

Electronic Supplementary Information

Type-I van der Waals Heterostructure Formed by MoS₂ and ReS₂ Monolayers

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1. Confirmation of Monolayer Thickness

In order to confirm the monolayer (ML) thickness of both the MoS₂ and ReS₂ layers, we utilize a combination of optical contrast, Raman, and photoluminescence (PL) measurements

One of the simplest methods to confirm the layer thickness of a sample has been to utilize optical contrast in different color channels of red, green, and blue. From previously reported studies [1,2], it has been established that ML transition metal dichalcogenides (TMDs) have a green channel contrast of about 12-13%. Figure S1(a) shows the green channel of an optical image of the heterostructure sample. The brightness of the image as a function of position along the yellow line was measured extracted from the image. From the brightness, the contrast was calculated as

$$C(x) = \frac{I_{sub} - I(x)}{I_{sub}},$$

where I_{sub} and $I(x)$ are the brightness of the substrate and brightness at position, x , on the yellow line, respectively. These contrast results are plotted in the inset of Fig. S1(a). It can be seen that both MLs have a similar contrast of about 12-13% range, consistent with previous studies [1,2].

In addition to the contrast measurements, photoluminescence (PL) spectroscopy was used to confirm that the MoS₂ sample is monolayer. It is well known that MoS₂ has a direct bandgap in monolayer that transitions to indirect in bilayer and thicker samples. Figure S1(b) shows the PL spectrum measured for the MoS₂ sample. The PL peak position and overall strength of the signal is consistent with that of ML MoS₂ thus confirming ML thickness of the MoS₂ layer. ReS₂ does not exhibit the direct to indirect transition and thus PL cannot be used to confirm its thickness.

Raman measurements were also used to further confirm monolayer thickness of both layers. MoS₂ is known to have a very noticeable shift in the E_{2g}¹ and A_{1g} peaks that progresses from monolayer to bulk. From Li *et al.* [3], in ML, these peaks should be at 384.7 and 402.8 cm⁻¹, respectively, when excited with a 488-nm laser. These two peaks in bilayer MoS₂ are expected to be at 383.3 and 405.5 cm⁻¹, respectively, and become further separated with an increased number of layers. To account for any offsets in measurements, the difference between the E_{2g} and A_{1g} peaks can be taken to indicate thickness of MoS₂ samples.

In our data, we fit the two MoS₂ peaks with Lorentzian functions, shown in Fig. S1(c). The resulting peak positions for monolayer were found to be 385.05 ± 0.05 and 404.10 ± 0.04 cm⁻¹. The difference between our measured peaks, 19.05 cm⁻¹, is consistent with the expectation for ML. The measured results for bulk MoS₂ are also presented in Fig. S1(c), again matching the expected results.

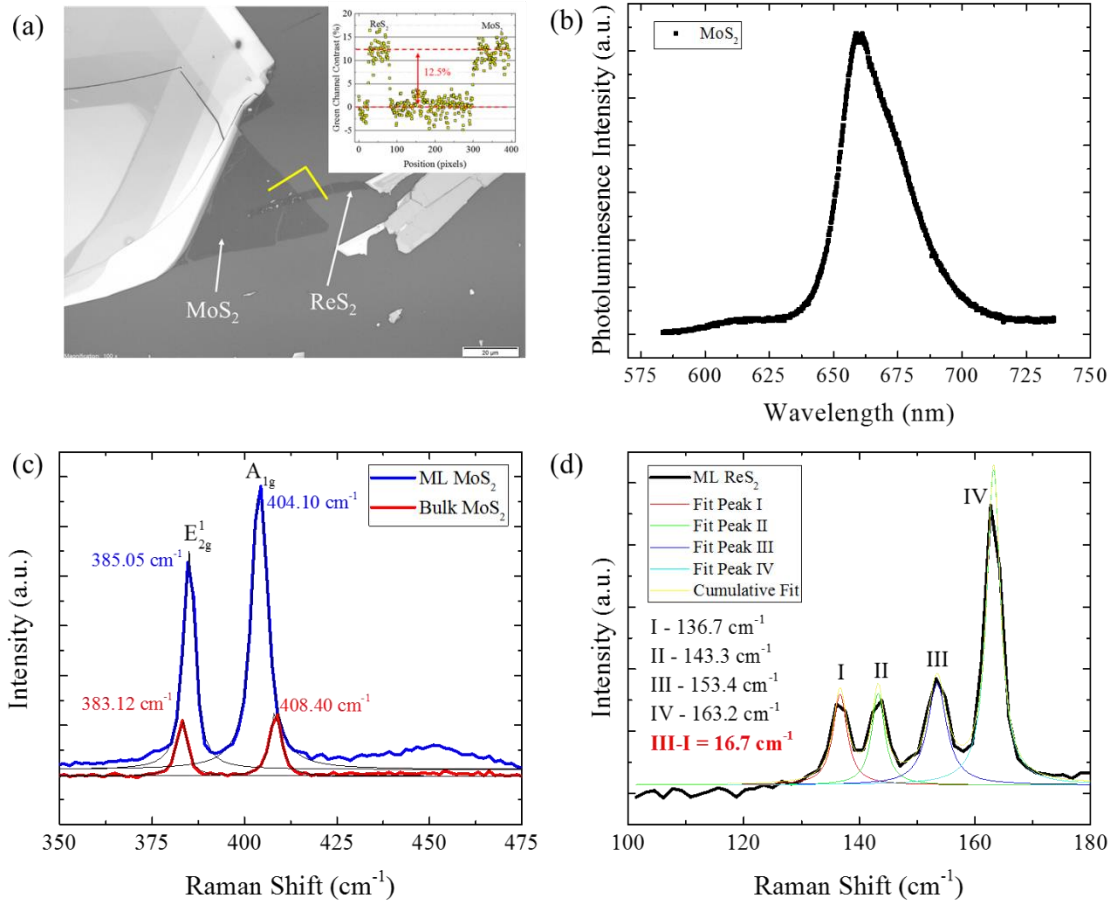


Figure S1. (a) Green channel image of the ReS₂-MoS₂ heterostructure. The inset shows the contrast of the sample along the yellow line. (b) Photoluminescence spectrum of MoS₂ sample. (c) Measured Raman spectra of ML and bulk MoS₂. (d) Measured Raman spectrum of ML ReS₂ between 100 cm⁻¹ and 180 cm⁻¹ (Full spectral range is shown in main text).

The Raman measurements for ReS₂ sample are presented in Fig. S1(d). As discussed in Chenet *et al* [4], four characteristic peaks in ReS₂ between 120-170 cm⁻¹ can be used to identify sample thickness. These peaks are labeled I - IV in Fig. S1(d). Taking the difference between the positions of peaks I and III leads to a significant difference between monolayer and thicker ReS₂. They observed a difference of 16.8 cm⁻¹ between these peaks in ML that drops to 14.9 cm⁻¹ in BL and drops further in thicker samples [4]. Our result of 16.7 cm⁻¹ closely matches their result for ML. Other reports on ReS₂ Raman spectra also show the presence of low frequency, shear and breathing, modes that occur in bilayer ReS₂, but are not present in monolayer [5,6].

2. Power Dependence

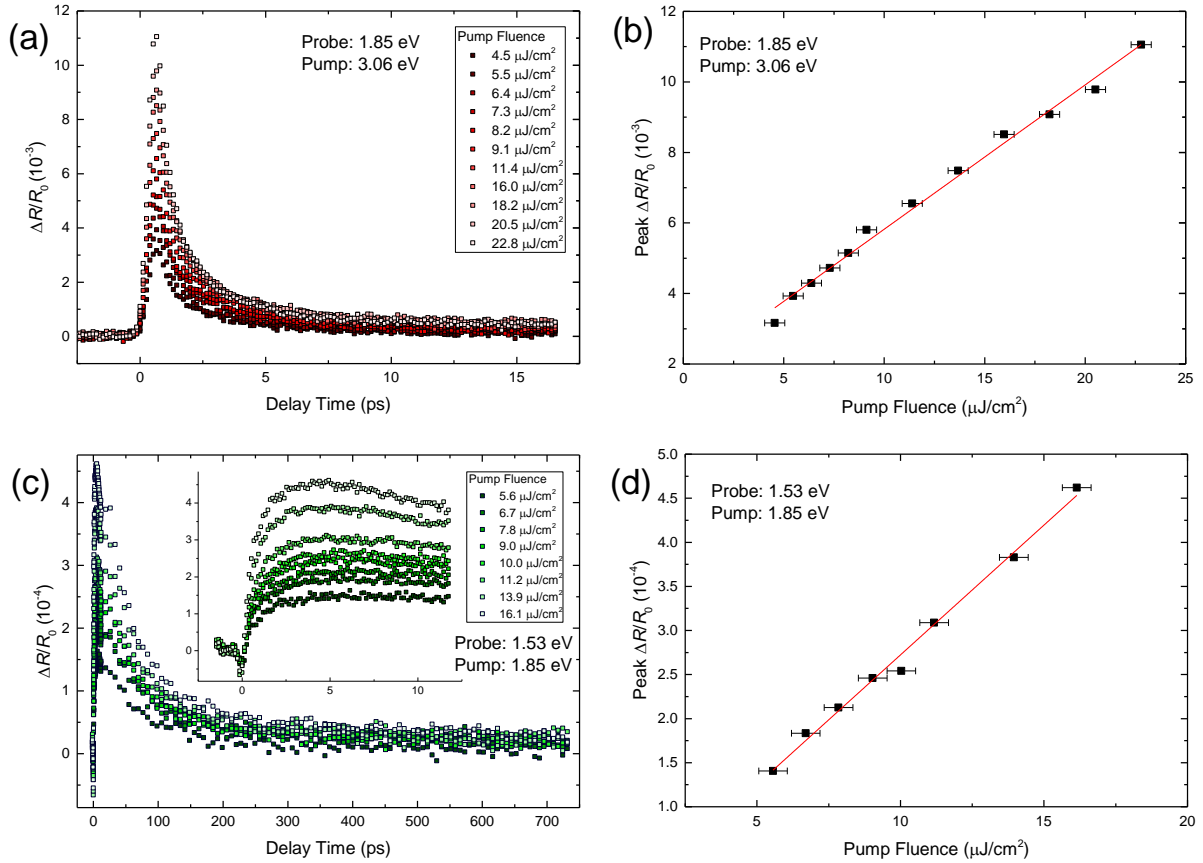


Figure S2. (a) Transient absorption measurements on the MoS₂-ReS₂ heterostructure for various peak fluences of a 3.06 eV pump, using a 1.85 eV probe (near MoS₂ A-exciton resonance). (b) Dependence of the peak $\Delta R/R_0$ signal on the pump fluence of the 3.06 eV pump. The trend line (red) is a linear fit to the data. (c) Transient absorption measurements on the heterostructure for various peak fluences of a 1.85 eV pump (near MoS₂ A-exciton resonance), probing at 1.53 eV (near ReS₂ A-exciton resonance). The inset shows the data on a shorter time scale. (d) Dependence of the peak $\Delta R/R_0$ signal on the pump fluence of the 1.85 eV pump. The trend line (red) is a linear fit to the data.

3. Computation Method:

All computations were carried out within the density functional theory (DFT) methods implemented in the Vienna *ab initio* simulation package (VASP 5.3). The generalized gradient approximation (GGA) within the form of Becke86 was adopted for exchange-correlation functional. Furthermore, the optB86b-vdW functional was used for structure optimization to account for the weak van der Waals (vdW) interactions. The interaction between the core electrons and valence electrons was modeled using the projector augmented wave (PAW) method. Energy cutoff for the plane-wave expansion was set to 500 eV. Brillouin zone sampling was performed with Monkhorst-Pack (MP) special k-point meshes. The 2D Brillouin zone integration using the Γ -center scheme was sampled with a $9 \times 9 \times 1$ grid for transition-metal dichalcogenides. A vacuum

layer greater than 2 nm was taken so that interaction between adjacent images can be neglected. The lattice constants of ReS₂/MoS₂ heterobilayer were fixed during the structure optimization until the computed Hellmann-Feynman force on each atom was less than 0.1 eV/nm. The convergence threshold was set as 10⁻⁵ eV in energy for the structure optimization. For computing the band structure of the ReS₂/MoS₂ heterobilayer, the Brillouin zone of Γ -K-M- Γ was selected.

References

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