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## I.

**Table 1** Comparison of  $O_2$  and Ni properties with experiments in isolated and adsorbed states.

	<b>^</b>		
Ni	EXP.	$O_2$	EXP.
3.52 Å	3.52 Å [1]		
0.010Å	0.007 <u>+</u> 0.003Å [2]		
0.60 µ <sub>B</sub>	0.60-0.62 μ <sub>B</sub> [1,3]		
		1.23 Å	1.21[4]
		1.96 μ <sub>B</sub>	
	O atom on Ni(111)		
1	1.86-1.89Å 1.85 ± 0		05 Å [5]
	1.20 Å 1.21 ± 0.09Å [5]		09Å [5]
	3.52 Å 0.010Å 0.60 μ <sub>B</sub>	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$

## References

[1] F Starrost, H Kim, S C Watson, E Kaxiras and E A Carter, *Phys. Rev. B*, 2001, 64, 235105.
[2] T Okazawa, F Takeuchi and Y. Kido *Phys. Rev. B*, 2005, 72, 075408.
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equation (1), we get the

## II.

Further discussions for Equation (1) in the manuscript:

The  $F_I(r-r_I)$  is a function with a norm (or value) 1 inside the sphere  $\Omega_I$  and smoothly goes

to zero at the boundary so that when it is multiplied to the magnetization density, m(r) in

$$M_I = \int_{\Omega_I} m(r) d^3 r$$

inside the  $\Omega_I$  and 0 at the boundary as expected.

 $\sin(x)$ The appropriate form of  $F_I$  is x because the limit of this function as  $x \rightarrow 0$  is 1 and 0 at  $x = \pi$  and from 0 to  $\pi$ , it decreases smoothly (monotonically). The  $x = \pi (|r - r_I|)/R_I$ , where  $R_I$  is the radius of the sphere,  $\Omega_I$ , can describe such condition for  $F_I$ . If r tends to  $r_I$  (inside

sphere), 
$$x \rightarrow 0$$
,  $F_I(r - r_I) = 1$  and  $M_I = \int_{\Omega_I} m(r) d^3 r$ .  
 $F_I(r - r_I) = 0$ , so  $M_I = 0$ .