

## Competitive cation binding to Phosphatidylinositol-4,5-bisphosphate domains revealed by X-ray fluorescence

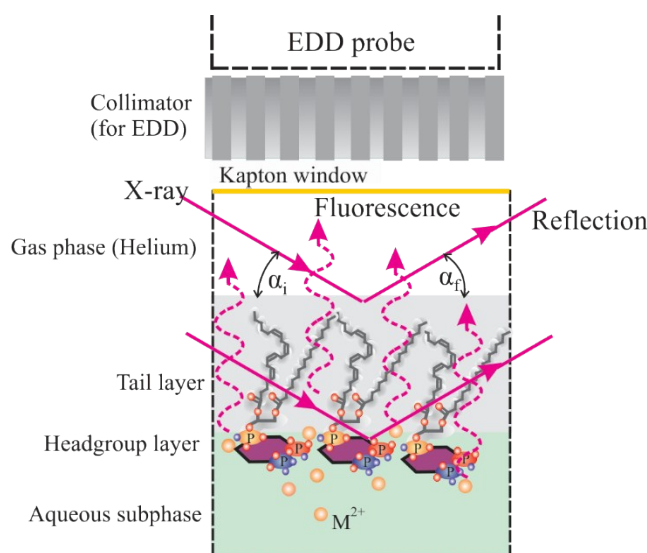
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### 1 Detailed Methods and Data Analysis

For completeness of this Section, parts of it also appear in the main text.

#### 1.1 Experimental set-up and monolayer details

X-ray reflectivity (XR) and near-total-reflection fluorescence (XNTRF) measurements were conducted on the liquid surface spectrometer (LSS) at beamline 9ID-C, Advanced Photon Source (APS), Argonne National Laboratory. Figure S1 provides a sketch of the experimental setup. The monolayer is formed by drop-wise deposit (at molecular area 85-100 Å<sup>2</sup>/molecule) of PIP<sub>2</sub> stock solution onto the aqueous surface of the subphase solutions that is contained in a Langmuir trough and maintained at constant temperature of 20 °C. The Langmuir trough is encapsulated in an enclosure with thin Kapton windows and purged with water-saturated helium gas to minimize background scattering from air and potential X-ray radiation damage to the samples. An oxygen sensor (S101, Qubit System Inc.) in the trough is used to monitor oxygen levels in the enclosure. The monolayer is compressed by a Teflon barrier to a desired surface pressure of 30 ± 2 mN/m. Surface pressure is monitored by a microbalance made of a Wilhelmy microbalance and filter paper plate. The X-ray measurements are started after the oxygen-helium exchange reaches equilibrium, where the oxygen content inside the enclosure is reduced by a factor of at least 100 with respect to that of the ambient atmosphere. The monolayer is maintained at  $\pi \approx 30$  mN/m during the course of the experiments. Radiation damage is monitored by systematically translating the trough laterally to probe fresh surfaces and examine measurement reproducibility. All reflectivity measurements were found to be reproducible.



**Figure S1: Setup.** Illustration of the X-ray fluorescence and reflectivity setup for the PIP<sub>2</sub> monolayer over a subphase containing divalent ions M<sup>2+</sup>, where M represents Ca or Mg. The box enclosed by the dashed lines is a side view of the interfaces of a vapor phase (water-saturated helium), monolayer and aqueous bulk. The X-ray beam impinges on the surface at an incident angle  $\alpha_i$  and reflects at an exit angle  $\alpha_f$  and  $\alpha_i = \alpha_f$  in a specular reflectivity mode. The probe of the EDD subtends the center of the sample surface with a collimator in front of it to define the EDD footprint that is smaller than the X-ray beam footprint and only accepts

X-ray photons of the fluorescence signals and background scattering (emanated from bulk water and helium) restricted to the surface normal direction with 1 degree acceptance angle.

## 1.2 X-ray experiments

The highly monochromated and collimated X-ray beam ( $E = 8.0$  keV; wavelength  $\lambda = 1.5497$  Å) passes through an ionization chamber (to monitor incident beam intensity) and variable beam attenuator (multilayers of nickel foil, for attenuation of the primary incident beam), and is steered onto the aqueous surface by a second monochromator to a desired incident angle  $\alpha_i$  with respect to the surface. The specular reflection is collected by a Bicron scintillation detector at an X-ray exit angle  $\alpha_f = \alpha_i$ . The angular-dependence of both XRR and XNTRF are expressed as functions of  $Q_z$ , the z-component (surface normal) of the scattering (wave vector transfer) vector  $Q$ .  $Q_z$  is related to the incident angle  $\alpha_i$  via  $Q_z = (4\pi/\lambda) \sin\alpha_i$ , and  $Q_c$  represents the  $Q_z$  that corresponds to the critical angle for total reflection,  $\alpha_c (= 0.154^\circ$  at  $E = 8\text{keV}$ ).

The fluorescence signals are collected at a series of  $\alpha_i$  near the critical angle  $\alpha_c$  by a Vortex energy dispersive detector (EDD, silicon-drift, Vortex-90EX) with a collimator in the front end of the EDD probe, which accepts the X-ray fluorescence of emitted photons in the direction along the surface normal at  $\sim 1^\circ$  angular resolution.

Fluorescence signals from pure water subphase are measured as background to be subtracted from those of the relevant samples. The EDD spectra also contain the primary beam signals at  $E = 8.0$  keV mainly from Thomson elastic scattering of bulk water and Compton inelastic scattering (from bulk water and water-saturated helium, at a lower energy shifted from  $E = 8.0$  keV by  $\leq 0.1$  keV). Both fluorescence and reflectivity data are normalized to incident beam intensity (more details can be found elsewhere <sup>1, 2</sup>).

## 1.3 X-ray reflectivity data analysis

X-ray reflectivity data,  $R(Q_z)$ , are normalized to the calculated Fresnel reflectivity ( $R_F$ ), for an ideally smooth, flat air-water interface <sup>3,4</sup>. The  $R/R_F$  data for a surface monolayer on an aqueous surface can be accounted for in terms of a simplistic, *two-box* (i.e., two-layer) structural model, where one *box* contains the hydrophilic head group, and the other contains the hydrocarbon tail <sup>1,3,5</sup>. Each *box* is characterized by a uniform electron density ( $\rho$ ) and vertical height ( $l$ ) relative to the interface. Given  $(\rho_H, l_H)$  and  $(\rho_T, l_T)$ , in which subscripts "H" and "T" represent the headgroup and the hydrocarbon tail respectively, a step-like, discrete ED profile across the interfaces,  $\rho_0(z)$ , is constructed and its corresponding reflectivity  $R_0(Q_z)$  is calculated by the Parratt recursion method <sup>1,3</sup>. The realistic ED profile,  $\rho(z)$  is obtained by smearing  $\rho_0(z)$  with an interfacial roughness,  $\sigma$ , mainly arising from capillary waves <sup>4</sup>, which lowers the  $R_0(Q_z)$  by a "Debye-Waller-like" factor,  $\exp(-Q_z^2 \sigma^2)$  as follows: <sup>1,3,4</sup>

$$R/R_F = (R_0/R_F) e^{-Q_z^2 \sigma^2} \quad (S1)$$

The structural refinements are carried out through a least squares optimization method <sup>1</sup>.

### 1.3.1 Calculation of molecular area from reflectivity data

The molecular area,  $A_{\text{mol}}$ , of  $\text{PIP}_2$  can be calculated from these reflectivity data by assuming that a  $\text{PIP}_2$  molecule has an acyl-chain length corresponding to the stearic acid tail at the sn-1 position of natural  $\text{PIP}_2$ . The maximum acyl chain length of stearic acid,  $l_{\text{Tmax}}$ , can be estimated as  $\sim 22$  Å since each carbon (17 in stearic acid, excluding the carbonyl) contributes  $\sim 1.27$  Å <sup>6,7</sup>. The molecular area of the lipid  $A_{\text{mol}}$ , can be determined based on this maximum ( $l_{\text{Tmax}}$ ) and the actual acyl-chain length,  $l_T$ , using  $2A_0 l_{\text{Tmax}} / l_T$ , where  $A_0$  is the minimum area available to a single, untilted alkyl chain ( $\sim 20$  Å<sup>2</sup>, according to measurements with arachidic acid) <sup>6,7</sup>. Here we assume that the arachidonoyl chain of natural  $\text{PIP}_2$  contributes an equal area, as the stearic acid chain, to the overall area of each  $\text{PIP}_2$ . Although the organization of arachidonoyl chains is largely unknown, this is a reasonable assumption and leads to a lower bound of the molecular area of  $\text{PIP}_2$  ( $A_{\text{mol}}$  in Table S1). These values are consistent with those obtained from isotherms ( $A_{\text{mol}2}$ ).

### 1.3.2 Calculation of # of divalent cations bound to $\text{PIP}_2$

The reflectivity data allows for an estimation of the number of divalent cations ( $\text{M}^{2+}$ ) bound to the headgroup of  $\text{PIP}_2$ , based on a space-filling model<sup>1, 8, 9</sup>. The model assumes the headgroup volume is the sum of volumes of the  $\text{K}^+$  ions,  $\text{M}^{2+}$  ions, water molecules, and the  $\text{PIP}_2$  head group. The conservation of electron numbers and constrains on the total volume contained in the two *boxes* (one associated with the acyl chains and the other, more complex, includes the headgroup, water molecules and ions) can be expressed as follows

$$A_{\text{mol}} I_H = V_{\text{H,dry}} + n_{\text{H}_2\text{O}} V_{\text{H}_2\text{O}} + n_{\text{K}^+} V_{\text{K}^+} + n_{\text{M}^{2+}} V_{\text{M}^{2+}}, \quad (\text{S2})$$

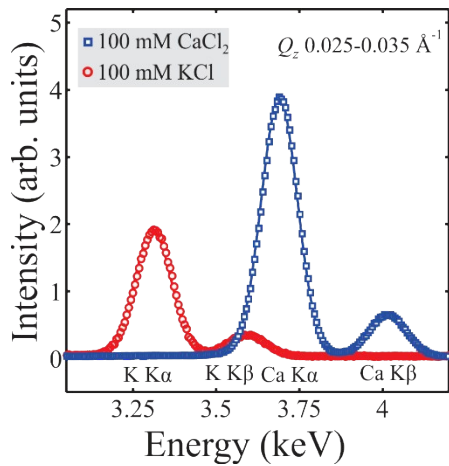
$$A_{\text{mol}} \rho_{\text{H}} I_H = N_{\text{H,dry}}^e + N_{\text{H}_2\text{O}}^e n_{\text{H}_2\text{O}} + N_{\text{K}^+}^e n_{\text{K}^+} + n_{\text{M}^{2+}} N_{\text{M}^{2+}}^e, \quad (\text{S3})$$

Assuming the number of bound  $\text{K}^+$  per  $\text{PIP}_2$  is unchanged prior to and after  $\text{M}^{2+}$  binding, we obtain the  $n_{\text{M}^{2+}}$  the number of  $\text{M}^{2+}$  per  $\text{PIP}_2$  molecule (see Table 1 and 2 in the main manuscript)<sup>9</sup>. These binding numbers serve as lower limits of actual binding number of  $\text{Ca}^{2+}$  and  $\text{Mg}^{2+}$  per  $\text{PIP}_2$ . In this case we also assume that the divalent ions are dehydrated when they bind to  $\text{PIP}_2$ , however, Slochow et al. suggested that  $\text{Mg}^{2+}$  may remain hydrated when it binds to  $\text{PIP}_2$ <sup>10</sup>. If we assume that a hydration shell of six molecules remains around  $\text{Mg}^{2+}$ , this reduces the calculated number of bound  $\text{Mg}^{2+}$  ions to  $\sim 0.6$  ions/ $\text{PIP}_2$ .

#### 1.4 X-ray fluorescence data analysis

Surface fluorescence studies have been performed in the past<sup>11, 12</sup> and analytic routines to quantify the surface enrichment of ions have been established<sup>8, 9</sup> and are briefly discussed below. At  $\alpha_i < \alpha_c$ , the X-rays penetrate into the liquid by a very shallow depth ( $< 10$  nm) in the direction of the surface normal<sup>1, 3</sup>. At  $\alpha_i > \alpha_c$ , the X-rays penetrate into the bulk by a few micrometers (within the X-ray energy of 8-16 keV typically used for liquid surface scattering). The combination of shallow X-ray penetration depth and surface intensity enhancement arising from evanescence when  $\alpha_i < \alpha_c$  provides the surface sensitivity of the XNTRF method. In this study, the fluorescence measurements were conducted over the range  $0.5\alpha_c < \alpha_i < 1.5\alpha_c$ , corresponding to  $Q_z = 0.01\text{-}0.035 \text{ \AA}^{-1}$ , with a spacing of  $0.001 \text{ \AA}^{-1}$  between consecutive  $Q_z$  points. The fluorescence data collected at  $Q_z < Q_c$  and  $Q_z > Q_c$  are referred to as surface fluorescence and bulk fluorescence signal, respectively.

To calibrate the spectral intensity to known quantities of fluorescing ions, we compare fluorescence data from two bulk solutions of  $\text{CaCl}_2$  (pre-acidified with a small amount of HCl to prevent potential formation and precipitation of  $\text{Ca(OH)}_2$  at this high  $\text{CaCl}_2$  concentration) and KCl both prepared at 100 mM with a bare air/water interface. Figure S2 shows the individual bulk fluorescence spectra for  $\text{Ca}^{2+}$  and  $\text{K}^+$ . Each spectrum within the energy range shown in Figure S2 is characterized by two emission lines,  $\text{K}_\alpha$  and  $\text{K}_\beta$ , each of which is further profile-fit with a Gaussian-like function centered on the emission line energy. The finite spread of the emission lines is a measure of the EDD energy resolution at each emission line. The bulk fluorescence spectra shown in Figure S2 provide the “standard spectral profile” for each element, which characterizes (1) the relative intensity ratio of  $\text{K}_\alpha$  to  $\text{K}_\beta$  and (2) the relative line(s) intensity ratio of  $\text{Ca}^{2+}$  to  $\text{K}^+$ , given the same bulk concentrations. As 100 mM KCl is present in the relevant samples and the  $\text{K}^+$   $\text{K}_\beta$  line partially overlaps with the  $\text{Ca}^{2+}$   $\text{K}_\alpha$  line, the standard spectral profiles serve to discriminate the  $\text{K}^+$  and  $\text{Ca}^{2+}$  contribution when both are present in the spectra.



**Figure S2: Bulk Fluorescence Spectra.** The bulk fluorescence spectra for a pure KCl and a  $\text{CaCl}_2$  solution as indicated. Both solutions were prepared at 100 mM.  $\text{CaCl}_2$  solutions were acidified to prevent  $\text{Ca}^{2+}$  precipitation as  $\text{Ca(OH)}_2$  (due to the high concentration of  $\text{CaCl}_2$ ). The spectra were integrated over  $Q_z = 0.025\text{-}0.035 \text{ \AA}^{-1}$  to demonstrate the relative intensity of emission lines of  $\text{Ca}^{2+}$  relative to  $\text{K}^+$  detected by the EDD.

The  $\text{Mg}^{2+}$  ion emission lines are  $K_\alpha = 1.25$  keV and  $K_\beta = 1.07$  keV. These low energy emission lines are outside the sensitivity range of the EDD. At each  $Q_z$ , the fluorescence intensity of a specific emission line is integrated over its spread in the spectrum and expressed as a function of  $\sin \alpha_i$  (i.e., proportional to  $Q_z$ ). For a solution of fluorescing ions with a bare surface the fluorescence intensity,  $I_b(\alpha_i)$ , is proportional to its bulk concentration  $n_b$  and the volume illuminated by X-ray, and can be expressed as follows:

$$I_b(\alpha_i) = C |t_F(\alpha_i)|^2 D(\alpha_i) n_b \quad (S4)$$

where  $C$  is an element-specific scale factor that accounts for elemental fluorescence yield and detector efficiency,  $t_F(\alpha_i)$  is the X-ray Fresnel transmission coefficient at a sharp, flat vapor/liquid interface,  $D(\alpha_i)$  is the X-ray penetration depth normal to the surface, and  $n_b$  is the bulk concentration of the ion<sup>1,3</sup>. For a thin layer of surface excess ions, their fluorescence contribution,  $I_s(\alpha_i)$ , can be expressed as follows:

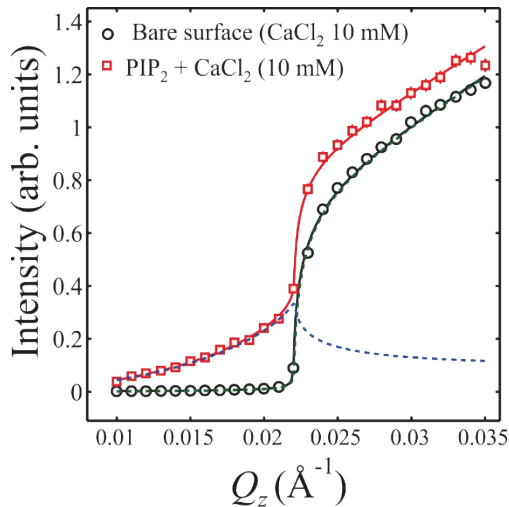
$$I_s(\alpha_i) = C |t_F(\alpha_i)|^2 n_s e^{-|z_{\text{ion}}|/D(\alpha_i)} \quad (S5)$$

where  $z_{\text{ion}}$  is the z-coordinate of the ion-enriched layer with  $z = 0$  at the monolayer/vapor interface.

For a solution with relatively low cation concentration ( $< 0.01\text{M}$ ) with a bare surface (i.e., no monolayer), the fluorescence comes almost exclusively from the bulk (see the black data points in Figure S3). In this case the ions are randomly distributed throughout the solution, and the very few ions that are observed in the 5 nm surface region are practically beyond the detection limit of the EDD. Thus, for the situation where  $Q_z < Q_c$  there is very little to almost no fluorescence signal. For  $Q_z > Q_c$ , the ions in the bulk solution are excited and we observe detectable fluorescence intensity. For a cation enriched surface bound to  $\text{PIP}_2$ , the fluorescence signal is observed both below and above  $Q_c$  (see Figure S3 in red –signal containing both the surface and bulk contribution). The surface density of cations for a surface covered with a  $\text{PIP}_2$  monolayer is given by:

$$n_s = \frac{I_s(\alpha_i)}{I_b(\alpha_i)} D(\alpha_i) n_b e^{z_{\text{ion}}/D(\alpha_i)} \quad (S6)$$

For known mean molecular area of the lipids, the number of ions per lipid can be calculated by simply multiplying the ion density by the area per lipid. This molecular area is obtained from the isotherms recorded during the preparation of the  $\text{PIP}_2$  monolayer. Briefly,  $\text{PIP}_2$  was spread at a molecular area between  $85\text{-}100 \text{ \AA}^2/\text{molecule}$  (i.e. at negligible surface pressure) on each of the subphases examined and compressed to a pressure of  $30 \pm 2 \text{ mN/m}$ , monitored by a Wilhelmy balance. The molecular areas obtained from the isotherms are the same (within experimental error) to those extracted from the analysis of the X-ray reflectivity data, see tables S1 and S2 below.



**Figure S3: Fluorescence  $Q_z$  Dependence.** Fluorescence intensities of  $\text{Ca}^{2+}$  emission line(s) as a function of  $Q_z$  for a solution of  $\text{CaCl}_2$  at concentration of 10 mM covered with a  $\text{PIP}_2$  monolayer (red) and an equivalent 10 mM  $\text{CaCl}_2$  solution with a bare surface (black). Each data point represents the intensity integrated exclusively over the  $\text{Ca K}_\alpha$  emission line in the spectrum. The solid lines are best-fit profiles calculated in terms of Eqn. 3 for the bulk contribution and the sum of Eqn. 3 and Eqn. 4 for the solution covered with the  $\text{PIP}_2$  monolayer. The dashed line is solely due to the surface contribution as expressed in Eqn. 3.

## References

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