

## Supplementary information

### Anhydrous proton conduction in liquid crystals containing benzimidazole moieties

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#### Structural Characterization

<sup>1</sup>H nuclear magnetic resonance (NMR) spectra of **4a** and **1a** are shown in Fig. S1. In the <sup>1</sup>H NMR spectrum of **4a**, the peak at 3.44 ppm was ascribed to the alkyl protons linked to the bromo atom. In the spectrum of **1a**, the peak was at 3.35 ppm arose from the alkyl protons linked to the sulfur atom<sup>1,2</sup>. The signal intensities of the peaks at 7.57, 7.51, 6.97 ppm corresponding to the chemical shifts of protons in aromatic rings increased in the <sup>1</sup>H NMR spectrum of **1a**. The results suggested the successful substitution reaction.

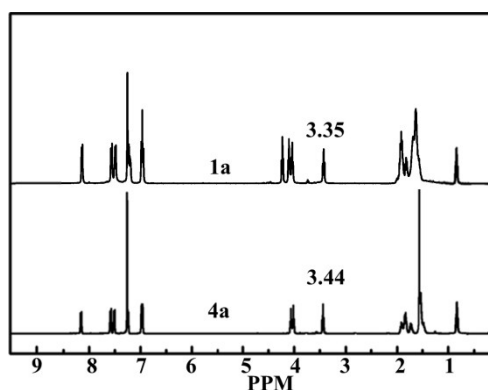


Fig. S1 <sup>1</sup>H NMR spectra of **1a** and **4a** in CDCl<sub>3</sub>

Fourier transform infrared (FT-IR) spectra of **4a**, **1a**, **1b** and **1c** are shown in Fig. S2 and S3. The absorption band at 550cm<sup>-1</sup> in the spectrum of **4a** was assigned to C-Br stretching<sup>4</sup>. In the spectrum of **1a**, the bands at 1346 and 600cm<sup>-1</sup> were ascribed to the C=N stretching and N-H out-of plane bending vibrations of benzimidazole moieties, respectively<sup>3</sup>. The absorption band at 1440cm<sup>-1</sup> was attributed to C-S stretching. A broad absorption centered at 3070cm<sup>-1</sup> was attributed to the intermolecular hydrogen-bonded N-H stretching vibration between benzimidazole moieties<sup>5</sup>.

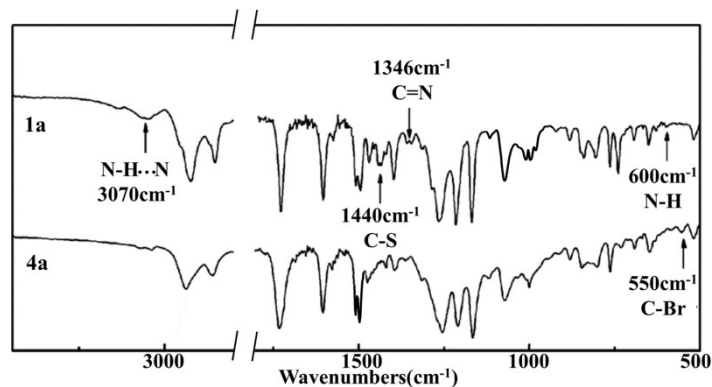


Fig. S2 FT-IR spectra of **1a** and **4a** at room temperature

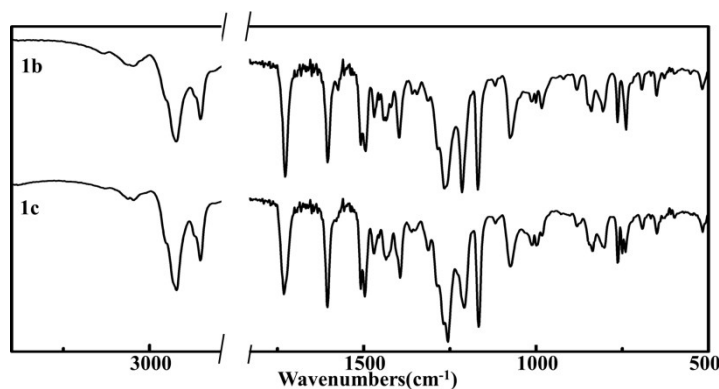


Fig. S3 FT-IR spectra of **1b** and **1c** at room temperature

## Thermal characterization

The differential scanning calorimetry (DSC) traces of **1** during the first cooling and second heating scans are shown in Fig. S4. The multiple peaks indicated appearance of mesophases. XRD patterns of **1a**, **1b** and **1c** in the smectic phase during the cooling process are shown in Fig. S5.

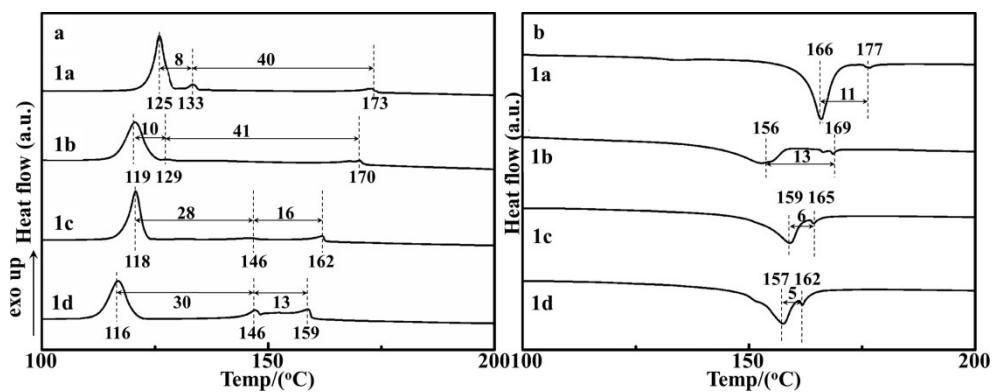


Fig. S4 DSC traces of **1** during the first cooling (a) and second heating (b)

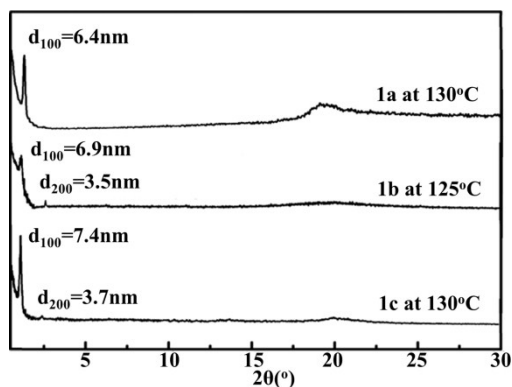


Fig. S5 XRD patterns of **1a**, **1b** and **1c** in the smectic phase during the cooling process

## Electrochemical characterization

Samples were dried under vacuum overnight and sandwiched between two indium tin oxide (ITO) glass slides. The sample thickness ( $d$ ) and cross sectional area ( $A$ ) were controlled by a 150 $\mu$ m thick Teflon spacer with a 0.2 cm diameter hole. The impedance spectrum of **1d** at 140 $^{\circ}$ C is shown in Fig. S6. The bulk resistance was derived from the low intersection of the high-frequency semicircle on a complex impedance plane with the real components axis. Multiple measurements were performed for each sample and the relative standard deviation of the resultant bulk resistance was less than 5%.

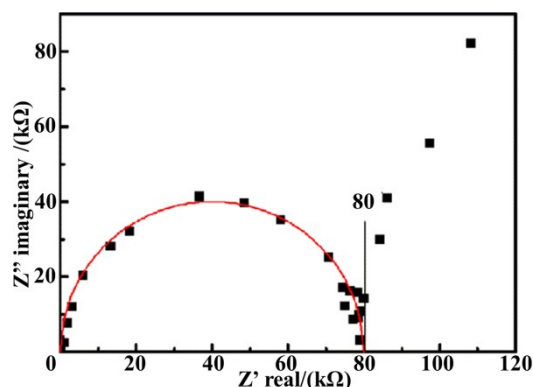


Fig. S6 Impedance spectrum for **1d** at 140 $^{\circ}$ C

The transference numbers corresponding to ion and electron transfer were calculated using Wagner's DC polarization technique with ITO electrodes<sup>7</sup>. Current relaxation curve of **1d** at 140 $^{\circ}$ C during potentiostatic polarization is shown in Fig. S7. The ion transference numbers ( $t_p$ ) were calculated from the curve using the following equation:

$$t_p = \frac{I_s - I_f}{I_s} \quad (\text{S1})$$

Where  $I_s$  is the initial current ( $\mu$ A) and  $I_f$  is the steady state current ( $\mu$ A). For all the samples, the ionic transference numbers were found to be in the range of 0.98~0.99.

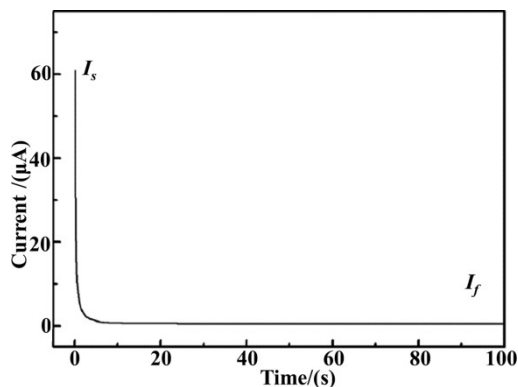


Fig. S7 Current relaxation curve of **1d** at 140°C

Proton conductivity  $\sigma$  ( $\text{S cm}^{-1}$ ) was obtained from the following equation:

$$\sigma = d \cdot R_b^{-1} \cdot A^{-1} \quad (\text{S2})$$

where  $R_b$ ,  $d$  and  $A$  are the bulk resistance ( $\Omega$ ), the sample thickness (cm), and the sample area ( $\text{cm}^2$ ), respectively.

The activation energies ( $E_a$ ) for proton conduction were estimated from the Arrhenius plots.

$$\ln \sigma = \ln A - E_a \cdot R \cdot (T + 273)^{-1} \quad (\text{S3})$$

where  $A$  is the pre-exponential factor ( $\text{S cm}^{-1}$ );  $R$  is known as molar gas constant ( $\text{J}^\circ \text{K}^{-1} \text{mol}^{-1}$ );  $E_a$  is the activation energy ( $\text{kJ mol}^{-1}$ ).

## References

1. A. K. Khalil, Phosphorus Sulfur., 2007, **182**, 815.
2. J. P. Jayachandran, M. Wang, Synth. Commun., 1999, **29**, 4087.
3. V. LUPTÁKOVÁ, G. PLESCH, Clay Miner. 40(2005)295-302.
4. J. Coates, Interpretation of Infrared Spectra, A Practical Approach, in: R.A. Meyers (Ed.), Encyclopedia of Analytical Chemistry, John Wiley & Sons Ltd., Chichester, 2000, pp. 10815.
5. A. Kawada, A. R. McGhie, M. M. Labes, J. Chem. Phys., 1970, 52, 3121.
6. C. Perchard, A. Novak, J. Chem. Phys., 1968, **48**, 3079.
7. S. Selvasekarapandian, R. Baskaran, M. Hema, Phys. B., 2005, **357**, 412.