

## **Supporting Information Available:**

### **Interlayer Coupling and Optoelectronic Properties of Ultrathin Two-Dimensional Heterostructures based on Graphene, MoS<sub>2</sub> and WS<sub>2</sub>**

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## Optical properties of individual MoS<sub>2</sub>, WS<sub>2</sub> and graphene

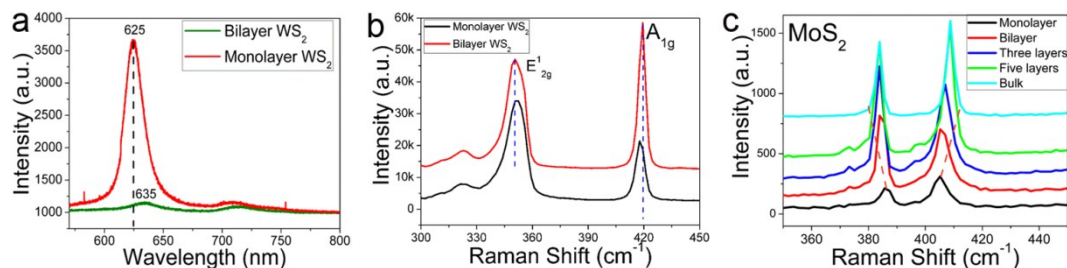


Fig. S1 (a) Photoluminescence (PL) and (b) Raman spectra of monolayer and bilayer WS<sub>2</sub>. (c) Raman spectra of MoS<sub>2</sub> with increasing number of layers (from monolayer to bulk).

The PL spectra of WS<sub>2</sub> layers are shown in Fig. S1a, where the significantly increased PL intensity and blue-shift of PL peak from bilayer to monolayer indicate the indirect-to-direct bandgap transition. Furthermore, the phonon frequency of atomically thin MS<sub>2</sub> (M = W or Mo) layers also exhibit unique thickness dependence, that is the Raman active mode A<sub>1g</sub> (out-of-plane displacement of S atoms) stiffens and the E<sub>2g</sub> (in-plane displacement of M and S atoms) softens with increasing number of layers as seen in Fig. S1b for WS<sub>2</sub> and Fig. S1c for MoS<sub>2</sub> which is consistent with previous reports<sup>S1,S2</sup>.

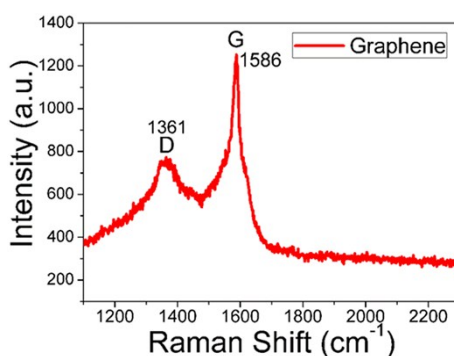


Fig. S2 Raman spectra of graphene, showing the typical Raman modes of graphene: G peak at 1586 cm<sup>-1</sup> and D peak at 1361 cm<sup>-1</sup>.

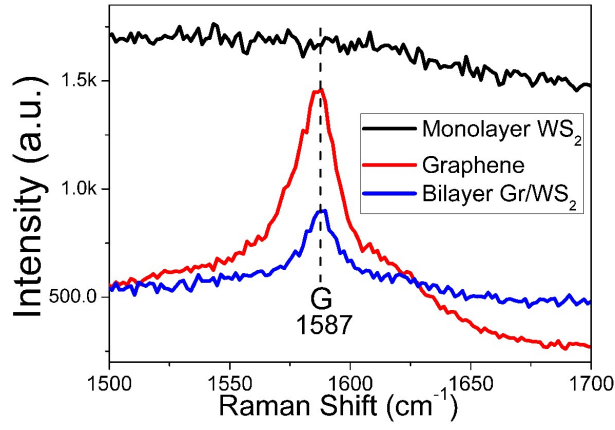


Fig. S3 Raman spectra of monolayer WS<sub>2</sub>, graphene, and bilayer Gr/WS<sub>2</sub> from 1500 to 1700 cm<sup>-1</sup>.

Though the the Raman modes  $A_{1g}$  and  $E_{2g}^1$  of WS<sub>2</sub> in this heterostructure much stiffer compared to that in isolated WS<sub>2</sub> monolayer due to strong interlayer coupling between graphene and WS<sub>2</sub>, the G peak of graphene is unchanged suggesting the strong in-plane  $sp^2$  bond in graphene is not much affected by the interlayer interaction.

## Electrical and photoresponse properties of individual MoS<sub>2</sub>, WS<sub>2</sub>

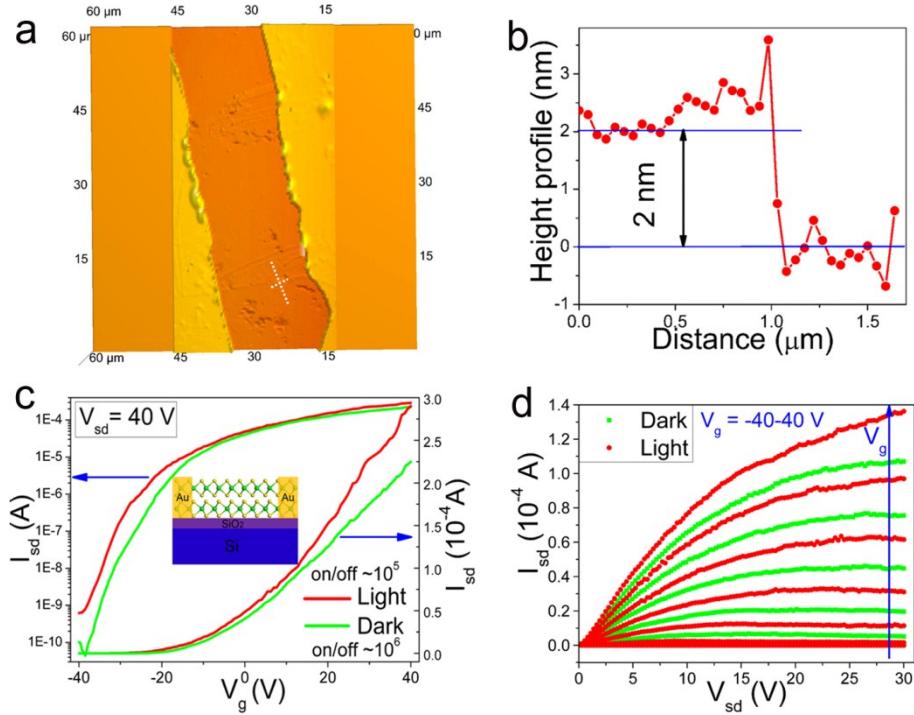


Fig. S4 AFM images of atomically thin MoS<sub>2</sub> transistors (a) with thickness of 2 nm from the cross-sectional plot (b), corresponding to 3-4 layers of MoS<sub>2</sub>. (c) Transfer and (d) output characteristic of ultrathin MoS<sub>2</sub> under dark and light illumination (173  $\mu\text{W}/\text{cm}^2$ ), the inset of (c) corresponds to the schematic diagram of transistor based on few-layer MoS<sub>2</sub>.

Transistors made of ultrathin MoS<sub>2</sub> and WS<sub>2</sub> layers are fabricated. The thickness of MoS<sub>2</sub> is 2 nm from atomic force microscope (AFM) images (Fig. S4a and S4b) corresponding to 3-4 layers of MoS<sub>2</sub>, and the schematic diagram of few-layer MoS<sub>2</sub> based transistor is shown in the inset of Fig. S4c. From the measured transfer and output characteristics shown in Fig. S4c and S4d, the few-layer MoS<sub>2</sub> transistor exhibits n-type behavior with the on/off ratio as high as  $5 \times 10^6$ . The electron mobility ( $\mu$ ) can be obtained from the equation  $\mu = \frac{\partial I_{sd}}{\partial V_g} \left( \frac{L}{WC_i V_{sd}} \right)$ , where L is the channel length, W is the channel width, and  $C_i$  is the gate capacitance between the channel and the silicon back gate per unit area, which can be given by equation  $C_i = \epsilon_o \epsilon_r / d$ ,  $\epsilon_o$  ( $8.85 \times 10^{-12}$  F/m) is vacuum dielectric constant, and  $\epsilon_r$  (3.9) and d (300 nm) are dielectric constant and thickness of SiO<sub>2</sub>, respectively. The calculated electron mobility is 10.3 cm<sup>2</sup>/Vs, which is comparable and even higher than previous studied few- or monolayer MoS<sub>2</sub><sup>S3-S5</sup>. Under visible light illumination, the source-drain current  $I_{sd}$  is significantly increased at the same gate voltage  $V_g$  (Fig. S4c) or source-drain voltage  $V_{sd}$  (Fig. S4d), indicating the few-layer MoS<sub>2</sub> is sensitive to visible light. We also notice that the photocurrent is more obvious and the photosensitive on/off ratio (defined as  $I_{\text{light}}/I_{\text{dark}}$ ) is larger under negative  $V_g$  compared to that under positive  $V_g$ . At Off-state (negative  $V_g$ ), the carriers in MoS<sub>2</sub> layers are consumed so the photo-

generated electron and hole carriers are dominating, leading to high photosensitivity (high on/off ratio). Whereas, at On-state (positive  $V_g$ ), the photo-generated carriers are not prominent because of the existence of high concentration of gate-induced carriers. As a result, the field-effect on/off ratio under light ( $5 \times 10^5$ ) is decreased by one order of magnitude compared to that under dark.

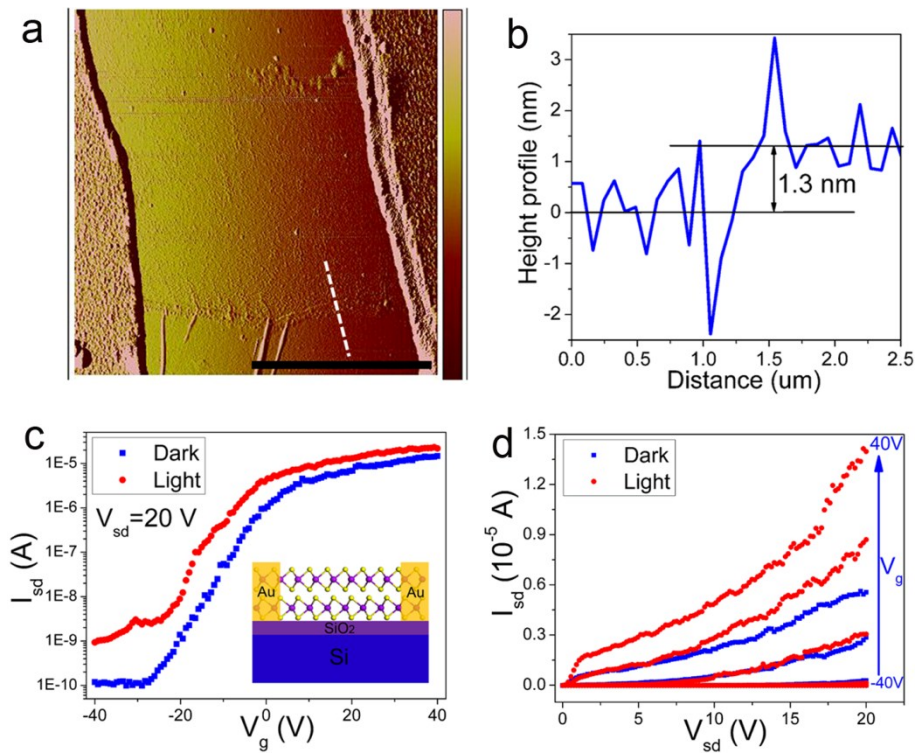


Fig. S5 AFM images of atomically thin  $WS_2$  transistors (a) with thickness of 1.3 nm from the cross-sectional plot (b), corresponding to 2 layers of  $WS_2$ . The scale bar in (a) is 5  $\mu m$ . (c) Transfer and (d) output characteristic of ultrathin  $WS_2$  under dark and light illumination (173  $\mu W/cm^2$ ), the inset of (c) is the schematic diagram of few-layer  $WS_2$  transistor.

The thickness of  $WS_2$  is 1.3 nm corresponding to 2 layers of  $WS_2$  from AFM images (Fig. S5a and S5b), and the inset of Fig. S5c shows the schematic diagram of the device. The transfer and output characteristics of few-layer  $WS_2$  transistors are

shown in Fig. S5c and S5d, respectively. The bilayer WS<sub>2</sub> also exhibits an n-type behavior with the on/off ratio as high as 10<sup>5</sup> under dark. The calculated electron mobility is 1.86 cm<sup>2</sup>/Vs. Similar to MoS<sub>2</sub>, WS<sub>2</sub> can also sensitively respond to the visible light and the photocurrent also become more obvious at Off-state, indicating the gate modulation of photocurrent.

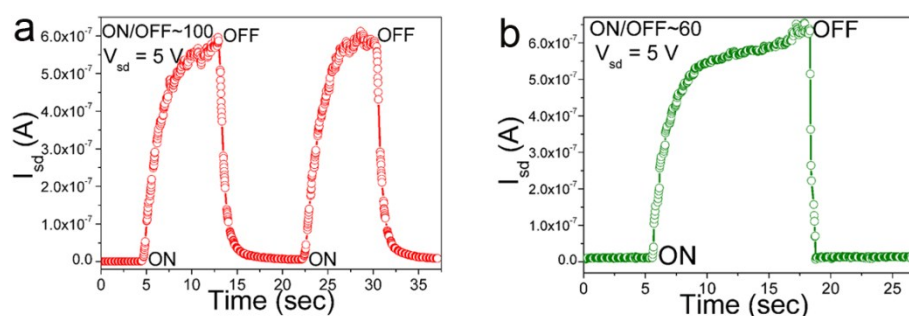


Fig. S6 Time dependences of source drain current  $I_{sd}$  of (a) MoS<sub>2</sub> and (b) WS<sub>2</sub> transistors during the visible light (1.36 mW/cm<sup>2</sup>) switching on/off at  $V_{sd} = 5$  V without gate voltage.

The dynamic response of the transistors based on ultrathin MoS<sub>2</sub> and WS<sub>2</sub> are also investigated during the light switching on/off as shown in Fig. S6a and S6b, respectively. Under light illumination, photocurrent can be generated resulting in the “On” state of the devices, and the devices can work between “On” and “Off” states fast and reversibly during light turned on/off with photosensitive on/off ratio of about 100 times for MoS<sub>2</sub> and 60 times for WS<sub>2</sub>. The response time is about 7 s and 5s for MoS<sub>2</sub> and WS<sub>2</sub>, respectively.

### P-type behavior of graphene

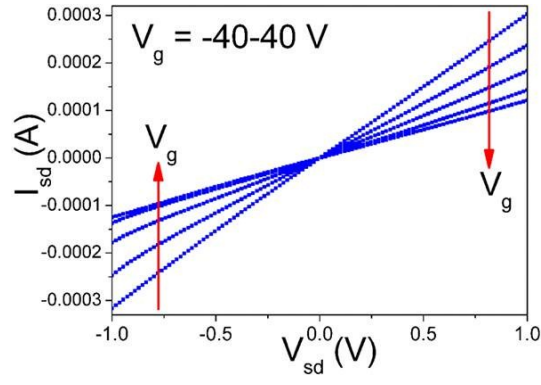


Fig. S7 Output characteristics of graphene, indicating the p-type behavior of graphene.

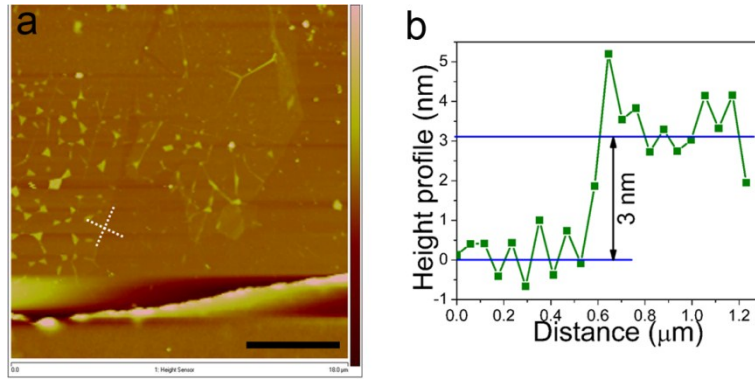


Fig. S8 AFM images of Gr/WS<sub>2</sub> heterostructure (a), and the cross-sectional plot (b) determines the thickness of the heterostructures to be 3 nm. The scale bar in (a) is 5 μm.

### Bipolar behavior of Gr/WS<sub>2</sub> heterostructure

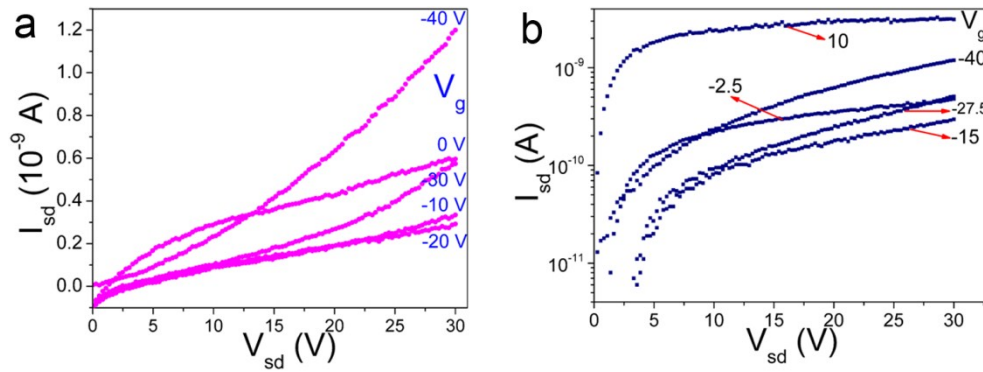


Fig. S9 Output characteristics of Gr/WS<sub>2</sub> heterostructure on linear scale (a) and log scale (b), exhibiting the typical bipolar behavior.



## Novel field-effect and photoresponsive properties of MoS<sub>2</sub>/WS<sub>2</sub> heterostructures

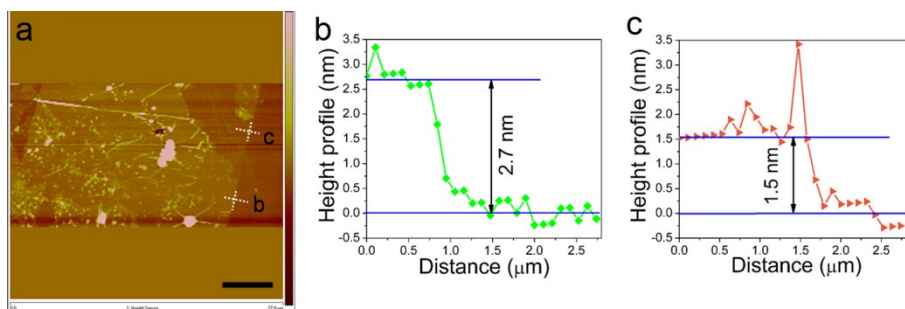


Fig. S10 AFM images of MoS<sub>2</sub>/WS<sub>2</sub> heterostructure (a). The total thickness of the heterostructure is 2.7 nm (b), and the thickness of bottom WS<sub>2</sub> layers is 1.5 nm (c), corresponding to 2-3 layers of WS<sub>2</sub> and 2 layers of MoS<sub>2</sub> in the heterostructure. The scale bar in (a) is 5 μm.

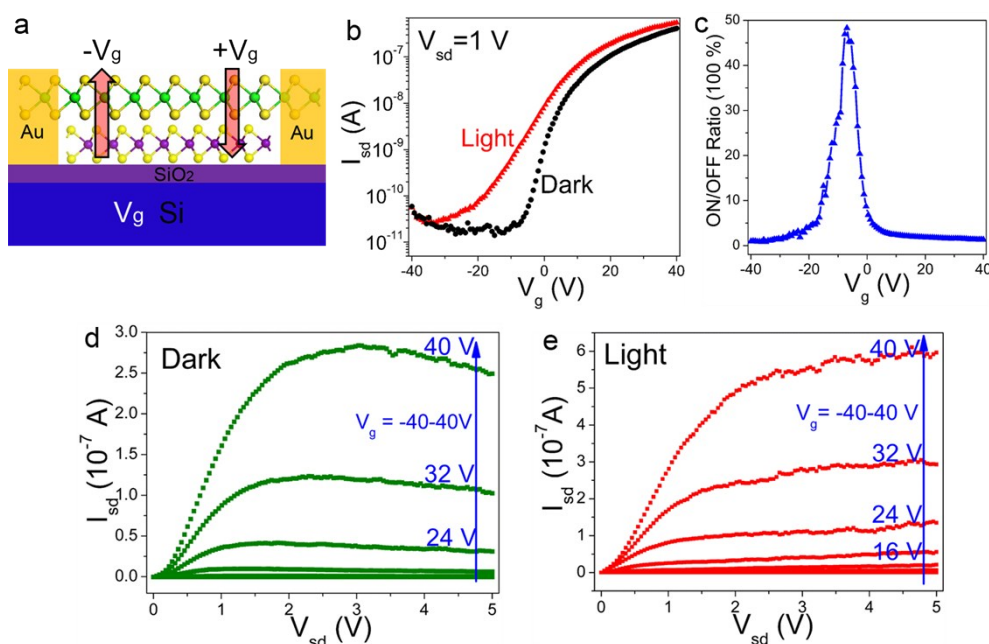


Fig. S11 (a) Schematic diagram of the transistors based on MoS<sub>2</sub>/WS<sub>2</sub> heterostructure with both Au electrodes located on top of MoS<sub>2</sub>. (b) Transfer characteristics of the device under dark and light illumination (173 μW/cm<sup>2</sup>). (c) Photocurrent on/off ratio of the device defined as  $I_{photo}/I_{dark}$  under different gate voltage, indicating the gate-tunable photocurrent. Output characteristics of



MoS<sub>2</sub>/WS<sub>2</sub> heterostructures under dark (d) and light illumination (e), showing the carrier leakage from top MoS<sub>2</sub> into bottom WS<sub>2</sub> under dark induced by the positive  $V_g$  and the photocurrent from WS<sub>2</sub> to MoS<sub>2</sub> can complement this  $V_g$ -induced electrons leakage.

There are lots of heterostructure-induced effects in the MoS<sub>2</sub>/WS<sub>2</sub> heterostructure based transistors. In the fabrication process, WS<sub>2</sub> with a thickness of about 1.5 nm (~2-3 layers) is covered by 1.2 nm MoS<sub>2</sub> (~2 layers), and the source-drain electrodes are located on top of heterostructure as shown in the AFM images (Fig. S10). A schematic diagram of the device can be seen in Fig. S11a. The total electron density in the heterostructure is determined by the applied  $V_g$ , but the current density would dominantly come from the electrons in the top MoS<sub>2</sub> layer. The charge transfer between the two layers can be tuned by  $V_g$ , and this induces an interesting electrical transport property. By applying the positive  $V_g$ , the gate-induced electrons in MoS<sub>2</sub> will transfer into the bottom WS<sub>2</sub> acting as trap and leading to the reduced On-state current. Hence the on/off ratio of heterostructures ( $>10^4$ ) is decreased compared with that in individual MoS<sub>2</sub> layers (as shown in the transfer curves, Fig. S11b). However, the threshold voltage ( $V_T$ ) of the heterostructures is increased to -7 V, indicating that the heterostructures facilitate the carrier depletion with the presence of WS<sub>2</sub> trap. From the output characteristics (Fig. S11d), the source drain current  $I_{sd}$  under dark reaches maximum at first, and then drops down under positive  $V_g$  with increasing  $V_{sd}$ . Because the positive  $V_g$  induces electrons transfer from MoS<sub>2</sub> to WS<sub>2</sub>. It is interesting to note that the  $I_{sd}$  can reach saturation under light illumination (Fig. S11e). This is benefit from the generation of photocurrent from WS<sub>2</sub> to MoS<sub>2</sub> which complements

the  $V_g$ -induced electrons leakage. On the contrary, the negative  $V_g$  induces the electrons transfer from  $WS_2$  into  $MoS_2$ , resulting in the “quasi-bipolar” transport behavior. Under high negative  $V_g$ , some electrons in  $WS_2$  layers are repelled into the  $MoS_2$  layers, leading to the slightly increased  $I_{sd}$  with increasing negative  $V_g$ . Thus the “p-type” behavior is obtained (Actually, the conducting carriers are still electrons).

Furthermore, the photocurrent can also be modulated by  $V_g$ . The photocurrent is non-significant under high (negative or positive)  $V_g$ , and the photosensitive on/off ratio can reach the maximum at  $V_T$  (Fig. S11c). When  $V_g > V_T$ , large amounts of carriers are induced by increased  $V_g$  leading to high dark current and large total electrons density, together with the unobvious photo-excited carriers and low photosensitive on/off ratio. When  $V_g < V_T$ , the device is turned off, so dark current is low and slightly increased due to the electrons transfer from bottom  $WS_2$  as discussed above. But the photocurrent under light illumination still drops down with decreased  $V_g$ . Thus, the photosensitive on/off ratio is also decreased with increased negative  $V_g$ .

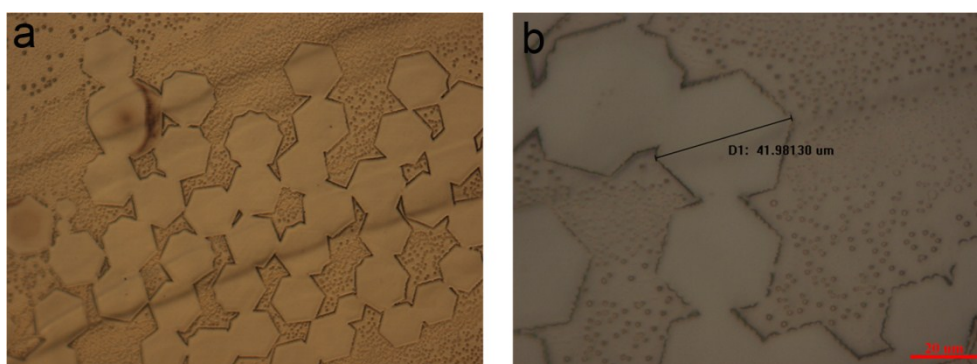


Fig. S12 Optical microscope images of hexagonal single-crystalline graphene with large-size of about 40  $\mu m$  synthesized by liquid copper CVD method.

## References

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