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SUPPORTING INFORMATION FOR

The Temperature Dependence of the Hofmeister Series:

Thermodynamic Fingerprints of Cosolute-Protein Interactions**

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S.1 Analysis of DSC data

S.1.1 Calculation of partial heat capacity

To correct for heat capacity differences between the sample and reference cell a buffer scan (buffer vs. buffer) was subtracted from the protein scan (protein vs. buffer) to obtain $\Delta C_p^{app}(T)$. The partial heat capacity of RNase A $C_{p,pr}^{exp}(T)$ was then calculated according to eq. 1, where *T* is the absolute temperature, $C_{p,water}$ the heat capacity of water, V_{pr}^{part} is the partial specific volume of RNase A in the sample, V_{tot} is the total volume of the sample solution and *n* is the amount of protein in mol. V_{pr}^{part} of RNase A was calculated based on the partial molar volumes of the amino acids at 25 °C as described by Makhatadze *et al.*¹

$$C_{\rm p,pr}^{\rm exp}(T) = \left(\Delta C_{\rm p}^{\rm app}(T) + C_{\rm p,water}(T) \frac{V_{\rm pr}^{\rm part}}{V_{\rm tot}}\right) / n$$
¹

S.1.2 Two-state transition model for DSC analysis

All DSC scans proved to be reversible indicating an equilibrium between folded and unfolded state at any temperature. This allows a meaningful extraction of equilibrium thermodynamic data from the thermograms. In the two-state protein folding model the melting temperature, T_m , is defined as the temperature at which the populations of the folded (F_N) and unfolded state (F_U) are equal. Both populations can be related to the equilibrium constant, K (eq. 2), where R is the gas constant and ΔG_u the Gibbs free energy change of unfolding,

$$K = e^{-\Delta G_{\rm u}/RT}$$
 2

yielding

$$F_{\rm N}(T) = 1/(1+K)$$
 3

$$F_{\rm U}(T) = K/(1+K) \tag{4}$$

Standard thermodynamics relates ΔG_u to the heat capacities of the folded and unfolded protein ($C_{p,N}$ and $C_{p,U}$, respectively) by the integrated form of the Gibbs-Helmholtz equation (eq. 5). Assuming a temperature-independent heat capacity change upon unfolding, $\Delta C_p = C_{p,U} - C_{p,N}$, the integrated Gibbs-Helmholtz equation is given by

$$\Delta G_{\rm u} = T_{\rm m} - T/T_{\rm m} * \Delta H_{\rm u,fit}(T_{\rm m}) + \Delta C_{\rm p} * (T - T_{\rm m}) + T * \Delta C_{\rm p} * \ln(T_{\rm m}/T), \qquad 5$$

Where $\Delta H_{u,fit}(T_m)$ is the enthalpy change upon unfolding at T_m . Likewise, the temperature dependence for the enthalpy reads

$$\Delta H_{\rm u} = \Delta H_{\rm u,fit} / (T_{\rm m}) + \Delta C_{\rm p} * (T - T_{\rm m}), \qquad 6$$

resulting in an excess heat capacity C_p^{exc} .^{2, 3}

$$C_{\rm p}^{\rm exc}(T) = [\Delta H_{\rm u}(T)^2 / (R * T^2)] * [K / (1 + K)^2],$$
7

Taken together, we obtain the final expression for the two-state model

$$C_{p,pr}(T) = F_{N}(T) * C_{p,N}(T) + C_{p}^{exc}(T) + F_{U}(T) * C_{p,U}(T),$$
8

which can be fitted to $C_{p,pr}^{exp}(T)$ of RNase A to obtain T_m , $\Delta H_{u,fit}(T_m)$, and $\Delta C_p(T_m)$. Here, $C_{p,N}(T)$ and $C_{p,U}(T)$ of RNase A are represented by linear equations (eqs. 9&10).

$$C_{\rm p,N}(T) = A_{\rm N} * (T - 273.15) + B_{\rm N}$$
9

$$C_{\rm p,U}(T) = A_{\rm U} * (T - 273.15) + B_{\rm U}$$
 10

S.1.3 Calculation of excess heat capacity profile

The so-called excess heat capacity function $\langle C_p(T) \rangle^{\text{exc}}$ is caused solely by the heat required to unfold the protein. It is calculated by subtracting the progress heat capacity function $C_{p,pr}^{prg}(T)$. The respective equations are

$$C_{p,pr}^{prg}(T) = F_{N}(T) * C_{p,N}(T) + F_{U}(T) * C_{p,U}(T),$$
11

and

$$< C_{\rm p}(T) >^{\rm exc} = C_{\rm p,pr}^{\rm exp}(T) - C_{\rm p,pr}^{\rm prg}(T), \qquad 12$$

The excess heat capacity of RNase A is then used to calculate the calorimetric enthalpy change at $T_{\rm m}$, $\Delta H_{\rm u,cal}(T_{\rm m})$ (eq. 13), and the van't Hoff enthalpy change at $T_{\rm m}$, $\Delta H_{\rm u,vH}(T_{\rm m})$ (eq. 14), where $< C_{\rm p}(T_{\rm m}) >^{\rm exc,max}$ is the maximum of $< C_{\rm p}(T) >^{\rm exc}$. Note, $\Delta H_{\rm u,fit}$ is a van't Hoff enthalpy as well.

$$\Delta H_{\rm cal}(T_{\rm m}) = \int_0^\infty \langle C_{\rm p}(T) \rangle^{\rm exc} \, \mathrm{d}T$$
¹³

$$\Delta H_{\rm vH}(T_{\rm m}) = \frac{4RT_{\rm m}^2 < C_{\rm p}(T_{\rm m}) >^{\rm exc,max}}{\Delta H_{\rm cal}(T_{\rm m})}$$
¹⁴

S.1.4 Calculation of excess thermodynamic parameters

The thermodynamic excess parameters $\Delta\Delta G_u$, $\Delta\Delta H_u$, and $T\Delta\Delta S_u$ are defined as $\Delta\Delta X_u = \Delta X_{u,cosolute} - \Delta X_{u,buffer}$ with X = G, H, S. The excess parameters were calculated at $T_{m,buffer}$ as a common temperature. $\Delta H_{u,fit}$ of each dataset was used for the calculations. No significant changes of ΔC_p in the presence of cosolutes could be observed experimentally (see S.2 for details). Therefore, the reported temperature dependence (278-398 K) of ΔC_p in a cosolute-free solution by Privalov and Makhatadze⁴ was fit to a second order polynomial ($\Delta C_p(T) = -35.26 + 0.2603T - 0.0004190T^2$) and used for the calculation of $\Delta\Delta G_u$, $\Delta\Delta H_u$, and $T\Delta\Delta S_u$. The uncertainties of the excess functions were calculated according to Gaussian error propagation using the systematic error estimation of $T_{\rm m}$ and $\Delta H_{\rm u,fit}$ from three repeated DSC measurements in buffer (Table S1). For a few scans the statistical error of $T_{\rm m}$ and $\Delta H_{\rm u,fit}$ was larger than the systematic error estimated from the buffer scans. In these cases the larger statistical error was used for the error propagation.

Table S1: T_m and $\Delta H_{u,fit}(T_m)$ obtained from the fit of eq. 8 to $C_{p,pr}^{exp}(T)$ and the results for $\Delta H_{u,cal}(T_m)$ and $\Delta H_{u,vH}(T_m)$ (eqs. 13&14, respectively) are listed. Different buffer scans were used as a reference because of different experimental conditions (VP-DSC instrument^{n,d}, Capillary-DSC instrument^{h,c}, protein concentration, buffer, and scan rate). For all scans the respective reference measurement (a-d) is given with a superscript. For reference a, errors of T_m and ΔH_{fit} are representative of the 90% confidence interval of the mean. This error of the mean was derived from three repeated measurements using a Student's t-distribution, which considers the limited repetition of measurements, errors of $\Delta H_{u,fit}$ are the statistical errors obtained from the fit. The statistical error of most of the melting temperatures were in the region of 0.01-0.05 K and therefore not listed. ^e: Enthalpy could not be estimated reliably from the first protein scan. The given value was extrapolated (linear function or second-order polynomial) based on the data of the subsequent protein scans. Because of that an error estimation (Gaussian error propagation of the standard errors of the fit) is given for ΔH_{vH} in these cases as well.

Solution	Conc. (M)	<i>T</i> _m , (K)	$\Delta H_{\rm fit}$, (kJ/mol)	ΔH_{cal} , (kJ/mol)	$\Delta H_{\rm vH}$, (kJ/mol)
50 mM citrate	(a)	338.89 ± 0.08	514 ± 12	509 ± 9	515 ± 16
buffer, pH 5.0,	(b)	338.7	505.4 ± 0.6	499	517
$[pr.]= 0.5 \text{ mg mL}^{-1},$	(c)	338.6	473.6 ± 0.6	467	496
Scan rate: 90 K h ⁻¹		338.5	477.1 ± 0.6	471	501

10 mM phosphate	(d)	335.7	459.9 ± 0.9	461	478
buffer, pH 5.5					
[pr.]= 5.0 mg mL ⁻¹ ,					
Scan rate: 60 K h ⁻¹					
LiCl ^c	0.25	335.8	432.9 ± 0.8	429	456
(Sigma-Aldrich, St.	0.5	335.0	448.9 ± 0.9	444	464
Louis, USA)	1	334.5	442.8 ± 1.3	440	456
	2	334.4	441.1 ± 0.6	440	444
	4	326.8	374.9 ± 2.2	368	391
NaCl ^a	0.25	337.1	491.2 ± 1.3	488	492
(Carl Roth, Karls-	0.5	336.8	485.5 ± 1.3	481	486
ruhe, Germany)	1	337.8	485.7 ± 1.8	481	487
	2	339.5	456.1 ± 0.8	446	479
	4	342.7	467.2 ± 3.4	461	475
KCl ^a	0.25	337.0	488.7 ± 1.2	486	489
(Mallinckrodt	0.5	336.6	486.2 ± 1.4	483	485
Baker, Deventer,	1	337.2	472.9 ± 2.4	468	476
Netherlands)	2	339.8	457.0 ± 2.2	454	462
CsCl ^a	0.25	336.3	477.1 ± 4.2	470	487
(Sigma-Aldrich, St.	0.5	335.1	446.4 ± 0.9	436	476
Louis, USA)	1 ^e	335.8	461.3 ± 8.5	459 ± 15	458 ± 15
	2 ^e	337.6	445.3 ± 4.7	442 ± 8	455 ± 8
	4 ^e	339.9	425.5 ± 6.3	423 ± 10	435 ± 10
NaBr ^a	0.25	335.5	487.2 ± 2.1	478	489
(ABCR, Karlsruhe,	0.5	334.1	459.0 ± 0.9	455	461
Germany)	1	332.9	446.1 ± 1.8	442	448
	2	332.0	423.2 ± 0.6	420	424
	4	324.8	360.1 ± 4.9	355	367
NaClO4 ^a	0.25	332.4	457.5 ± 0.9	453	460

(Sigma-Aldrich, St.	0.5	329.2	433.3 ± 0.7	430	435
Louis, USA)	1	324.3	401.0 ± 1.2	398	405
	2	316.4	360.5 ± 0.8	359	366
	4	301.9	300.1 ± 1.6	279	360
[chol]Cl ^a	0.25	336.8	489.6 ± 1.3	485	492
(TCI, Tokyo, Japan)	0.5	335.8	476.2 ± 1.0	472	478
	1	335.6	470.9 ± 1.1	467	473
	2	337.0	482.0 ± 2.0	476	484
	4	339.1	480.5 ± 1.0	477	481
[Me ₄ N]Cl ^a	0.25	336.7	496.0 ± 1.6	490	499
(Sigma-Aldrich, St.	0.5	335.6	476.3 ± 1.6	463	500
Louis, USA)	1	336.2	485.8 ± 1.6	480	487
	2	337.5	491.0 ± 1.9	485	491
	4	337.6	488.3 ± 0.7	453	488
K[dhp] ^{a,b}	0.125ª	340.2	511.9 ± 1.8	507	516
(VWR, Leuven,	0.25 ^a	341.5	527.6 ± 3.6	518	536
Belgium)	0.5 ^a	343.9	536.1 ± 2.3	532	537
	1 ^{b,e}	347.4	507.0 ± 1.5	507 ± 3.5	517 ± 3.5
Ethanol ^{a,b}	0.25 ^b	338.1	513.8 ± 1.1	509	520
(Sigma-Aldrich, St.	0.5 ^b	337.4	521.6 ± 1.2	516	528
Louis, USA)	1 ^a	335.6	502.9 ± 0.9	493	526
	2 ^a	333.6	546.5 ± 3.7	538	552
	4 ^a	328.6	556.1 ± 4.5	542	565
Ethylene glycol ^a	0.25	338.7	519.9 ± 2.2	512	523
(Sigma-Aldrich, St.	0.5	338.6	520.7 ± 2.1	513	523
Louis, USA)	1	338.5	519.8 ± 1.6	515	520
	2	338.3	531.0 ± 1.8	525	534
	4	337.7	546.8 ± 4.3	539	553
Glycerol ^a	0.25	339.1	509.7 ± 3.0	504	514

(Sigma-Aldrich, St.	0.5	339.1	501.7 ± 2.4	497	509
Louis, USA)	1	339.6	522.6 ± 2.6	520	523
	2	340.0	519.5 ± 2.4	502	548
	4	342.9	555.2 ± 2.0	550	557
[emim]Br ^d	0.25	331.9	448.3 ± 1.2	442	484
(IoLiTec, Denz-lin-	0.75	326.8 ± 0.3	412 ± 13	402	457
gen, Germany)	1	324.7 ± 0.4	398 ± 13	388	448
[bmpyrr]Br ^d	0.25	331.7	460.7 ± 1.4	462	475
(IoLiTec, Denz-lin-	0.5	328.2 ± 0.2	428.6 ± 5.1	418	472
gen, Germany)	1	323.1 ± 0.2	407 ± 11	392	471
	1.5	319.9 ± 0.2	426.3 ± 4.9	419	469
[bmim]Br ^d	0.25	330.9	490.0 ± 0.7	483	507
(IoLiTec, Denz-lin-	0.5	326.5	473.1 ± 2.2	462	498
gen, Germany)	1	319.5	448.7 ± 1.5	437	485
	1.5	316.6	438.9 ± 1.3	428	477
[Bu4N]Br ^d	0.25	327.0	472.0 ± 1.1	473	493
(IoLiTec, Denz-lin-	0.5	320.1	477.6 ± 1.6	478	502
gen, Germany)	1	310.2	400.9 ± 1.9	389	432
[hmim]Br ^d	0.25	328.3	440.6 ± 1.2	440	460
(IoLiTec, Denz-lin-	0.5	320.8	398.2 ± 1.6	395	438
gen, Germany)	0.75	316.0	387.2 ± 4.1	379	430
	1	311.9	376.9 ± 2.3	371	415
	1.5	305.39	353.9 ± 1.9	355	392
[bmim][BF ₄] ^d	0.5	326.6	452.5 ± 0.8	446	477
(IoLiTec, Denz-lin-	1	322.1	426.7 ± 1.2	425	448
gen, Germany)	1.5	317.8	398.5 ± 1.0	386	444
[emim]Cl ^d	0.5	331.5	458.5 ± 1.4	457	476
(IoLiTec, Denz-lin-	1	328.5 ± 0.2	424.9 ± 5.1	414	471
gen, Germany)					

[emim][EtSO4] ^d	0.25	332.8	454.1 ± 0.9	453	474
(IoLiTec, Denz-lin-	0.5	330.5	451.7 ± 1.0	452	472
gen, Germany)	0.75	328.9	447.4 ± 1.1	456	461
	1	327.6	447.9 ± 1.1	447	469
[emim][dca] ^d	1	314.8	349.5 ± 2.1	345	388
(IoLiTec, Denz-lin-					
gen, Germany)					
[emim][SCN] ^d	0.25	328.5	443.6 ± 0.7	434	480
(IoLiTec, Denz-lin-	0.5	322.4	398.6 ± 0.7	387	456
gen, Germany)	0.75	318.2	392.3 ± 0.8	387	420
	1	314.7 ± 0.2	373.4 ± 4.3	365	410
	1.5	308.5	363.7 ± 1.7	351	409
[chol][dhp] ^d	0.5	341.4	511.4 ± 3.8	499	563
(IoLiTec, Denz-lin-	1	344.2	518.9 ± 1.6	523	526
gen, Germany)	1.5	347.5	528.0 ± 1.2	527	537
	3	353.2	519.4 ± 3.9	522	536
	4	356.6 ± 0.2	501.4 ± 6.4	495	537

S.2 Estimation of ΔC_p in the presence of cosolutes

S.2.1 Estimation of ΔC_p of RNase A from linear extrapolation of $C_{p,\text{folded}}$ and $C_{p,\text{unfolded}}$

The success of measuring ΔC_p depends largely on the data quality of the partial molar heat capacity of the folded and unfolded protein. High concentrations of cosolute can cause small distortions of the baseline making it impossible to reliably determine ΔC_p from the fit described in S.1.2. A more accurate way is to determine ΔC_p independently from the fit of the two-state model to the $C_{p,part.mol.}$. $C_{p,folded}$ and $C_{p,unfolded}$ can be described by linear equations (eqs. 9&10). The difference between both linear functions results in a linear function for $\Delta C_p(T)$ (Fig. SI). Since $C_{p,unfolded}$ has a nonlinear temperature dependence⁵ the linear extrapolation holds only for a small temperature range and should therefore only be used to estimate $\Delta C_p(T_m)$ and not to estimate the full temperature dependence of ΔC_p .



Figure SI: Estimation of ΔC_p from linear extrapolations of $C_{p,\text{folded}}$ and $C_{p,\text{unfolded}}$. $C_{p,\text{folded}}$ and $C_{p,\text{unfolded}}$ are linear fits to the blue data points. Data points shown in red describe the excess heat capacity due to the unfolding of the protein and where not used here to determine ΔC_p .

Using this approach, we determined $\Delta C_p(T_m)$ for all DSC scans for which the quality of the data of $C_{p,\text{folded}}$ and $C_{p,\text{unfolded}}$ allowed a linear extrapolation. Data for $\Delta C_p(T_m)$ are shown in Table S1. We observed no significant trends of $\Delta C_p(T_m)$ with increasing cosolute concentration or type of cosolute. $\Delta C_p(T_m)$ remained rather constant with respect to the buffer solution considering the relatively high uncertainty of estimating ΔC_p .

Table S2: Estimation of $\Delta C_p(T_m)$ using a linear extrapolation of $C_{p,\text{folded}}$ and $C_{p,\text{unfolded}}$.

Solution	Conc. (M)	$\Delta \boldsymbol{\mathcal{C}}_{\boldsymbol{p}}(\boldsymbol{T}_{\mathbf{m}}),$ (kJ mol ⁻¹ K ⁻¹)
50 mM citrate		3.8
buffer, pH 5.0		3.7
		4.2
		3.7

10 mM phosphate		
buffer, pH 5.5		4.3
Privalov ⁴		
323 K		5.3
348 K		4.3
NaCl	0.25	3.9
	0.5	3.8
	1	3.7
	2	4.5
	4	5.0
KCl	0.25	3.5
	0.5	3.4
	1	3.4
	2	3.0
NaBr	0.25	4.3
	0.5	3.6
	1	3.6
	4	3.2
NaClO ₄	0.25	4.8
	0.5	3.7
	1	6.1
	2	4.6
[chol]Cl	0.25	4.5
	0.5	3.4
	1	3.9
	2	3.8
	4	2.3
[Me ₄ N]Cl	0.25	5.8
	0.5	3.6

	1	3.9
	2	4.4
K[dhp]	0.125	3.9
	0.25	3.8
	0.5	4.2
Ethylene glycol	0.25	3.6
	0.5	4.7
	1	2.9
	2	4.0
glycerol	0.25	4.0
	0.5	4.5
[bmim]Br	0.25	5.4
[Bu4N]Br	0.25	4.6
[hmim]Br	0.25	4.0
[emim][EtSO ₄]	0.25	5.3
	0.5	6.4
	0.75	6.2
	1	5.3
[emim][dca]	1	4.5
[emim][SCN]	0.75	4.4

S.2.2 Estimation of ΔC_p of ubiquitin in ethanol solutions

To further check the influence of cosolutes on ΔC_p we choose ubiquitin as a model protein because it has a very high ΔC_p in comparison to its number of amino acids (0.075 kJ mol⁻¹ K⁻¹ N⁻¹ at 25°C).⁶ In contrast, RNase A is at the lower end of this scale (0.042 kJ mol⁻¹ K⁻¹ N⁻¹ at 25°C).⁴ Therefore, changes of ΔC_p are easier detectable using ubiquitin compared to RNase A. We have chosen 4 M ethanol to test the influence of a hydrophobic cosolute on ΔC_p since ΔC_p is related to the hydration of apolar residues.⁵ Thus, it might be influenced by the presence of a hydrophobic cosolute. Since ΔC_p in the presence of 4 M ethanol could not be reliably determined from the fit of the two-state transition model to the data or from the linear extrapolations of $C_{p,\text{folded}}$ and $C_{p,\text{unfolded}}$, we used a different procedure to obtain ΔC_p . Running DSC scans to very high temperatures (> 120°C) increases the irreversibility of the folding reaction. Thus, after a series of scans only the partial molar heat capacity of the unfolded protein is measured (Fig. SIIa). Subtraction of $C_{p,\text{folded}}$ (measured at the first scan) from $C_{p,unfolded}$ results in $\Delta C_p(T)$ (Fig. SIIb). Interestingly, even in 4 M ethanol ΔC_p seems not to be affected by the presence of the cosolute (at least in the temperature range shown in Fig. SIIb).



Figure SII: a) DSC scans of ubiquitin (1.0 mg mL⁻¹) in H2O at pH 3. b) $\Delta C_{p,\text{water}}$ and $\Delta C_{p,4M \text{ EtOH}}$ of ubiquitin as a function of temperature in the low temperature regime where the protein is folded. $\Delta C_{p,\text{water}}$ and $\Delta C_{p,4M \text{ EtOH}}$ are within the error bars of each other (standard error of the mean of three repeated measurements).

S.3 Correlation between $\Delta \Delta G_{\rm u}$ and $\Delta T_{\rm m}$

Fig. S4 shows a plot of $\Delta\Delta G_u$ versus ΔT_m for all cosolute data listed in Table S1. A linear regression of 97 data points imposing a zero intercept resulted in a slope of 1.39 kJ mol-1 K-1 with a R^2 of 0.996. Thereby, the slope corresponds to $\Delta H(T_m)/T_m$ of RNase A in the cosolute-free solution which follows from the derivative of $\Delta G(T)$ with respect to T at T_m (eq. 15).



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Figure S1: Correlation plot of $\Delta\Delta G_u$ vs. ΔT_m for all cosolute data listed in Table S1. The blue line corresponds to the linear fit to all data points (y = 1.39x). The red line corresponds to a linear curve with a slope of $\Delta H_u(T_m)/T_m$ of RNase A in the cosolute-free solution (y = 1.45x, calculated via a weighted average of all buffer data (Table S1)).

S.4 Influence of the temperature dependence of ΔC_p on the stability curve

While the heat capacity of the folded state, $C_{p,\text{folded}}$, can be described by a linear equation, the heat capacity of the unfolded state, $C_{p,\text{unfolded}}$, is best described by a second-order polynomial.⁵ Therefore, ΔC_p can be described by a second order polynomial as well. However, extrapolation of the second-order polynomial ΔC_p to extreme temperatures leads to a strong decrease of ΔC_p at low and high temperatures and eventually to negative ΔC_p values. This leads to an increase of protein stability at very low and very high temperatures (Fig. S2b) which is not physically meaningful. In order to estimate the temperature of maximum stability, T_s , and the temperature of cold denaturation, T_c , $\Delta C_p = \text{const.}$ has to be used to calculate $\Delta G_u(T)$ at these low temperatures. In a previous study, we have shown that using $\Delta C_p = \text{const.}$ can lead to large errors for $\Delta \Delta H_u$, and $T \Delta \Delta S_u$.⁷ However, because of strong enthalpy-entropy compensation effects (errors cancel out to a large extent) the simplification of $\Delta C_p = \text{const.}$ leads to meaningful results for $\Delta G_u(T)$ over wide temperature range (Fig. S2) as also observed by Privalov.⁸



Figure S2: Protein stability curves of RNase A in 50 mM citrate buffer (pH 5.0) calculated with $\Delta C_p(T) = -35.26 + 0.2603T - 0.0004190T^2$ (fit of second-order polynomial to the ΔC_p data published by Privalov and Makhatadze⁴) and $\Delta C_p = 5$ kJ mol⁻¹K⁻¹. a) Stability curve shown in the measured temperature range of ΔC_p (5-125 °C). b) Extrapolation to extreme temperatures leads to an increase in protein stability at very low temperatures using $\Delta C_p(T)$ because $\Delta C_p(T)$ turns negative.



Figure S3: Concentration dependence of $\Delta T_{\rm m}$ for nonelectrolytes. Error bars were calculated by Gaussian error propagation of the primary error estimates (see S.1) and are smaller than the point size. The data for sorbitol and urea are taken from Ravindra and Winter.⁹

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