

## **Extraordinary Optical Transmission in Nanopatterned Ultrathin Metal Film without Holes**

Akshit Peer<sup>1,3</sup> and Rana Biswas<sup>\*,1,2,3</sup>

<sup>1</sup>Department of Electrical and Computer Engineering, and <sup>2</sup>Department of Physics and  
Astronomy, Iowa State University, Ames, Iowa 50011, USA

<sup>3</sup>Ames Laboratory, Ames, Iowa 50011, USA

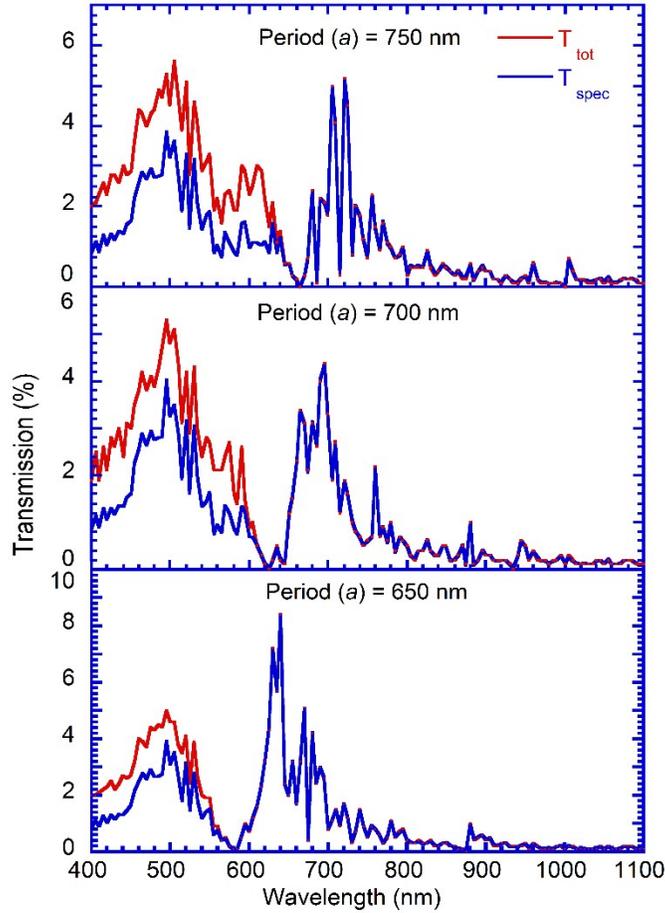
\*Corresponding Author: [biswasr@iastate.edu](mailto:biswasr@iastate.edu)

### **Supplementary Information**

- I. Period and  $R/a$  Variation
- II. Transverse Electric Field Profile
- III. Convergence with Number of Bloch Waves

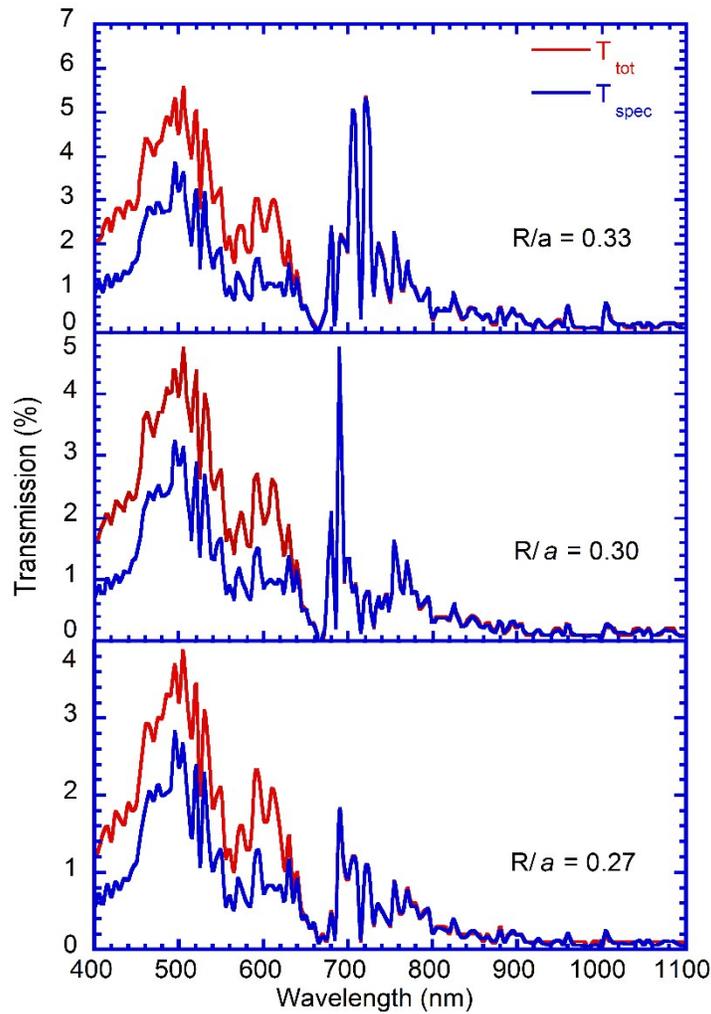
## I. PERIOD AND R/a VARIATION

**Period Variation:** We present the simulation results for different periods ( $a$ ) = 650 nm, 700 nm and 750 nm for constant  $R/a = 0.33$  in Fig. S1. It can be seen that the surface plasmon resonance peak of gold appears at same wavelength  $\lambda_{sp} \sim 504$  nm for all the three cases. The EOT peak appears at a wavelength slightly smaller than  $a$ . Notably, the position of the transmission peak scales with the period of the structure, that is, the peak transmission wavelength increases as  $a$  increases. The Wood's anomaly where  $T = 0$  is also seen at  $\lambda_W = \sqrt{3}a/2 \sim 580$  nm, 610 nm and 650 nm for  $a = 650$  nm, 700 nm and 750 nm, respectively. Corresponding to  $a = 650$  nm, 700 nm and 750 nm, the EOT peak appears at wavelength  $\lambda_1 \sim 625$ -640 nm, 665-695 nm and 707-720 nm, respectively. For our analysis, we choose  $a = 750$  nm since it is closest to the experimental structure. The precise shape of the transmission curves does not scale with  $a$  since the dielectric function of gold is wavelength-dependent.



**Figure S1.** Simulated total transmission and specular transmission spectra for the nanocup structure with different periods ( $a$ ) = 650 nm, 700 nm, and 750 nm. The value of  $R/a = 0.33$  for all the cases.

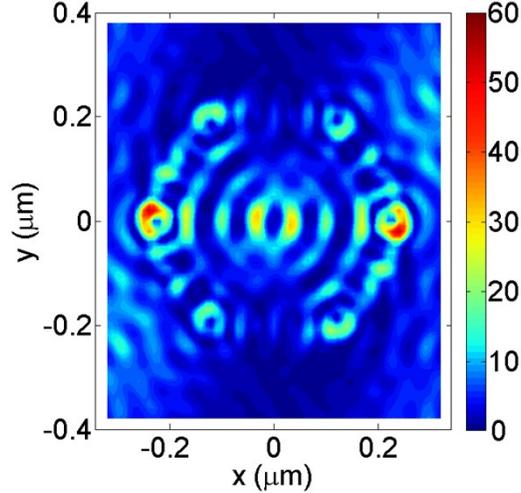
**R/a Variation:** We present the simulation results for different  $R/a = 0.27, 0.30,$  and  $0.33$  for constant pitch  $a = 750$  nm in Fig. S2. Similar to Fig. S1, the surface plasmon resonance peak of gold appears at same wavelength  $\lambda_{sp} \sim 504$  nm for all three cases. Since the period is constant for all the cases, the EOT peak ( $\sim 700$  nm) and the Wood's anomaly point ( $\sim 650$  nm) appear at nearly the same wavelength. Notably, when  $R/a$  is decreased the EOT peak becomes narrower and the transmission intensity also reduces due to the decrease in aperture of the nanocup. We present the results in the manuscript for  $R/a = 0.33$ , which corresponds to the experimental structure.



**Figure S2.** Simulated total and specular transmission spectra for the nanocup structure with different  $R/a = 0.27, 0.30,$  and  $0.33$ . The value of period ( $a$ ) = 750 nm for all the cases.

## II. TRANSVERSE ELECTRIC FIELD PROFILE

The electric field intensity  $|\mathbf{E}|^2$  at the EOT wavelength  $\lambda_1 \sim 700$  nm in the  $x$ - $y$  plane (Fig. S3) just below the surface of the nanocup shows an intensity distribution complementary to that in the  $xz$ -plane (Fig. 6a). The results are shown for incident field with  $x$  polarization. The  $|\mathbf{E}|^2$  has peaks in small regions at the edges of the circular aperture where  $|\mathbf{E}|^2$  is enhanced by  $\sim 50$ - $60$  in both the  $xy$  and  $xz$  cross-sections. Dipolar charge distributions are formed at the edges of the nanocup. This is typical of surface plasmon propagation along the  $x$ -axis (ref. 26). There are a pair of additional sub-peaks on either side of the center of the nanocup surface (Fig. S3 and Fig. 6a), where the  $|\mathbf{E}|^2$  is enhanced by  $\sim 50$ . There is a node of the electric field at the center of the nanocup surface as seen in both Fig. S3 and Fig. 6a. Since the nanocup geometry can be envisaged as a combination of many holes of decreasing radii, we used five layers of holes with decreasing radii ( $R, 0.8R, \dots, 0.2R$ ) to model the nanocup field intensity. We can see that the electric field plot in Fig. S3 is actually the superposition of well-known field intensity from a single hole (ref. 26) with field maxima lying along  $x$ -axis for incident field with  $x$  polarization.

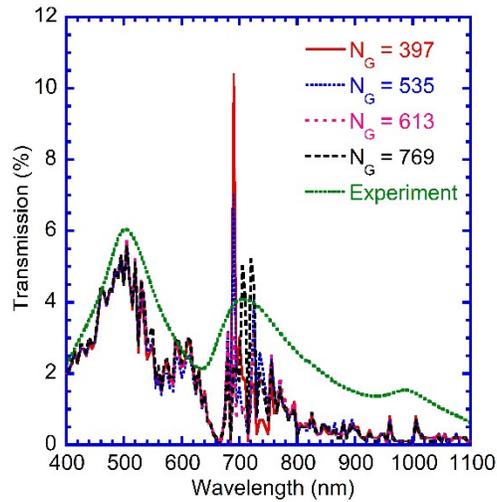


**Figure S3.** Electric field intensity at  $\lambda_1 \sim 700$  nm showing the enhanced field in the  $xy$  plane just below the surface of nanocup.

## III. CONVERGENCE WITH NUMBER OF BLOCH WAVES

We have investigated the convergence of the transmission with increased number of Bloch waves ( $N_G$ ) per polarization as summarized in Fig. S4, using a 5 nm thick gold layer at the bottom of the nanocup corresponding to the geometry used for simulation in Fig. 4a. The scattering matrix has dimension of  $2 N_G \times 2 N_G$ . The  $N_G$  values correspond to closed shells of

reciprocal lattice  $\mathbf{G}$  vectors. The convergence of short wavelengths ( $\lambda < 660$  nm = Woods anomaly wavelength,  $\lambda_w$ ) or at longer wavelengths  $\lambda > 800$  nm is very good for all the chosen  $N_G$  values. In the EOT region 650-850 nm (Fig. S4), the transmission is more sensitive to the number of Bloch waves, and we find  $N_G = 535, 769$  are converged. Both  $N_G = 535$  and  $N_G = 769$  show two EOT peaks with a narrower feature around 700 nm and a broader feature (consisting of multiple peaks at 720-780 nm). These two features combine to produce the broader experimental peak centered around 704 nm.



**Figure S4.** Simulated transmission spectra with  $N_G = 397, 535, 613, 769$  compared to the experimental transmission.