

Supporting Information

High performance, visible to mid-infrared photodetector based on graphene nanoribbons passivated by HfO₂

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1. Methods.

Sample Preparation: Highly-ordered pyrolytic graphite (HOPG) was purchased from SPI supplies. Monolayer graphene was fabricated by a 3M tape-based mechanical exfoliation on SiO₂/Si wafer with 285 nm thermal-oxidized SiO₂ layer. Heavily p-doped silicon substrate was employed as a bottom gate electrode. Monolayer graphene was identified by the optical microscope and Raman spectroscopy. The single Lorenz 2D peak ($\sim 2700\text{ cm}^{-1}$), $I_{2D}/I_G > 2$ and absence of D peak ($\sim 1350\text{ cm}^{-1}$) demonstrated the quality of the single layer graphene.

EBL Process: Monolayer graphene sheet with lateral sizes of $\sim 4\text{ }\mu\text{m}$ was contacted with Ti/Au (20/80 nm) metal electrodes fabricated by e-beam evaporation after photolithography. Consequently, the graphene channel was patterned into nanoribbons by EBL process. As shown in Figure 1(a)-(d), the electron beam lithography (EBL) can be briefly described as follows. First, hydrogen silsesquioxane (HSQ) was spin-coated on the whole wafer, and the thickness was $\sim 30\text{ nm}$ by using the spin speed of 5000 rpm. Second, the HSQ covered sample was exposed by electron beam with a designed nanoribbon pattern. Third, the exposed sample was developed by NaOH/NaCl salty solution, and the un-exposed HSQ was washed away. Next, the sample was put in the RIE chamber and etched by O₂ plasma, the graphene area with HSQ covering was retained while the other uncovered graphene parts will be etched away by the O₂ plasma. Last, HSQ resist was washed away by immersing in HF solution ($\sim 2\%$) for 30 seconds.

ALD Process: Hafnia coating was deposited by atomic layer deposition (ALD) (Cambridge Nanotech), the chamber was kept at 50°C before graphene samples were placed inside, and then increased to 270 °C at low vacuum level (10^{-3} mTorr). One single cycle of an ALD is composed of 1-second pulse of TEMA (tetraakis (ethylmethylamido) hafnium) and successive 1-second pulse of water vapor followed by nitrogen purge of 5 seconds. Different thicknesses of hafnia, 1.8 nm, 3.7 nm, 6.2 nm and 7.8 nm were controlled by using 20, 40, 60 and 80 cycles. The thickness and surface roughness were measured by AFM (Bruker

Dimension FastScan, contact mode). The quality of the ALD film was characterized by X-ray Photoelectron spectroscopy (XPS, ESCALAB 250 Xi).

Electrical and Photoresponse Measurements: Electrical characteristics were conducted by a semiconductor device analyser (Agilent B1500A). Photoresponsivity measurement was performed in a digital deep level transient spectroscopy (DLTS, BIORAD) combined with a four-probe stations. Three light sources (visible: 632 nm semiconductor laser; near infrared: 1470 nm fiber laser; mid-infrared: tunable quantum cascade laser (QCL) from 9.75 μm to 10.48 μm) were employed to illuminate the whole device. The lasers were aligned by optical lens and the power densities of the lasers were assumed to be uniform. All the measurements were conducted at room temperature.

2. Characterization results

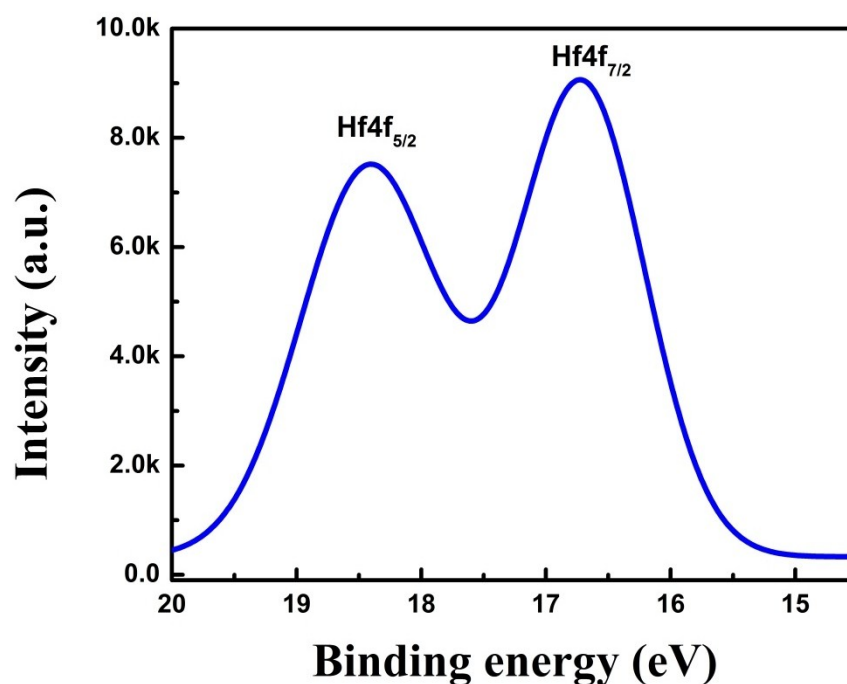


Figure S1. XPS spectra (Hf4f) of the ALD HfO₂ film. The intensity ratio of the Hf 4f 7/2 and Hf 4f 5/2 peaks are constrained to 4:3. The distance of the two peaks is $D = 1.71$ eV. The Hf 4f 7/2 binding energy for the metal is found at 16.6 eV and is fit with an asymmetric peak-shape. The Hf 4f 7/2 binding energy for the oxide is found ~ 18.3 eV.

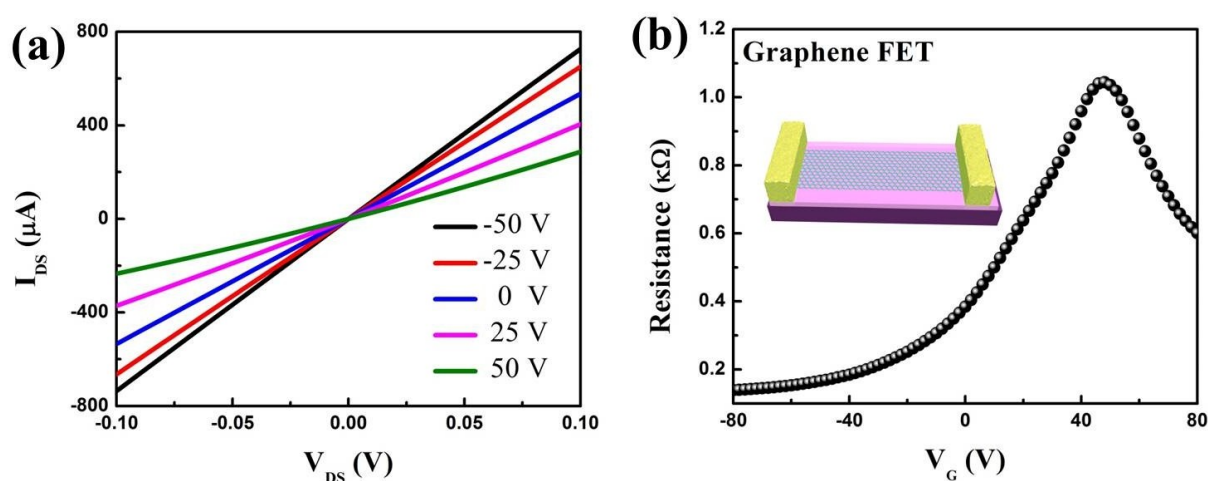


Figure S2. (a) I_{DS} - V_{DS} curves of graphene FET recorded at different gate voltages. (b)

Dependence of graphene FET on gate voltage recorded at $V_{DS}=10$ mV, indicating a semiconducting behavior curve.

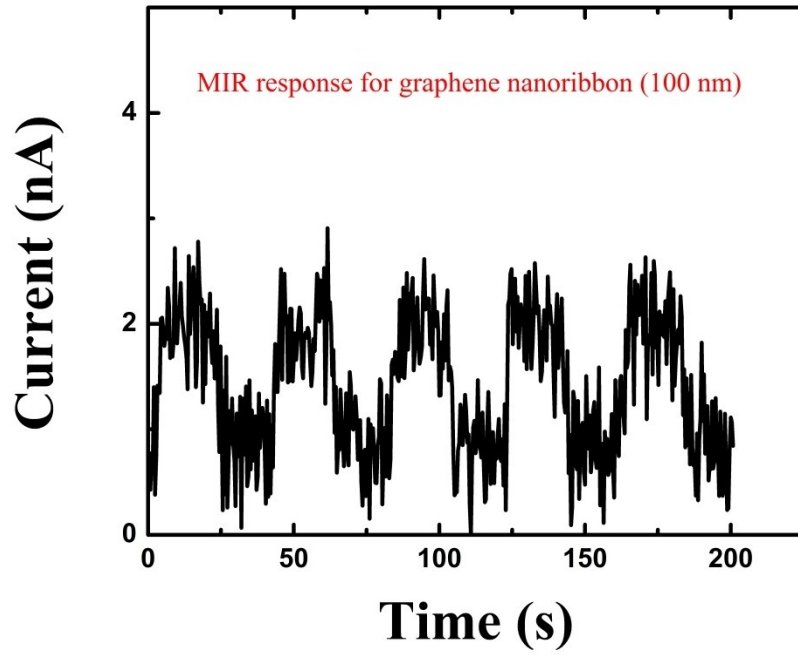


Figure S3. Photoresponse of graphene nanoribbon FET photodetector with a nanoribbon width of 100 nm under a 10 μm laser illumination at room temperature.

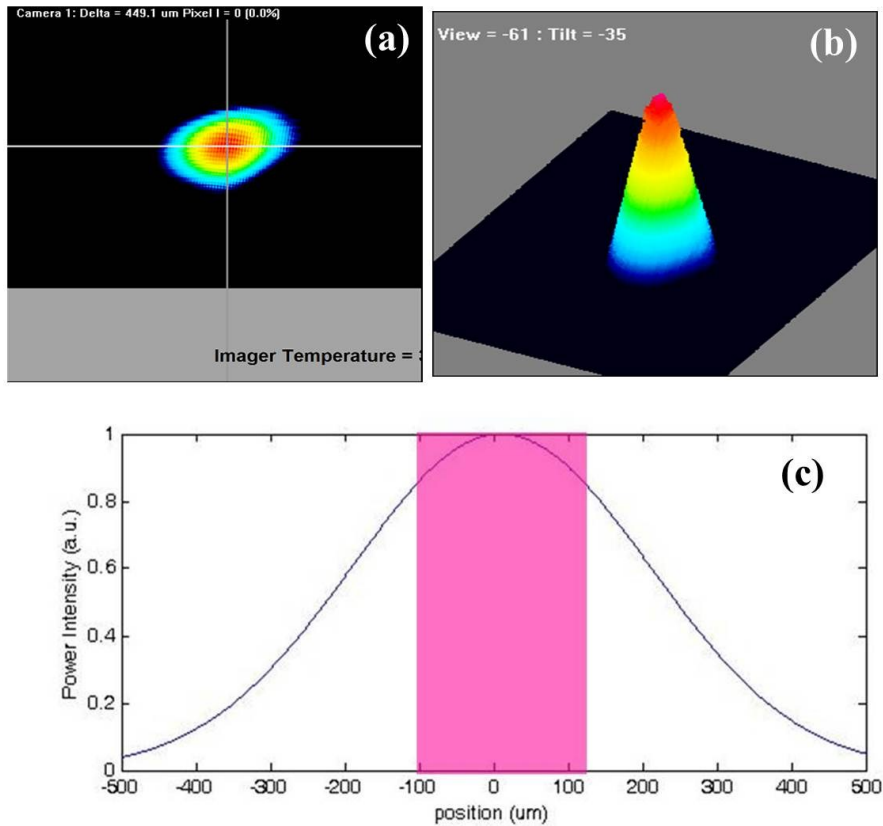


Figure S4. Beam profile of the mid-infrared laser used in our experiment.

3. Calculation of photoresponsivity, gain and detectivity

The photoresponsivity (R) are calculated as the ratio of photocurrent over the optical power on graphene:

$$R = I_{pc} / (P \times 2.3\%)$$

For example, the photon responsivity for this device operated in the visible can be calculated as $R = 300 \text{ nA} / (2.5 \text{ W/cm}^2 \times 4 \text{ } \mu\text{m} \times 16 \text{ } \mu\text{m} \times 10\%) = \sim 1.75 \text{ A/W}$. Using the same method, we can calculate the photoresponsivity of the near infrared (1.47 μm) and mid-infrared (10 μm) as 1.5 A/W and 0.18 A/W, respectively. The gain and detectivity are calculated as following:

$$G = \frac{I_{ph} / e}{\left(\frac{S_G}{S_L} \cdot P \cdot 2.3\%\right) / h\nu} \quad \left(\nu = \frac{c}{\lambda_{incident}}\right) \quad \text{and} \quad D = \frac{I_{ph}}{P \sqrt{2eI_{dark}}},$$

where I_{ph} and I_{dark} are the photocurrent dark current, S_G and S_L are the area of the graphene FET channel and light spot, P is the illumination power, ν and λ are the frequency and wavelength of the incident light, e is the electron charge, h is the Planck constant, and c is the speed of light.

4. Calculation of photoconductive gain

The numbers of the photogenerated electrons can be assumed as the consequence of excitation to the conduction band, relaxation from the conduction band, and recombination with holes in the valance band or defects. Thus, we have^[S1-S2]

$$\frac{dn_1}{dt} = \chi\beta - Rn_1 - \beta n_1 = 0 \quad (1)$$

$$\frac{dn_2}{dt} = \alpha n_1 - \frac{n_2}{\tau} = 0 \quad (2)$$

$$n_1 + n_2 = \frac{I}{e} \tau_t, \quad (3)$$

where e is the electron charge, n_1 and n_2 are the electron populations in the conduction band and the valance band, R is the recombination rate, I is the measured current, α is the capture rate of electrons, β is the electron generated rate, χ is the multipul excitation factor, τ and τ_t are the measured decay time and transient time. From equations (1) to (3), the capture rate can be expressed by

$$\alpha = \frac{\left(\frac{I}{e\chi\beta} \tau_t R - 1 \right)}{\left(\tau_t - \frac{I}{e\chi\beta} \tau \right)} \approx \frac{I}{e\chi\beta} \tau_t R / \tau,$$

and the quantum efficiency of carrier multiplication can be expressed by

$$\eta = \frac{\alpha n_1}{\beta} = \frac{\alpha}{\chi(\alpha + R)} \approx \frac{I}{e\chi\beta\tau}$$

References

- [S1] Zhang, Y., Liu, T., Meng, B., Li, X., Liang, G., Hu, X., Wang, Q. *J. Nat. Commun.* **2013**, *4*, 1811.
- [S2] Winzer, T.; Knorr, A.; Malic, E., Carrier Multiplication in Graphene. *Nano Lett.* **2010**, *10*, 4839.