

Electronic Supplementary Information to accompany

**Extreme population inversion in the fragments  
formed by UV photoinduced S–H bond fission in  
2-thiophenethiol**

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## Vibrational Wavenumber Calculations for 2-thiophenethiyl Radical

Table S.1: Harmonic and anharmonic normal mode wavenumbers calculated at the B3LYP/6-311G(*d,p*) level for the ground state of the 2-thiophenethiyl radical. Modes are numbered in Herzberg notation.

Mode	Harmonic / $\text{cm}^{-1}$	Anharmonic / $\text{cm}^{-1}$	Symmetry
$\nu_1$	3234	3088	A'
$\nu_2$	3221	3078	A'
$\nu_3$	3203	3061	A'
$\nu_4$	1467	1432	A'
$\nu_5$	1427	1398	A'
$\nu_6$	1314	1288	A'
$\nu_7$	1198	1173	A'
$\nu_8$	1117	1099	A'
$\nu_9$	1092	1075	A'
$\nu_{10}$	990	972	A'
$\nu_{11}$	854	838	A'
$\nu_{12}$	710	696	A'
$\nu_{13}$	622	609	A'
$\nu_{14}$	453	449	A'
$\nu_{15}$	261	260	A'
$\nu_{16}$	930	906	A''
$\nu_{17}$	834	813	A''
$\nu_{18}$	710	684	A''
$\nu_{19}$	567	559	A''
$\nu_{20}$	458	450	A''
$\nu_{21}$	193	188	A''

## CAS Active Spaces

Figure S.1: Active orbitals included in CASSCF calculations along  $R_{S-H}$ , numbered in order of increasing energy. The HOMO is orbital 30

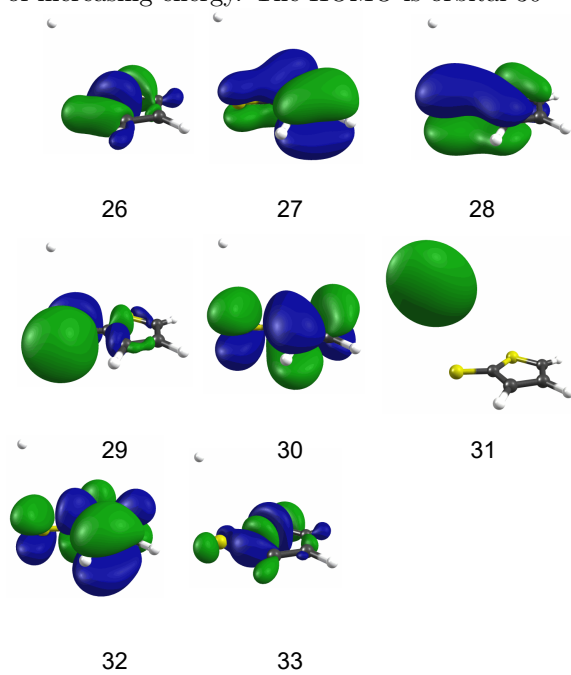


Figure S.2: Larger active space used for CASSCF calculations of transition dipole moments, oscillator strengths and vertical excitation energies. Orbitals are numbered in order of increasing energy. The HOMO is orbital 30

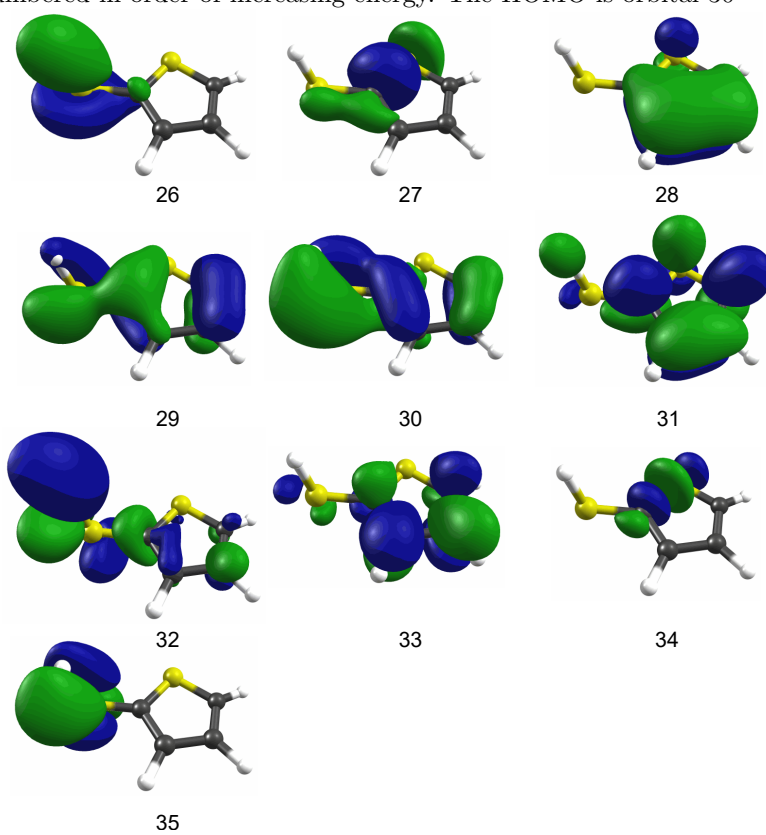
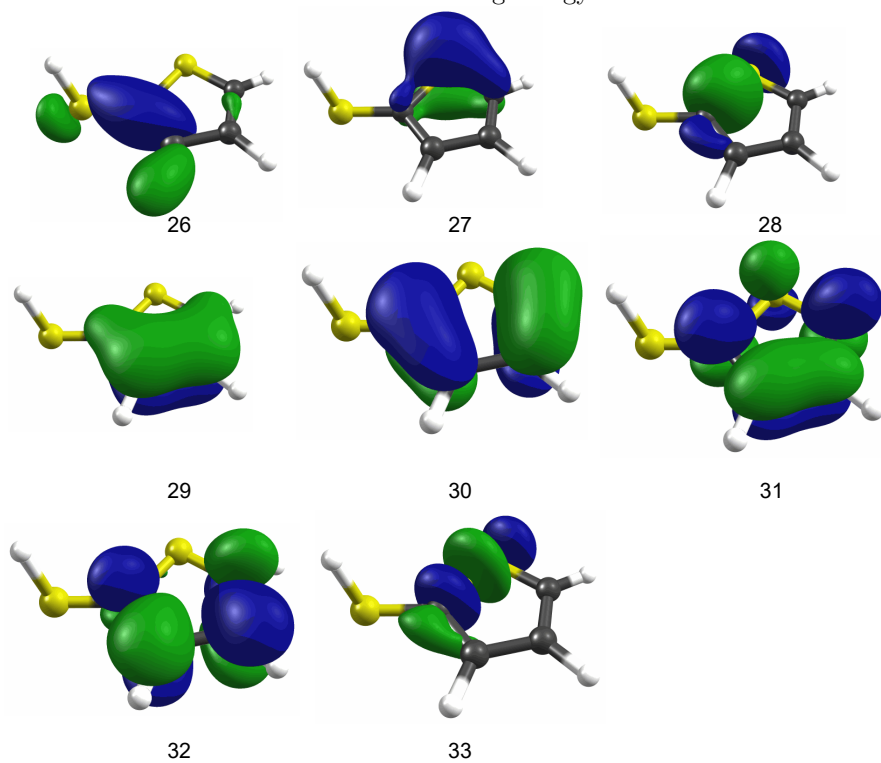


Figure S.3: 10/8 active orbitals included in CASSCF calculations along  $R_{C-S}$ . Orbitals are numbered in order of increasing energy. The HOMO is orbital 30



## Kinetic modelling of the H<sup>+</sup> signal transients

H<sup>+</sup> signal transients obtained from time-resolved velocity map ion imaging (TR-VMI) measurements are modelled using a combination of exponential rise and decay functions, which are convoluted with the Gaussian instrument response function,  $g(\Delta t)$ .

We use a combination of: (i) a series of exponential rises,  $f_n(\Delta t)$ , which describe the rise in signal at  $\Delta t > 0$  attributed to probed H-atoms generated through dissociation of the parent molecule; (ii) an exponential decay,  $d(\Delta t)$ , which accounts for the sharp ‘spike’ in signal around  $\Delta t = 0$ ; and (iii) a second exponential rise function at  $\Delta t < 0$ ,  $r(\Delta t)$ , which models the ‘reverse dynamics’ observed at all photoexcitation wavelengths. These functions are each convoluted with the Gaussian instrument response function,  $g(\Delta t)$ , and have the forms:

$$f_n(\Delta t) = g(\Delta t) * \left[ 1 - A_n \exp\left(-\frac{\Delta t}{\tau_{fn}}\right) \right]$$
$$d(\Delta t) = g(\Delta t) * \left[ B \exp\left(-\frac{\Delta t}{\tau_d}\right) \right]$$
$$r(\Delta t) = g(\Delta t) * \left[ 1 - C \exp\left(-\frac{\Delta t}{\tau_r}\right) \right]$$

where  $A_n$ ,  $B$  and  $C$  are the amplitudes of each function,  $\tau_f$ ,  $\tau_d$  and  $\tau_r$  are the time-constants associated with each exponential function and  $\Delta t$  is the pump-probe time delay. The overall fit is subsequently given as:

$$a(\Delta t) = \sum_n f_n(\Delta t) + d(\Delta t) + r(\Delta t)$$

where  $n$  is the minimum number of rise functions necessary to fit the data accurately.

Component fits presented in Figure 8 of the main manuscript have corresponding time-constants given as follows: (a)  $\tau_{260} = 133 \pm 24$  fs; (b)  $\tau_{238,\text{fast}} = 124 \pm 17$  fs; and (c)  $\tau_{238,\text{slow}} = 1.2 \pm 0.15$  ps, with  $\tau_d < 30$  fs and  $\tau_r < 30$  fs (see below).

## Explanation of reverse dynamics of the H<sup>+</sup> signal transients

The reverse dynamics observed in the H<sup>+</sup> signal transients correspond to single photon excitation with 243 nm, forming H-atoms which are then two-photon excited with 243 nm to the  $2s$  level and subsequently ionised with a single photon of the pump laser to yield H<sup>+</sup> ions. Precedence for such an explanation has been set in previous work by Roberts *et al.*<sup>1</sup> and Lippert *et al.*<sup>2</sup> Photodissociation to yield H-atoms and two-photon excitation of the  $2s \leftarrow 1s$  transition all needs to occur, therefore, within the temporal width of the 243 nm pulse (80 fs), hence giving rise to the very fast time-constant of  $\tau_r < 30$  fs (see above).

1. G. M. Roberts, C. A. Williams, H. Yu, A. S. Chatterley, J. D. Young, S. Ullrich and V. G. Stavros, *Faraday Discuss.*, 2013, **163**, 95–116
2. H. Lippert, H. H. Ritze, I. V. Hertel and W. Radloff, *ChemPhysChem*, 2004, **5**, 1423–1427