

## Supplementary material

### Formation and nonlinear optical properties of Ag nanocrystals capped with the conjugated ligand carbazolyl styryl terpyridine

Wan Sun <sup>a</sup>, Yingzhong Zhu <sup>a</sup>, Anran Wang<sup>a</sup>, Lin Kong<sup>a</sup>, Shengli Li <sup>a\*</sup>,  
Jieying Wu <sup>a</sup>, Yupeng Tian <sup>a, b, c\*</sup>

<sup>a</sup> (Department of Chemistry, Anhui Province Key Laboratory of Functional Inorganic Materials, Anhui University, Hefei 230039, China)

<sup>b</sup> (State Key Laboratory of Crystal Materials, Shandong University, Jinan 250100, China)

<sup>c</sup> (State Key Laboratory of Coordination Chemistry, Nanjing University, Nanjing 210093, China)

\* Corresponding author. Fax: +86-551-65107342; Tel: +86-551-65108151

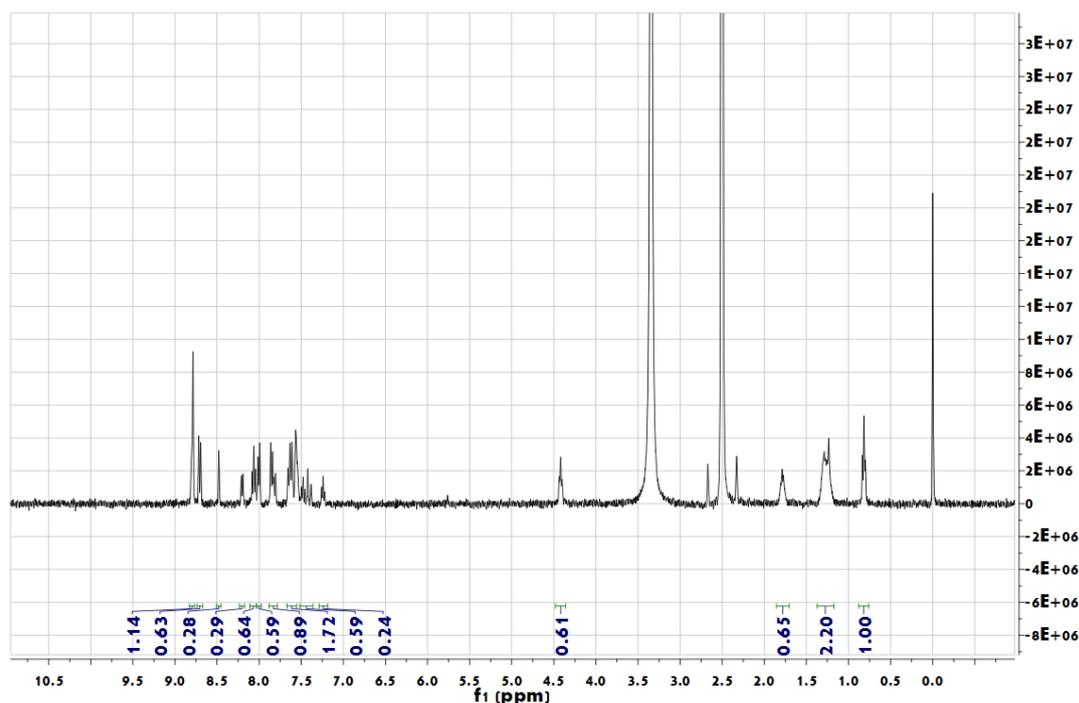
E-mail address: lsl1968@ahu.edu.cn; yptian@ahu.edu.cn

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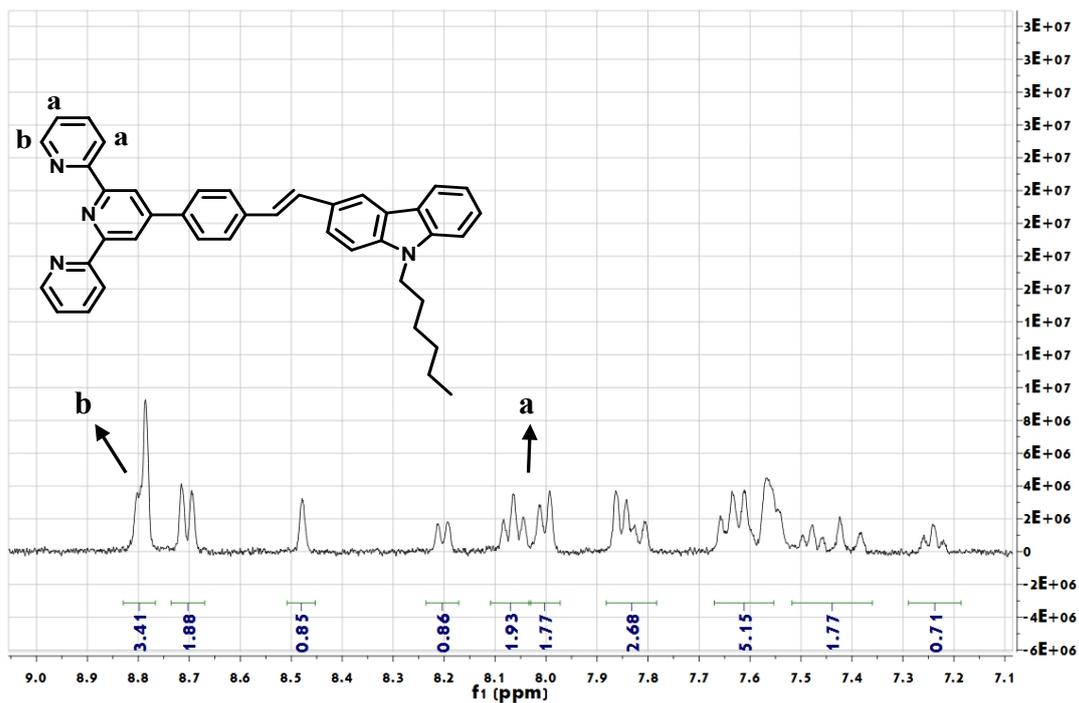
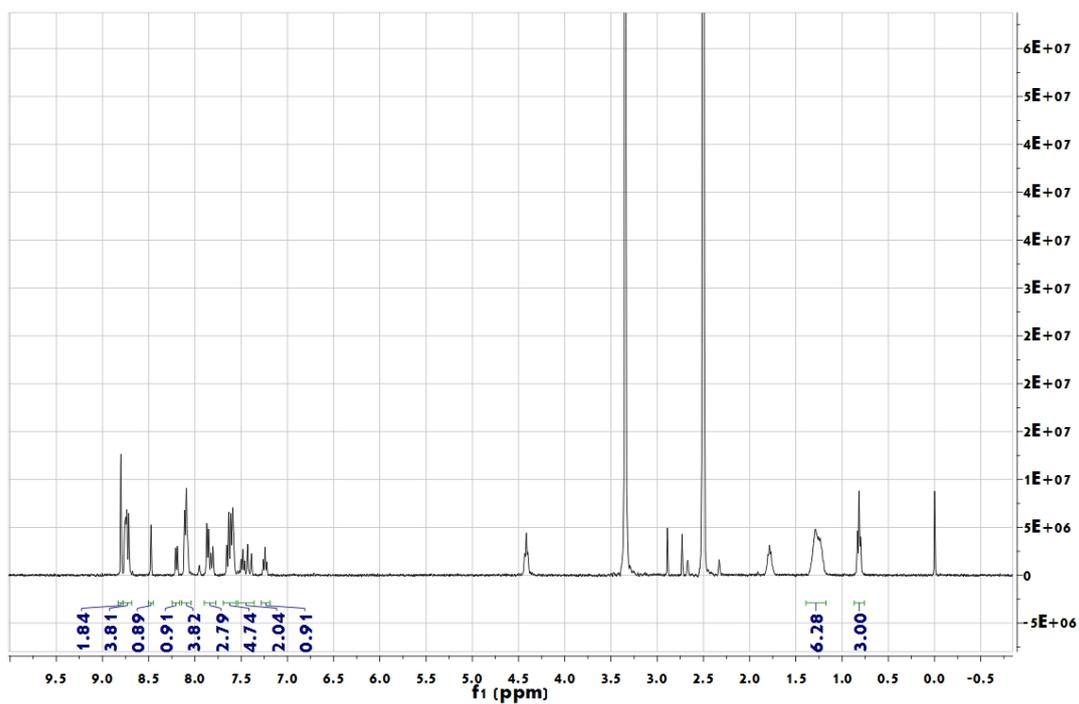


Fig. S1 Relevant peaks in the <sup>1</sup>H NMR spectra of pure L in d<sup>6</sup>-DMSO



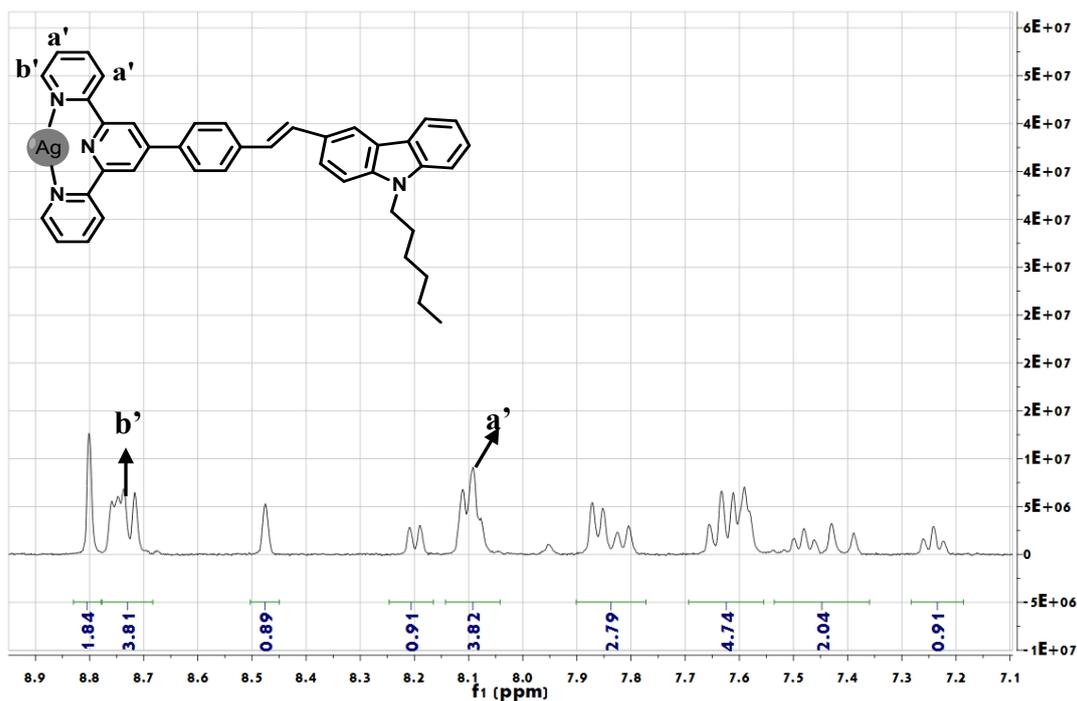


Fig. S2 Relevant peaks in the  $^1\text{H}$  NMR spectra of L-Ag NCs in  $\text{d}^6\text{-DMSO}$

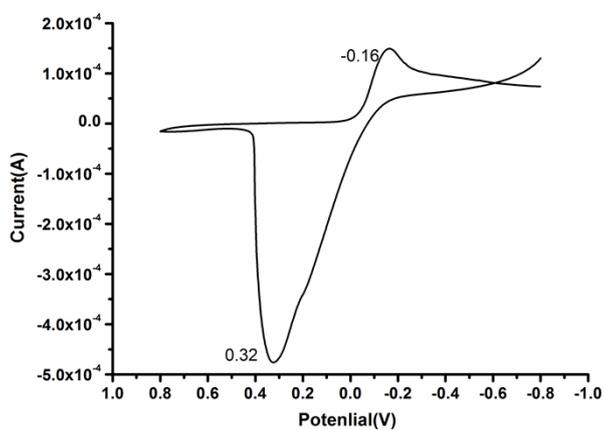


Fig. S3 The cyclic voltammetry curve of  $\text{AgNO}_3$

#### S4 The calculation method:

HOMO: Highest Occupied Molecular Orbital

LUMO: Lowest Unoccupied Molecular Orbital

$E_{\text{HOMO}} = -e(\text{Eonset} + 4.38) \text{ eV}$ , (Eonset: the first onset oxidation potential, Eonset=0.97V).

A HOMO-LUMO gap:  $E_{\text{opt}} = hc/\lambda$ , ( $\lambda = \lambda_{\text{edge}} = 589 \text{ nm}$ . The value of the intersection point of the tangent of the crest

of maximum wavelength and X-axis is  $\lambda_{\text{edge}}$ , as shown in Fig. 5,  $h$  is Planck constant  $4.136 \times 10^{-6}$ ,  $c$  is speed of light

$3.0 \times 10^8 \text{ m/s}$ ) so,  $E_{\text{opt}}$  is 2.11eV.)

$$E_{\text{LUMO}} = E_{\text{HOMO}} + E_{\text{opt}} = -5.35 + 2.11 = -3.24 \text{ eV}.$$

The oxidation potential of L is 0.97 V, which is higher than the electrode potential of  $\text{Ag}^+/\text{Ag}$  (-0.16 V) (shown in Fig. S3). This means that L can exist stably in  $\text{AgNO}_3$ -DMF solution.<sup>[25]</sup>

$E_{\text{HOMO}}$  was close to the Fermi level of Ag (-4.26 eV). The result indicates that the L combining with Ag nanoparticle is accompanied by electron redistribution, which probably leads to the optical property change.<sup>[26]</sup>

#### Open-aperture Z-scan:

The NLO absorption components were evaluated by Z-scan experiment under an open aperture configuration. The TPA coefficient  $\beta$  and TPA cross-sections ( $\sigma$ ) were determined by the OA Z-scan technique. The theoretical data were fitted using the following equations [33]:

$$T(z, s = 1) = \sum_{m=0}^{\infty} \frac{[-q_0(z)]^m}{(m+1)^{3/2}} \quad \text{for } |q_0| < 1 \quad (1)$$

$$q_0(z) = \frac{\beta I_0 L_{\text{eff}}}{1 + \chi^2} \quad (2)$$

$\beta$  is the nonlinear absorption (TPA) coefficient of the solution,  $I_0$  is the input intensity of laser beam at focus ( $z = 0$ ) divided by  $\pi\omega_0^2$ ,  $L_{\text{eff}} = [1 - \exp(-\alpha_0 L)]/\alpha_0$  is the effective length with  $\alpha_0$  the linear absorption coefficient and  $L$  the sample length.  $\chi = z/z_0$ ,  $z_0 = \pi\omega_0^2/\lambda$  is the diffraction length of the beam with  $\omega_0$  the spot size at focus,  $\lambda$  is the wavelength of the beam and  $z$  is the sample position. So the nonlinear TPA coefficient  $\beta$  (in units of  $\text{cm/GW}$ ) can be deduced. Furthermore, the  $\sigma$  could be determined by the following relationship [34]:

$$\sigma = \frac{h\gamma\beta}{N_A d} \times 10^{-3} \quad (3)$$

Here,  $h$  is the Planck constant,  $\gamma$  is the frequency of incident laser,  $\sigma$  is molecular TPA cross-section,  $N_A$  is the Avogadro number, and  $d$  is the concentration (in units of  $\text{mol}\cdot\text{L}^{-1}$ ). Based on equation(3), the molecular TPA cross-section  $\sigma$  can be calculated.

#### Close-aperture Z-scan:

For the closed aperture, the calculation of the nonlinear refractive index  $\gamma$  fitting can be done as in Equation (4).

$$\Delta T_{\text{p-v}} = 0.406(1-s)^{0.25} |\Delta\Phi_0| \quad \text{where } \Delta\Phi_0 = \omega(\Delta n)L = K \gamma I_0 L_{\text{eff}} \quad (4)$$

where  $\Delta T_{\text{p-v}}$  is the peak-valley transmittance difference from the closed-aperture scan. It can be seen that the difference between normalized transmittance values at valley and peak positions,  $\Delta T_{\text{p-v}}$ , was 0.73 for R and 1.49

for the nanohybrid.  $s$  is the fraction of the transmitted beam through the aperture (0.20 in our experiment).  $\Delta\Phi_0$  is the on-axis nonlinear phase shift and  $K$  is the wave vector ( $K=2\pi/\lambda$ ,  $\lambda$  was 790 nm for R and the nanohybrid).

The third-order nonlinear susceptibility ( $\chi^{(3)}$ ) was also determined through the closed-aperture Z-scan method. The value of the real part of the third-order nonlinear susceptibility,  $\text{Re}(\chi^{(3)})$ , can be calculated by the experimental measurements of  $g$  as in Equation (5).

$$\mathbf{R}_c\chi^{(3)}(\text{esu})=\mathbf{n}_0^2\epsilon_0c^2\gamma/\pi \quad (5)$$

where  $\epsilon_0$  is the vacuum permittivity,  $c$  is the velocity of light in a vacuum,  $n_0$  is the linear refractive index.

The value of the imaginary part of the third-order nonlinear susceptibility  $\text{Im}(\chi^{(3)})$  can be calculated from the value of  $b$  as given in Equation (6)

$$\mathbf{I}_m\chi^{(3)}=\mathbf{n}_0^2\epsilon_0c^2\lambda\beta/4\pi^2 \quad (6)$$

[33] Wan Sun, Pingping Sun, Shengli Li, *Dyes and Pigments* 2015, **115**, 110-119.