π ² π ¹ φ ¹ φ ¹ δ ¹ δ ¹), 9 (σ ² π ¹ π ¹ φ ² φ ¹ δ ² δ ¹)
	⁶ H	865	100 (σ ¹ π ¹ π ¹ φ ² φ ¹ δ ² δ ¹)
⁵ I	⁵ Φ	0	Ho(COT) ₂ ⁻ 75 (σ ¹ π ² π ² φ ² φ ¹ δ ¹ δ ¹), 16 (σ ¹ π ¹ π ¹ φ ² φ ¹ δ ² δ ²), 9 (σ ² π ¹ π ² φ ¹ φ ¹ δ ² δ ¹)
	⁵ Γ	33	53 (σ ¹ π ² π ¹ φ ² φ ¹ δ ² δ ¹), 47 (σ ² π ² π ¹ φ ² φ ¹ δ ¹ δ ¹)
	⁵ Δ	139	44 (σ ¹ π ² π ¹ φ ² φ ¹ δ ² δ ¹), 27 (σ ¹ π ² π ² φ ¹ φ ¹ δ ² δ ¹), 26 (σ ² π ² π ¹ φ ² φ ¹ δ ¹ δ ¹)
	⁵ H	266	100 (σ ² π ¹ π ¹ φ ² φ ¹ δ ² δ ¹)
	⁵ Π	303	38 (σ ² π ¹ π ¹ φ ² φ ¹ δ ² δ ¹), 26 (σ ¹ π ¹ π ² φ ¹ φ ¹ δ ² δ ²), 23 (σ ² π ² π ¹ φ ¹ φ ¹ δ ² δ ¹) 12 (σ ¹ π ² π ¹ φ ² φ ² δ ¹ δ ¹)
	⁵ Σ	373	55 (σ ² π ¹ π ¹ φ ¹ φ ¹ δ ² δ ²), 21 (σ ¹ π ² π ¹ φ ² φ ¹ δ ² δ ¹), 18 (σ ² π ¹ π ¹ φ ² φ ² δ ¹ δ ¹) 100 (σ ¹ π ² π ¹ φ ² φ ¹ δ ² δ ¹)
⁴ I	⁴ I	0	Er(COT) ₂ ⁻ 100 (σ ² π ² π ¹ φ ² φ ¹ δ ² δ ¹)
	⁴ Σ	103	56 (σ ¹ π ² π ² φ ² φ ² δ ¹ δ ¹), 21 (σ ² π ² π ¹ φ ² φ ¹ δ ² δ ¹), 17 (σ ¹ π ² π ² φ ¹ φ ¹ δ ² δ ²) 5 (σ ¹ π ¹ π ¹ φ ² φ ² δ ² δ ²)
	⁴ Π	163	37 (σ ¹ π ² π ² φ ² φ ¹ δ ² δ ¹), 28 (σ ² π ² π ¹ φ ² φ ² δ ² δ ²), 12 (σ ² π ² π ¹ φ ¹ φ ¹ δ ² δ ²)
	⁴ H	196	100 (σ ¹ π ² π ² φ ² φ ¹ δ ² δ ¹)
	⁴ Δ	300	46 (σ ² π ² π ¹ φ ² φ ¹ δ ² δ ¹), 27 (σ ² π ¹ π ¹ φ ² φ ² δ ¹ δ ²), 23 (σ ¹ π ² π ¹ φ ² φ ¹ δ ² δ ²)
	⁴ Γ	387	55 (σ ² π ² π ¹ φ ² φ ¹ δ ² δ ¹), 44 (σ ¹ π ² π ¹ φ ¹ φ ² δ ² δ ²)
	⁴ Φ	408	71 (σ ² π ¹ π ¹ φ ² φ ¹ δ ² δ ²), 20 (σ ² π ² π ² φ ² φ ¹ δ ² δ ²), 9 (σ ¹ π ² π ¹ φ ² φ ² δ ² δ ¹)
	³ H	³ H	0
³ Σ		500	53 (σ ² π ¹ π ¹ φ ² φ ² δ ² δ ²), 43 (σ ² π ² π ² φ ² φ ² δ ¹ δ ¹)
³ Π		533	64 (σ ² π ² π ¹ φ ² φ ² δ ² δ ¹), 21 (σ ¹ π ² π ¹ φ ² φ ² δ ² δ ²), 14 (σ ² π ² π ² φ ² φ ¹ δ ² δ ¹)
³ Δ		592	67 (σ ¹ π ² π ² φ ² φ ² δ ² δ ¹), 33 (σ ² π ² π ¹ φ ² φ ¹ δ ² δ ²)
³ Γ		615	100 (σ ² π ² π ¹ φ ² φ ¹ δ ² δ ²)
³ Φ		656	64 (σ ¹ π ² π ² φ ² φ ¹ δ ² δ ²), 36 (σ ² π ² π ¹ φ ² φ ² δ ² δ ¹)
² F	² Δ	0	Yb(COT) ₂ ⁻ 100(σ ² π ² π ² φ ² φ ² δ ² δ ¹)
	² Φ	182	100(σ ² π ² π ² φ ² φ ¹ δ ² δ ²)
	² Π	643	100(σ ² π ² π ¹ φ ² φ ² δ ² δ ²)
	² Σ	802	100(σ ¹ π ² π ¹ φ ² φ ² δ ² δ ²)

^a: The SF states characterized by the total angular momentum projection $M_L = \sum m_\ell$ are denoted using the $D_\infty h$ parent symmetry. ^b: Numbers indicate the weight of the configurations in %, only contributions larger than 5% are listed.

Energies and assignment of the lowest Spin-Orbit electronic states in $\text{Ln}(\text{COT})_2^-$

Table S7: Relative energies (cm^{-1}) and assignment (per-cent) of the SO states issued from the SO ground multiplet $^{2s+1}L_J$. CAS(n,7)SCF-SO calculations.

$^{2s+1}L_J$ Multiplet	M_J States	ΔE	Major Configurations ^a
$\text{Ce}(\text{COT})_2^-$			
$^2F_{5/2}$	$\pm 1/2$	0	54 $^2\Sigma$, 46 $^2\Pi$
	$\pm 5/2$	479	94 $^2\Phi$, 6 $^2\Delta$
	$\pm 3/2$	1040	59 $^2\Pi$, 41 $^2\Delta$
$\text{Pr}(\text{COT})_2^-$			
3H_4	± 3	0	65 $^3\Gamma$, 28 $^3\Phi$, 4 $^3\Delta$, 2 $^1\Phi$
	± 2	182	59 $^3\Phi$, 32 $^3\Delta$, 7 $^3\Pi$, 2 $^1\Delta$
	± 1	550	41 $^3\Delta$, 16 $^3\Sigma$, 18 $^3\Pi$, 2 $^1\Pi$
	0	645	44 $^3\Sigma$, 54 $^3\Pi$, 2 $^1\Sigma$
	± 4	894	6 $^3\Phi$, 37 $^3\Gamma$, 54 3H , 2 $^1\Gamma$
$\text{Nd}(\text{COT})_2^-$			
$^4I_{9/2}$	$\pm 5/2$	0	37 $^4\Phi$, 44 $^4\Gamma$, 14 $^4\Delta$
	$\pm 7/2$	156	12 $^4\Phi$, 36 $^4\Gamma$, 48, 4H
	$\pm 3/2$	204	38 $^4\Phi$, 38 $^4\Delta$, 18 $^4\Pi$, 4 $^4\Sigma$
	$\pm 1/2$	488	24 $^4\Delta$, 46 $^4\Pi$, 27 $^4\Sigma$
	$\pm 9/2$	552	7 $^4\Gamma$, 24 4H , 66 4I
$\text{Pm}(\text{COT})_2^-$			
5I_4	± 4	0	78 5I , 17 5H , 4 $^5\Gamma$
	0	42	34 $^5\Sigma$, 49 $^5\Pi$, 17 $^5\Delta$
	± 1	144	22 $^5\Sigma$, 39 $^5\Pi$, 27 $^5\Delta$, 11 $^5\Phi$
	± 3	339	52 5H , 30 $^5\Gamma$, 13 $^5\Phi$, 4 $^5\Delta$
	± 2	352	31 $^5\Phi$, 27 $^5\Delta$, 23 $^5\Gamma$, 15 $^5\Pi$, 4 $^5\Sigma$
$\text{Sm}(\text{COT})_2^-$			
$^6H_{5/2}$	$\pm 5/2$	0	81 6H , 14 $^6\Gamma$, 4 $^6\Phi$
	$\pm 1/2$	276	42 $^6\Pi$, 26 $^6\Sigma$, 24 $^6\Delta$, 8 $^6\Phi$
	$\pm 3/2$	432	25 $^6\Delta$, 25 $^6\Phi$, 22 $^6\Pi$, 17 $^6\Gamma$, 11 $^6\Sigma$
$\text{Eu}(\text{COT})_2^-$			
7F_0	0	0	41 $^7\Delta$, 36 $^7\Phi$, $^7\Pi$, $^7\Sigma$
$\text{Gd}(\text{COT})_2^-$			
$^8S_{7/2}$	$\pm 7/2$	0	100 $^8\Sigma$
	$\pm 5/2$	73	100 $^8\Sigma$
	$\pm 3/2$	121	100 $^8\Sigma$
	$\pm 1/2$	146	100 $^8\Sigma$
$\text{Tb}(\text{COT})_2^-$			
7F_6	0	0	47 $^7\Sigma$, 45 $^7\Pi$, 5 $^7\Delta$
	± 1	40	46 $^7\Pi$, 40 $^7\Sigma$, 7 $^7\Delta$
	± 2	161	49 $^7\Pi$, 33 $^7\Sigma$, 13 $^7\Delta$
	± 3	356	49 $^7\Pi$, 26 $^7\Delta$, 16 $^7\Sigma$, 6 $^7\Phi$
	± 6	473	97 $^7\Phi$
	± 4	585	43 $^7\Delta$, 33 $^7\Pi$, 21 $^7\Phi$
	± 5	676	59 $^7\Phi$, 38 $^7\Delta$

^a: Only the SF ^{2s+1}L configurations larger than 5% are given.

Table S8: Relative energies (cm⁻¹) and assignment (per-cent) of the SO states issued from the SO ground multiplet ^{2s+1}L_J. CAS(n,7)SCF-SO calculations.

^{2s+1} L _J Multiplet	M _J States	ΔE	Major Configurations ^a
⁶ H _{15/2}			Dy(COT) ₂ ⁻
	± ⁹ / ₂	0	48 ⁶ Φ, 24 ⁶ Δ, 21 ⁶ Γ
	± ¹¹ / ₂	25	49 ⁶ Φ, 45 ⁶ Γ, 6 ⁶ H
	± ⁷ / ₂	38	40 ⁶ Δ, 33 ⁶ Φ, 13 ⁶ Π, 8 ⁶ Γ
	± ⁵ / ₂	89	39 ⁶ Δ, 32 ⁶ Π, 15 ⁶ Φ, 7 ⁶ Σ, 2 ⁶ Γ
	± ³ / ₂	128	44 ⁶ Π, 24 ⁶ Δ, 23 ⁶ Σ, 5 ⁶ Φ
	± ¹ / ₂	149	47 ⁶ Π, 38 ⁶ Σ, 10 ⁶ Δ, 1 ⁶ Φ
	± ¹³ / ₂	206	73 ⁶ Γ, 23 ⁶ H
	± ¹⁵ / ₂	785	96 ⁶ H
⁵ I ₈			Ho(COT) ₂ ⁻
	±5	0	45 ⁵ Γ, 38 ⁵ Φ, 11 ⁵ H
	±4	1	46 ⁵ Φ, 25 ⁵ Δ, 21 ⁵ Γ, 2 ⁵ H
	±3	63	42 ⁵ Δ, 30 ⁵ Φ, 16 ⁵ Π, 6 ⁵ Γ
	±6	85	54 ⁵ Γ, 36 ⁵ H, 4 ⁵ I
	±2	149	36 ⁵ Δ, 36 ⁵ Π, 11 ⁵ Φ, 10 ⁵ Σ
	±1	221	46 ⁵ Π, 30 ⁵ Σ, 17 ⁵ Δ, 2 ⁵ Φ
	±7	247	72 ⁵ H, 22 ⁵ I
	0	248	47 ⁵ Π, 40 ⁵ Σ, 8 ⁵ Δ
±8	403	95 ⁵ I	
⁴ I _{15/2}			Er(COT) ₂ ⁻
	± ¹⁵ / ₂	0	98 ⁴ I
	± ¹ / ₂	160	48 ⁴ Π, 42 ⁴ Σ, 7 ⁴ Δ
	± ¹³ / ₂	165	77 ⁴ H, 20 ⁴ I
	± ³ / ₂	214	47 ⁴ Π, 28 ⁴ Δ, 18 ⁴ Σ, 4 ⁴ Φ
	± ⁵ / ₂	300	48 ⁴ Δ, 26 ⁴ Π, 16 ⁴ Φ, 2 ⁴ Γ
	± ¹¹ / ₂	320	60 ⁴ Γ, 35 ⁴ H, 3 ⁴ I
	± ⁷ / ₂	374	47 ⁴ Φ, 36 ⁴ Δ, 14 ⁴ Γ
± ⁹ / ₂	392	47 ⁴ Φ, 43 ⁴ Γ, 8 ⁴ H	
³ H ₆			Tm(COT) ₂ ⁻
	±6	0	99 ³ H
	±5	501	80 ³ Γ, 19 ³ H
	0	519	55 ³ Σ, 44 ³ Π
	±1	536	53 ³ Π, 32 ³ Σ, 14 ³ Δ
	±2	576	45 ³ Δ, 42 ³ Π, 8 ³ Φ
	±3	619	55 ³ Δ, 40 ³ Φ, 4 ³ Γ
±4	629	67 ³ Φ, 30 ³ Γ, 1 ³ H	
² F _{7/2}			Yb(COT) ₂ ⁻
	± ⁵ / ₂	0	86 ² Δ, 14 ² Φ
	± ⁷ / ₂	156	100 ² Φ
	± ³ / ₂	424	69 ² Π, 31 ² Δ
± ¹ / ₂	710	56 ² Σ, 44 ² Π	

^a: Only the SF ^{2s+1}L configurations larger than 5% are given.

Crystal Field Theory

The crystal field Hamiltonian for a f complex in a D_{8h} symmetry can be written as:

$$\hat{H}^{CF} = B_2^0 C_2^0 + B_4^0 C_4^0 + B_6^0 C_6^0 \quad (\text{S1})$$

with

$$C_k^q(\theta, \phi) = \sqrt{\frac{4\pi}{2k+1}} Y_k^q(\theta, \phi) \quad (\text{S2})$$

and B_k^q some coefficients that include an integral over the radial part of the wavefunctions.

If the spin-orbit coupling is relatively large compared to the crystal field interaction, the magnetic behavior of the lanthanide ion can be described by considering only the ground multiplet $^{2S+1}L_J$. This restriction to the ground multiplet has the advantage that only the matrix elements $\langle JM_J | \hat{H}^{CF} | JM_J' \rangle$ are needed for the calculation of the energies of the low energy spectrum. Due to the spherical tensor structure of the \hat{H}^{CF} operator, the matrix elements $\langle JM_J | \hat{H}^{CF} | JM_J' \rangle$ are easily evaluated by application of the so-called operator equivalent technique.¹³ This technique is based on the fact that CF operators $C_k^{(q)} = \sqrt{\frac{4\pi}{2k+1}} Y_k^q$, written in spatial coordinates in terms of spherical harmonics, are equivalent to the operators $\hat{J}_x, \hat{J}_y, \hat{J}_z$ of the total angular momentum apart from a factor called Stevens coefficients.^{14,15} These factors $\beta(q)$ depend on the order q of the operator and on the number of f electrons in the shell and are tabulated.¹⁶

The Hamiltonian defined in Eq. S1 written in terms of Stevens operators becomes

$$\hat{H}^{CF} = \beta_J(2)[B_2^0 \tilde{O}_2^0] + \beta_J(4)[B_4^0 \tilde{O}_4^0] + \beta_J(6)[B_6^0 \tilde{O}_6^0] \quad (\text{S3})$$

where the $\beta_J(k)$ are the Stevens coefficients and are obtained using the Wigner-Eckard theorem within the J -manifold.¹⁴ The matrix elements $\langle JM_J | \tilde{O}_k^q \hat{H}^{CF} | JM_J' \rangle$ have been listed by Abragam and Bleaney,¹⁶ Stevens^{14,15} and Judd¹⁷ for the $k = 0$ terms. The equivalent operators \tilde{O}_k^q are defined as follows:¹⁴⁻¹⁶

$$\tilde{O}_2^0 = 3J_z^2 - J(J+1) \quad (\text{S4a})$$

$$\tilde{O}_4^0 = 35J_z^4 - 30J(J+1)J_z^2 + 25J_z^2 - 6J(J+1) + 3J^2J(J+1)^2 \quad (\text{S4b})$$

$$\begin{aligned} \tilde{O}_6^0 = & 231J_z^6 - 315J(J+1)J_z^4 + 735J_z^4 + 105J^2(J+1)^2J_z^2 - 525J(J+1)J_z^2 \\ & + 294J_z^2 - 5J^3J(J+1)^3 + 40J^2J(J+1)^2 - 60J(J+1) \end{aligned} \quad (\text{S4c})$$

For the linear CF, the operator matrix representation in the $|J, M_J\rangle$ basis is diagonal. The eigenvalues are therefore trivially obtained.

The n -fold degeneracy of each ground multiplet $^{2S+1}L_J$ is split under the influence of the linear

crystal field and the energies of each m_J components as a function of the CF parameters b_2 , b_4 and b_6 are given in Tables S9 and S10.

For brevity, we define the new crystal field parameters b_2 , b_4 and b_6 as follows:

$$b_2 = \beta_J(2) \cdot F(2) \cdot B_2^0 \quad (\text{S5a})$$

$$b_4 = \beta_J(4) \cdot F(4) \cdot B_4^0 \quad (\text{S5b})$$

$$b_6 = \beta_J(6) \cdot F(6) \cdot B_6^0 \quad (\text{S5c})$$

where:

$$B_2^0 = A_2^0 \cdot \langle r^2 \rangle \quad (\text{S6a})$$

$$B_4^0 = A_4^0 \cdot \langle r^4 \rangle \quad (\text{S6b})$$

$$B_6^0 = A_6^0 \cdot \langle r^6 \rangle \quad (\text{S6c})$$

and the numbers $F(n)$ are multiplying factors common to all $|J, M_J\rangle$ basis. The crystal field parameters B_2^0 , B_4^0 and B_6^0 are resumed in Tables S11 and S12

Table S9: Stevens coefficients $\beta_J(n)$, multiplying factors $F(n)$, and energies of the m_J states of each ground multiplets $^{2S+1}L_J$.

	Stevens Coefficients $\beta_J(n)$	$F(n)$	m_J States	Energy of m_J States
$\text{Ce}(\text{COT})_2^-$	$\beta_J(2) = -\frac{2}{35}$	$F(2) = 2$	$m_J = \pm 1/2$	$E = -4b_2 + 2b_4$
	$\beta_J(4) = \frac{2}{315}$	$F(4) = 60$	$m_J = \pm 3/2$	$E = -1b_2 - 3b_4$
			$m_J = \pm 5/2$	$E = 5b_2 + 1b_4$
$\text{Pr}(\text{COT})_2^-$	$\beta_J(2) = -\frac{52}{2475}$	$F(2) = 1$	$m_J = 0$	$E = -20b_2 + 18b_4 - 20b_6$
	$\beta_J(4) = -\frac{4}{5445}$	$F(4) = 60$	$m_J = \pm 1$	$E = -17b_2 + 9b_4 + 1b_6$
	$\beta_J(6) = \frac{272}{4459455}$	$F(6) = 1260$	$m_J = \pm 2$	$E = -8b_2 - 11b_4 + 22b_6$
			$m_J = \pm 3$	$E = 7b_2 - 21b_4 - 17b_6$
			$m_J = \pm 4$	$E = 28b_2 + 14b_4 + 4b_6$
$\text{Nd}(\text{COT})_2^-$	$\beta_J(2) = -\frac{7}{1089}$	$F(2) = 6$	$m_J = \pm 1/2$	$E = -4b_2 + 18b_4 - 8b_6$
	$\beta_J(4) = -\frac{136}{467181}$	$F(4) = 84$	$m_J = \pm 3/2$	$E = -3b_2 + 3b_4 + 6b_6$
	$\beta_J(6) = -\frac{1615}{42513471}$	$F(6) = 5040$	$m_J = \pm 5/2$	$E = -1b_2 - 17b_4 + 10b_6$
			$m_J = \pm 7/2$	$E = 2b_2 - 22b_4 - 11b_6$
			$m_J = \pm 9/2$	$E = 6b_2 + 18b_4 + 3b_6$
$\text{Pm}(\text{COT})_2^-$	$\beta_J(2) = \frac{14}{1815}$	$F(2) = 1$	$m_J = 0$	$E = -20b_2 + 18b_4 - 20b_6$
	$\beta_J(4) = \frac{952}{2335905}$	$F(4) = 60$	$m_J = \pm 1$	$E = -17b_2 + 9b_4 + 1b_6$
	$\beta_J(6) = \frac{2584}{42513471}$	$F(6) = 1260$	$m_J = \pm 2$	$E = -8b_2 - 11b_4 + 22b_6$
			$m_J = \pm 3$	$E = 7b_2 - 21b_4 - 17b_6$
			$m_J = \pm 4$	$E = 28b_2 + 14b_4 + 4b_6$
$\text{Sm}(\text{COT})_2^-$	$\beta_J(2) = \frac{13}{315}$	$F(2) = 2$	$m_J = \pm 1/2$	$E = -4b_2 + 2b_4$
	$\beta_J(4) = \frac{26}{10395}$	$F(4) = 60$	$m_J = \pm 3/2$	$E = -1b_2 - 3b_4$
			$m_J = \pm 5/2$	$E = 5b_2 + 1b_4$
$\text{Tb}(\text{COT})_2^-$	$\beta_J(2) = -\frac{1}{99}$	$F(2) = 3$	$m_J = 0$	$E = -14b_2 + 84b_4 - 40b_6$
	$\beta_J(4) = \frac{2}{16335}$	$F(4) = 60$	$m_J = \pm 1$	$E = -13b_2 + 64b_4 - 20b_6$
	$\beta_J(6) = -\frac{1}{891891}$	$F(6) = 7560$	$m_J = \pm 2$	$E = -10b_2 + 11b_4 + 22b_6$
			$m_J = \pm 3$	$E = -5b_2 - 54b_4 + 43b_6$
			$m_J = \pm 4$	$E = 2b_2 - 96b_4 + 8b_6$
			$m_J = \pm 5$	$E = 11b_2 - 66b_4 - 55b_6$
		$m_J = \pm 6$	$E = 22b_2 + 99b_4 + 22b_6$	

Table S10: Stevens coefficients $\beta_J(n)$, multiplying factors $F(n)$, and energies of the m_J states of each ground multiplets $^{2S+1}L_J$.

	Stevens Coefficients $\beta_J(n)$	$F(n)$	m_J States	Energy of m_J States
$\text{Dy}(\text{COT})_2^-$	$\beta_J(2) = -\frac{2}{315}$	$F(2) = 3$	$m_J = \pm 1/2$	$E = -21b_2 + 189b_4 - 75b_6$
	$\beta_J(4) = -\frac{8}{135135}$	$F(4) = 60$	$m_J = \pm 3/2$	$E = -19b_2 + 129b_4 - 25b_6$
	$\beta_J(6) = \frac{4}{3864861}$	$F(6) = 13860$	$m_J = \pm 5/2$	$E = -15b_2 + 23b_4 + 45b_6$
			$m_J = \pm 7/2$	$E = -9b_2 - 101b_4 + 87b_6$
			$m_J = \pm 9/2$	$E = -1b_2 - 201b_4 + 59b_6$
			$m_J = \pm 11/2$	$E = 9b_2 - 221b_4 - 39b_6$
			$m_J = \pm 13/2$	$E = 21b_2 - 91b_4 - 117b_6$
$m_J = \pm 15/2$	$E = 35b_2 + 273b_4 + 65b_6$			
$\text{Ho}(\text{COT})_2^-$	$\beta_J(2) = -\frac{1}{450}$	$F(2) = 3$	$m_J = 0$	$E = -24b_2 + 36b_4 - 120b_6$
	$\beta_J(4) = -\frac{1}{30030}$	$F(4) = 420$	$m_J = \pm 1$	$E = -23b_2 + 31b_4 - 85b_6$
	$\beta_J(6) = -\frac{5}{3864861}$	$F(6) = 13860$	$m_J = \pm 2$	$E = -20b_2 + 17b_4 + 2b_6$
			$m_J = \pm 3$	$E = -15b_2 - 3b_4 + 93b_6$
			$m_J = \pm 4$	$E = -8b_2 - 24b_4 + 128b_6$
			$m_J = \pm 5$	$E = 1b_2 - 39b_4 + 65b_6$
			$m_J = \pm 6$	$E = 12b_2 - 39b_4 - 78b_6$
			$m_J = \pm 7$	$E = 25b_2 - 13b_4 - 169b_6$
$m_J = \pm 8$	$E = 40b_2 + 52b_4 + 104b_6$			
$\text{Er}(\text{COT})_2^-$	$\beta_J(2) = \frac{4}{1575}$	$F(2) = 3$	$m_J = \pm 1/2$	$E = -21b_2 + 189b_4 - 75b_6$
	$\beta_J(4) = \frac{2}{45045}$	$F(4) = 60$	$m_J = \pm 3/2$	$E = -19b_2 + 129b_4 - 25b_6$
	$\beta_J(6) = \frac{8}{3864861}$	$F(6) = 13860$	$m_J = \pm 5/2$	$E = -15b_2 + 23b_4 + 45b_6$
			$m_J = \pm 7/2$	$E = -9b_2 - 101b_4 + 87b_6$
			$m_J = \pm 9/2$	$E = -1b_2 - 201b_4 + 59b_6$
			$m_J = \pm 11/2$	$E = 9b_2 - 221b_4 - 39b_6$
			$m_J = \pm 13/2$	$E = 21b_2 - 91b_4 - 117b_6$
$m_J = \pm 15/2$	$E = 35b_2 + 273b_4 + 65b_6$			
$\text{Tm}(\text{COT})_2^-$	$\beta_J(2) = \frac{1}{99}$	$F(2) = 3$	$m_J = 0$	$E = -14b_2 + 84b_4 - 40b_6$
	$\beta_J(4) = \frac{8}{49005}$	$F(4) = 60$	$m_J = \pm 1$	$E = -13b_2 + 64b_4 - 20b_6$
	$\beta_J(6) = -\frac{5}{891891}$	$F(6) = 7560$	$m_J = \pm 2$	$E = -10b_2 + 11b_4 + 22b_6$
			$m_J = \pm 3$	$E = -5b_2 - 54b_4 + 43b_6$
			$m_J = \pm 4$	$E = 2b_2 - 96b_4 + 8b_6$
			$m_J = \pm 5$	$E = 11b_2 - 66b_4 - 55b_6$
$m_J = \pm 6$	$E = 22b_2 + 99b_4 + 22b_6$			
$\text{Yb}(\text{COT})_2^-$	$\beta_J(2) = \frac{2}{63}$	$F(2) = 3$	$m_J = \pm 1/2$	$E = -5b_2 + 9b_4 - 5b_6$
	$\beta_J(4) = -\frac{2}{1155}$	$F(4) = 60$	$m_J = \pm 3/2$	$E = -3b_2 - 3b_4 + 9b_6$
	$\beta_J(6) = \frac{4}{27027}$	$F(6) = 1260$	$m_J = \pm 5/2$	$E = 1b_2 - 13b_4 - 5b_6$
			$m_J = \pm 7/2$	$E = 7b_2 + 7b_4 + 1b_6$

Table S11: Crystal-field parameters (cm^{-1}) obtained by a least squares fit to the SCF-SO energies.

	$\text{Ce}(\text{COT})_2^-$	$\text{Pr}(\text{COT})_2^-$	$\text{Nd}(\text{COT})_2^-$	$\text{Pm}(\text{COT})_2^-$	$\text{Sm}(\text{COT})_2^-$
B_2^0	-282	-273	-342	-313	-443
B_4^0	-495	-477	-434	-379	-354
B_6^0		23	43	22	

Table S12: Crystal-field parameters (cm^{-1}) obtained by a least squares fit to the SCF-SO energies.

	$\text{Tb}(\text{COT})_2^-$	$\text{Dy}(\text{COT})_2^-$	$\text{Ho}(\text{COT})_2^-$	$\text{Er}(\text{COT})_2^-$	$\text{Tm}(\text{COT})_2^-$	$\text{Yb}(\text{COT})_2^-$
B_2^0	-494	-456	-439	-430	-420	-466
B_4^0	-273	-278	-239	-220	-194	-194
B_6^0	72	34	21	18	14	17

In both complexes, the B_2^0 and B_4^0 parameters have a large negative value while B_6^0 is much smaller in magnitude. According to Equations S3 and S4a, a negative value of B_2^0 favors a GS with large M_J value when the Stevens coefficient $\beta_J(2)$ is positive (Pm, Sm, Er, Tm and Yb). On the contrary, small M_J components are favored for negative value of the Stevens coefficient $\beta_J(2)$ (Ce, Pr, Nd, Tb, Dy and Ho). Furthermore, according to Eq. S4b, a large value of B_4^0 adds a quadratic term which favors intermediate values of M_J . In the $\text{Ln}(\text{COT})_2^-$ series, B_2^0 increases in absolute value in the series while B_4^0 decreases. Therefore, small M_J values are characterized for the GS of $\text{Ce}(\text{COT})_2^-$ and $\text{Tb}(\text{COT})_2^-$, whereas large M_J values are found for $\text{Pm}(\text{COT})_2^-$, $\text{Sm}(\text{COT})_2^-$, $\text{Er}(\text{COT})_2^-$ and $\text{Tm}(\text{COT})_2^-$. Intermediate M_J values are characterized for the rest of the series (see Tables S7 and S8).

Plots of $m_{\perp}^L(\mathbf{r})$ and $m_{\perp}^S(\mathbf{r})$

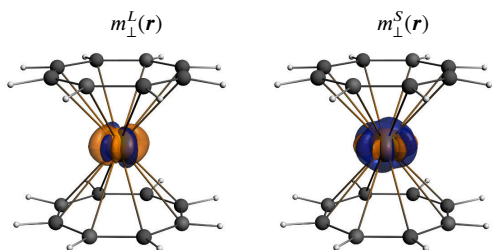


Figure S1: Orbital magnetization $m_{\perp}^L(\mathbf{r})$ and spin magnetization $m_{\perp}^S(\mathbf{r})$ density, along the \perp axis for $\text{Dy}(\text{COT})_2^-$. Isosurface values: ± 0.0001 au.

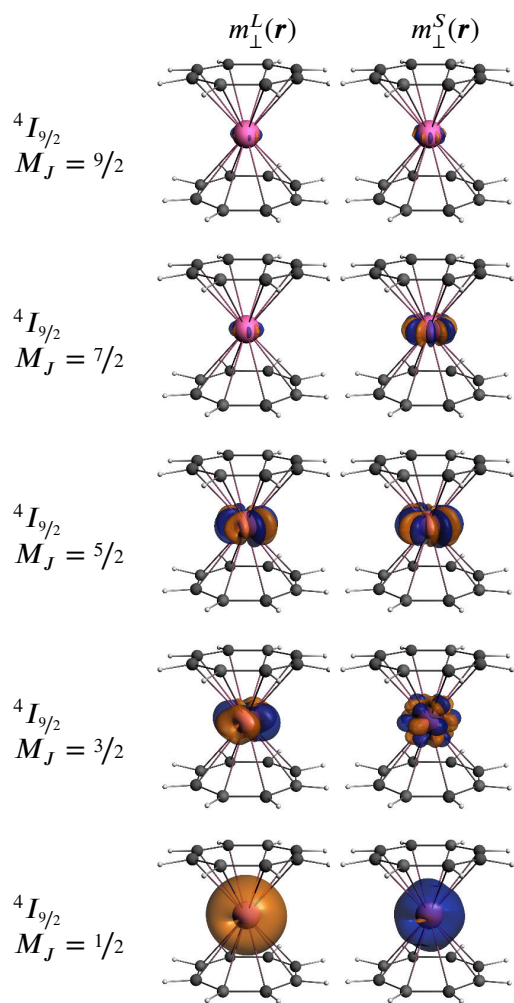


Figure S2: Orbital magnetization $m_{\perp}^L(\mathbf{r})$ and spin magnetization $m_{\perp}^S(\mathbf{r})$ density, along the \perp axis for $\text{Nd}(\text{COT})_2^-$. Isosurface values: ± 0.0001 au.

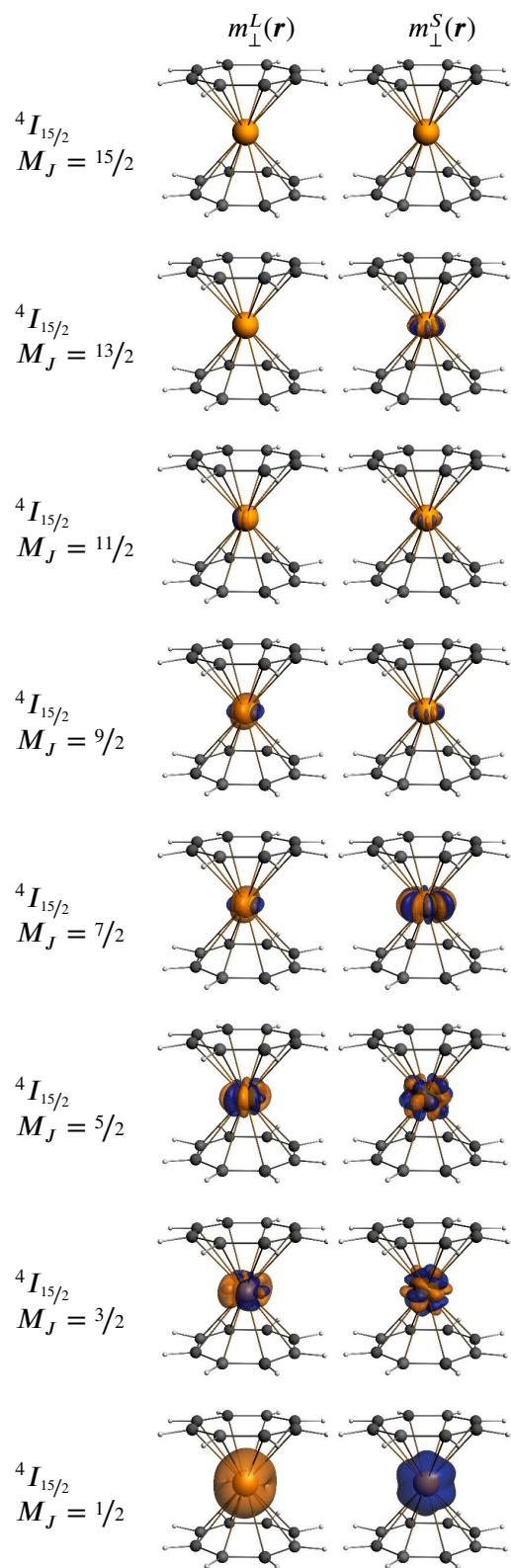


Figure S3: Orbital magnetization $m_{\perp}^L(\mathbf{r})$ and spin magnetization $m_{\perp}^S(\mathbf{r})$ density, along the \perp axis for $\text{Er}(\text{COT})_2^-$. Isosurface values: ± 0.0001 au.

Natural Orbitals of the ground states in $\text{Ln}(\text{COT})_2^-$

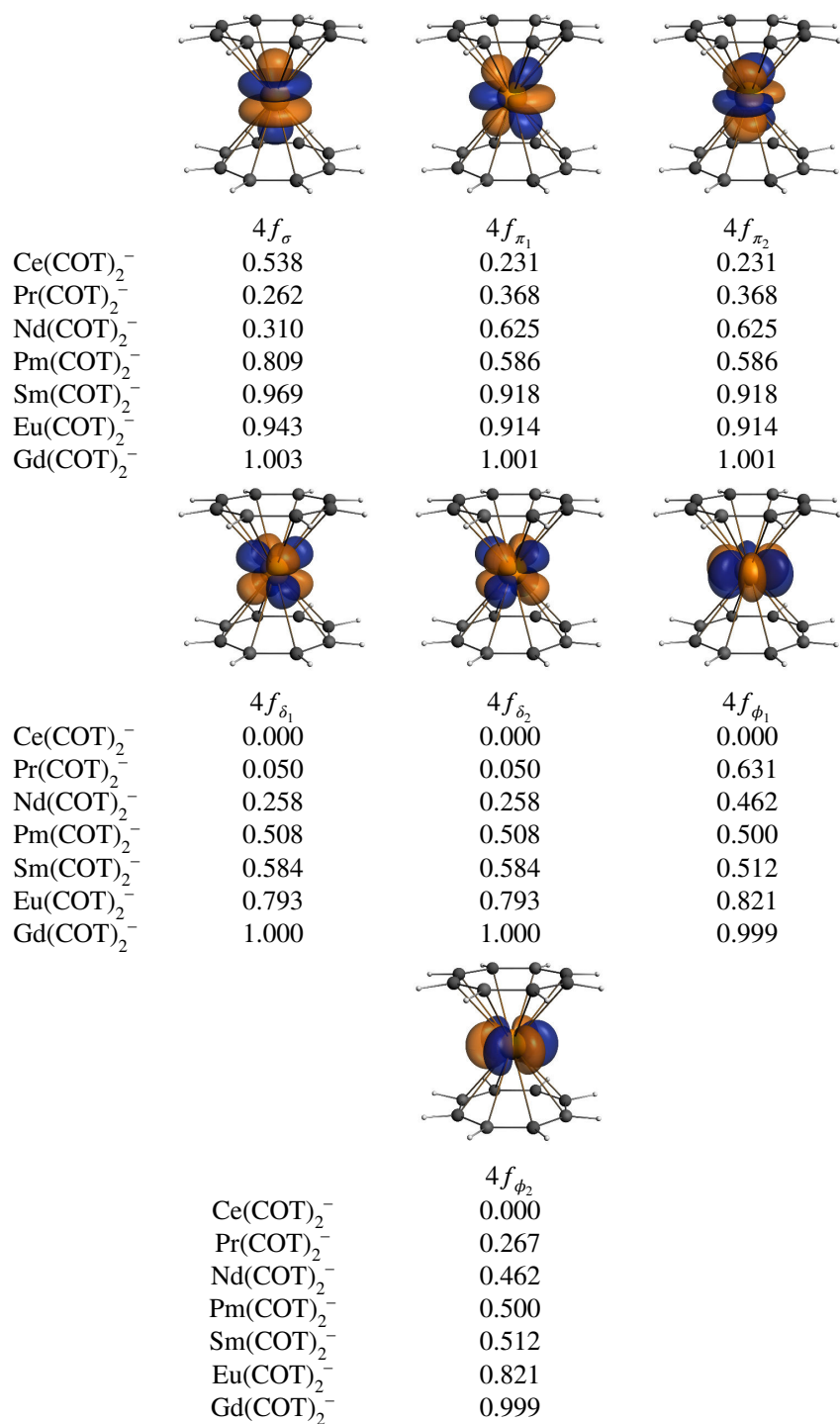


Figure S4: Selected NOs φ_p and occupation numbers n_p for $\text{Ln}(\text{COT})_2^-$. The figure shows the NOs from the SO calculation of $\text{Pr}(\text{COT})_2^-$. Isosurface values: ± 0.03 au.

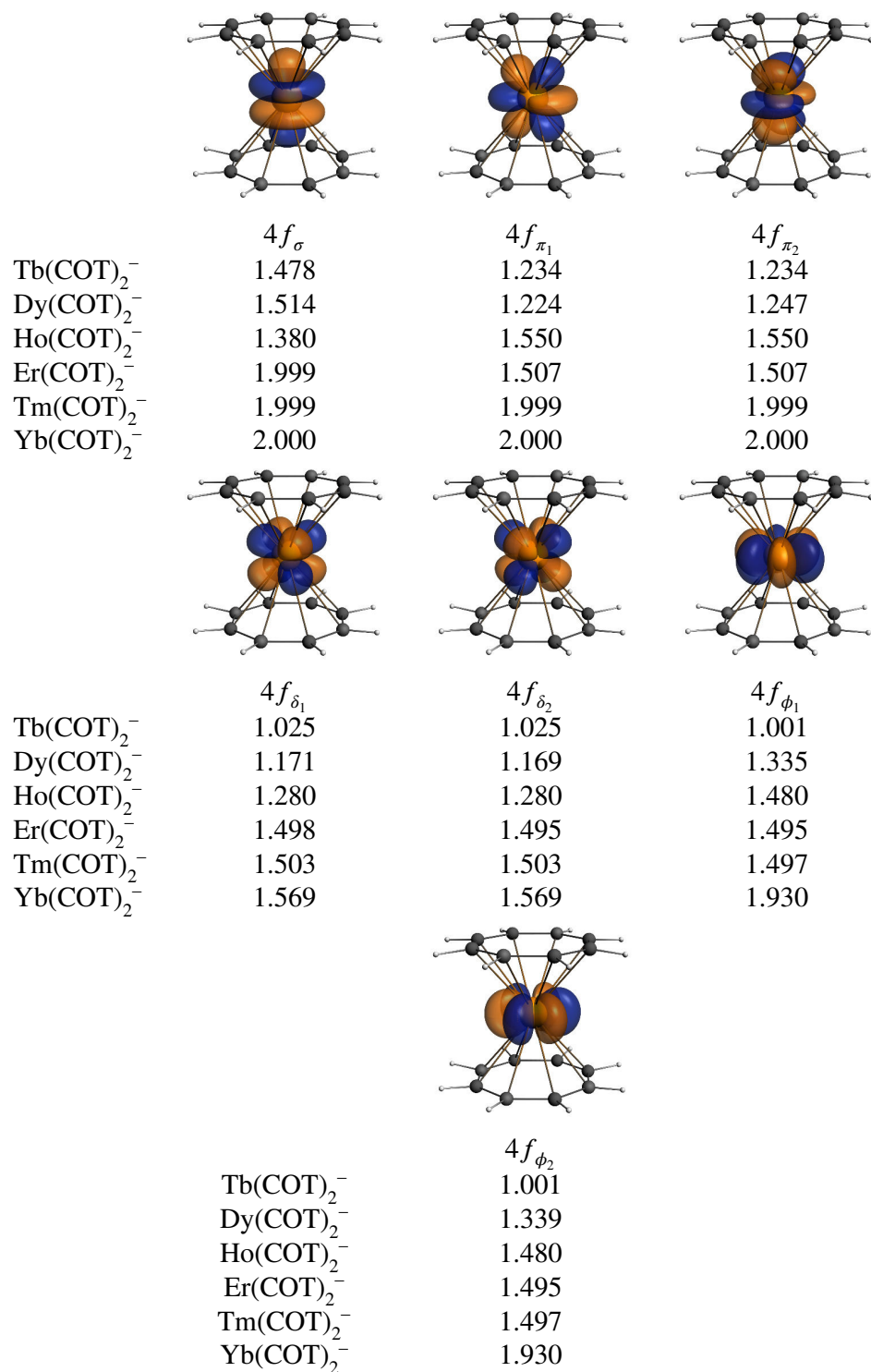


Figure S5: Selected NOs φ_p and occupation numbers n_p for Ln(COT)₂⁻. The figure shows the NOs from the SO calculation of Pr(COT)₂⁻. Isosurface values: ± 0.03 au.

Natural Spin Orbitals of the ground states in $\text{Ln}(\text{COT})_2^-$

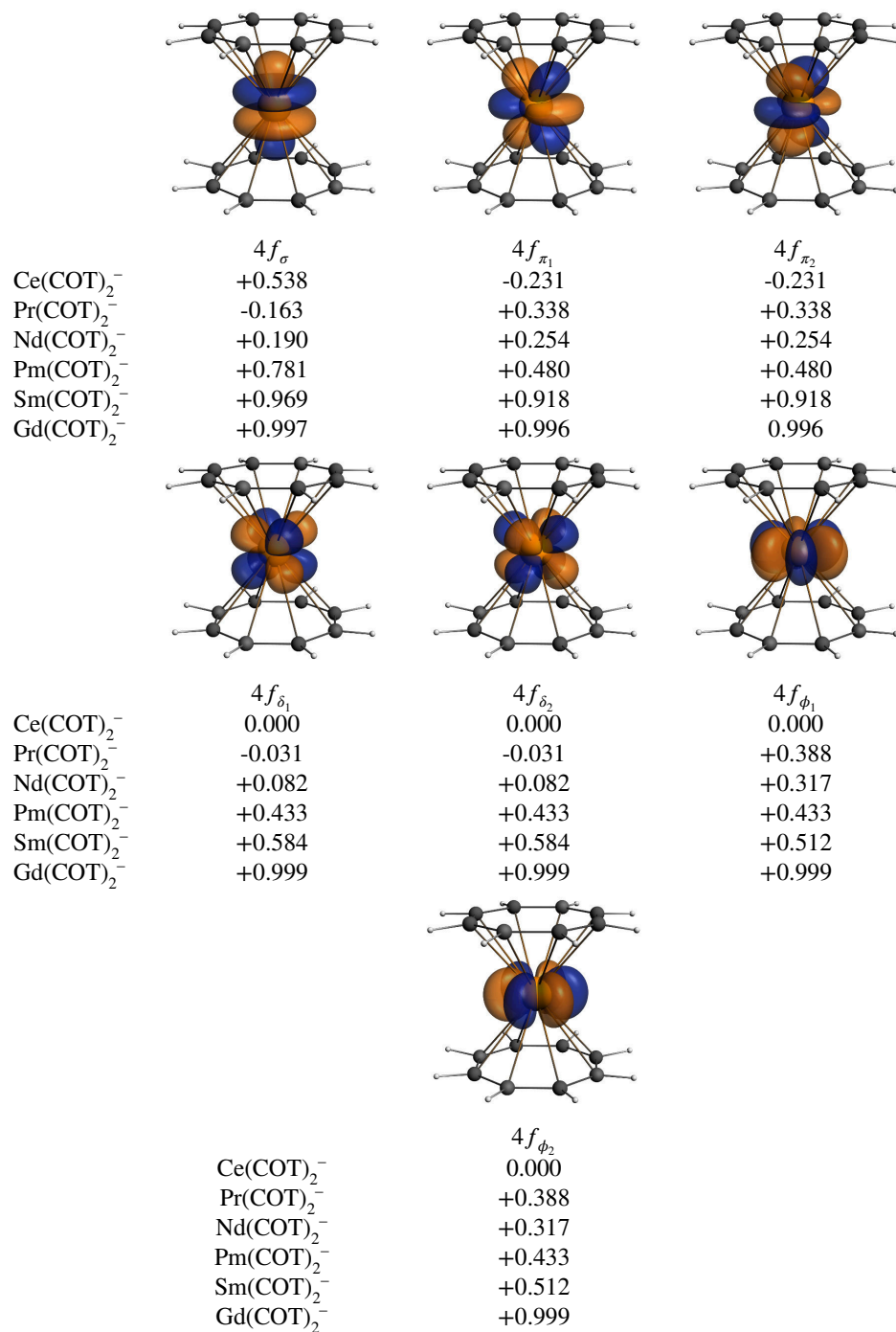


Figure S6: Selected NSOs φ_p^{\parallel} and contributions n_p^{\parallel} to $m_{\parallel}^S(\mathbf{r})$ for $\text{Ln}(\text{COT})_2^-$. The figure shows the NSOs from the SO calculation of $\text{Pr}(\text{COT})_2^-$. Isosurface values: ± 0.03 au.

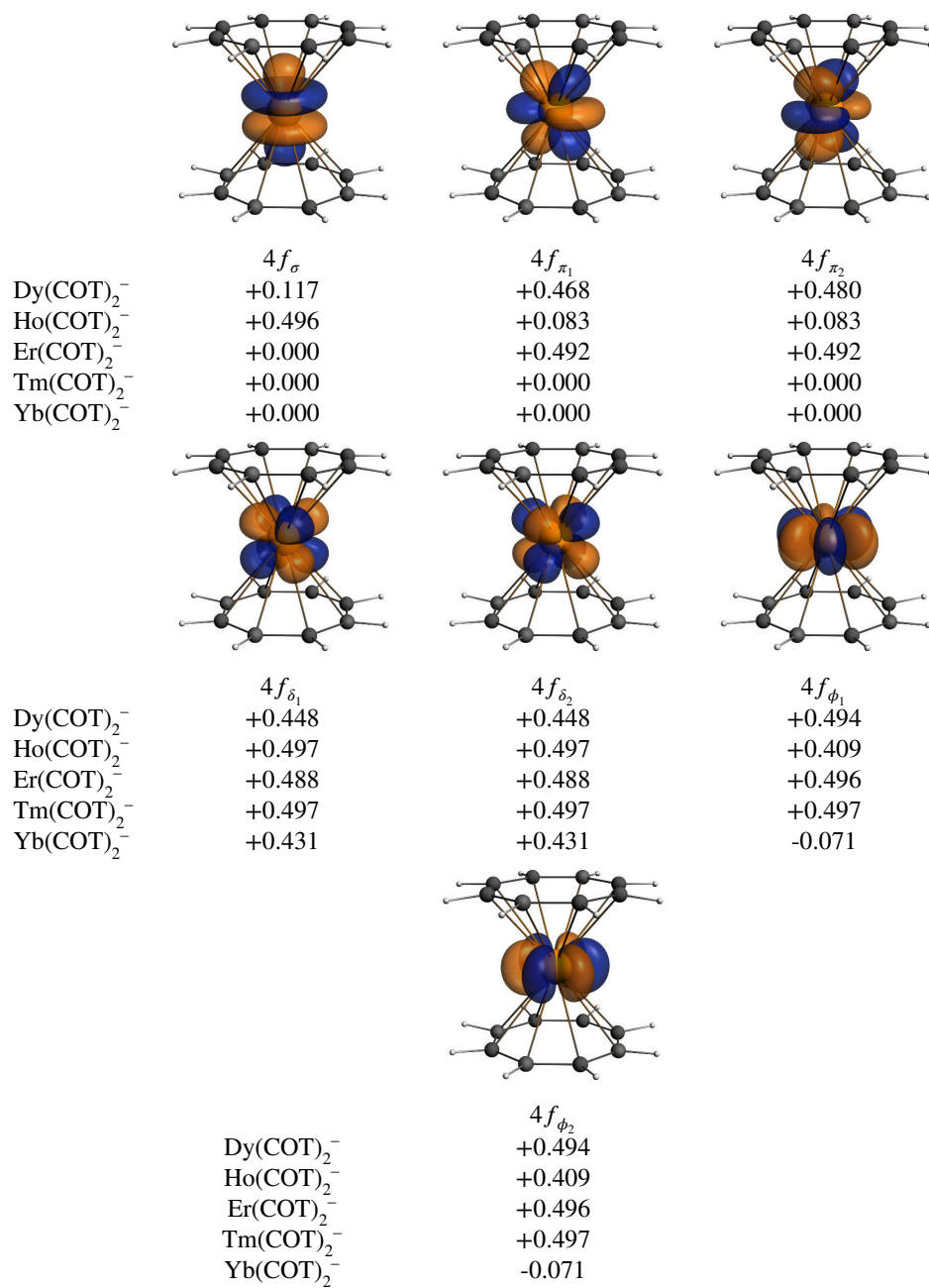


Figure S7: Selected NSOs φ_p^{\parallel} and contributions n_p^{\parallel} o $m_{\parallel}^S(\mathbf{r})$ for Ln(COT)₂⁻. The figure shows the NSOs from the SO calculation of Pr(COT)₂⁻. Isosurface values: ± 0.03 au.

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