

## Supplementary Information

# First-principles study of Zr-N crystalline phases: phase stability, electronic and mechanical properties

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## Supplementary Information

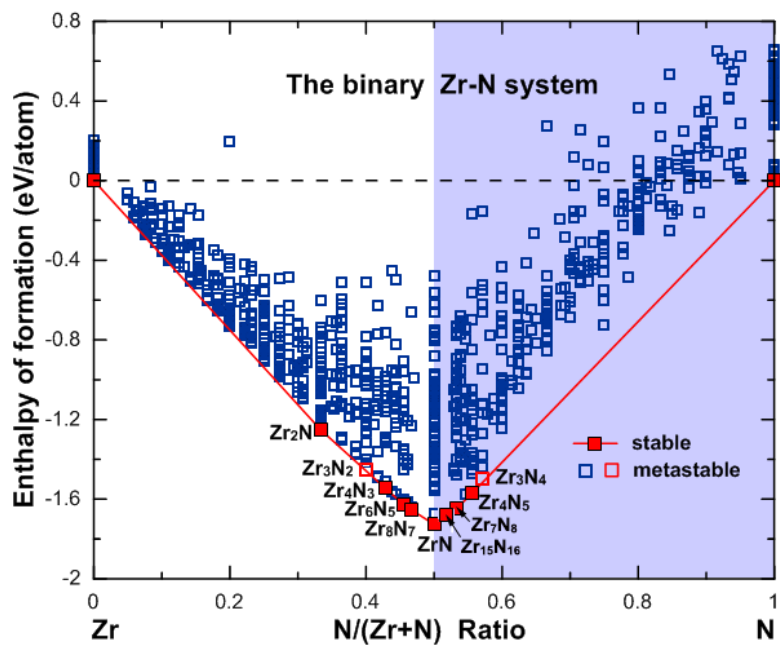


Figure S1 The convex hull diagram of the Zr-N system at 0 GPa.

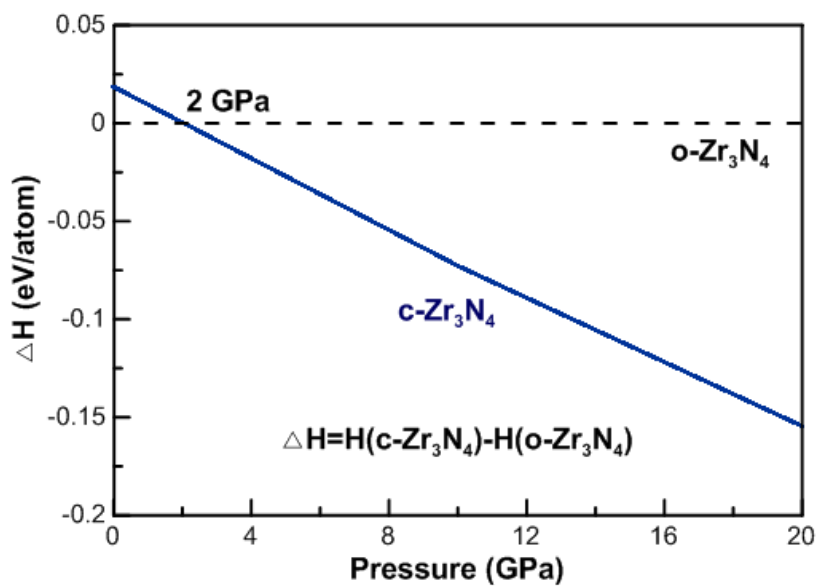
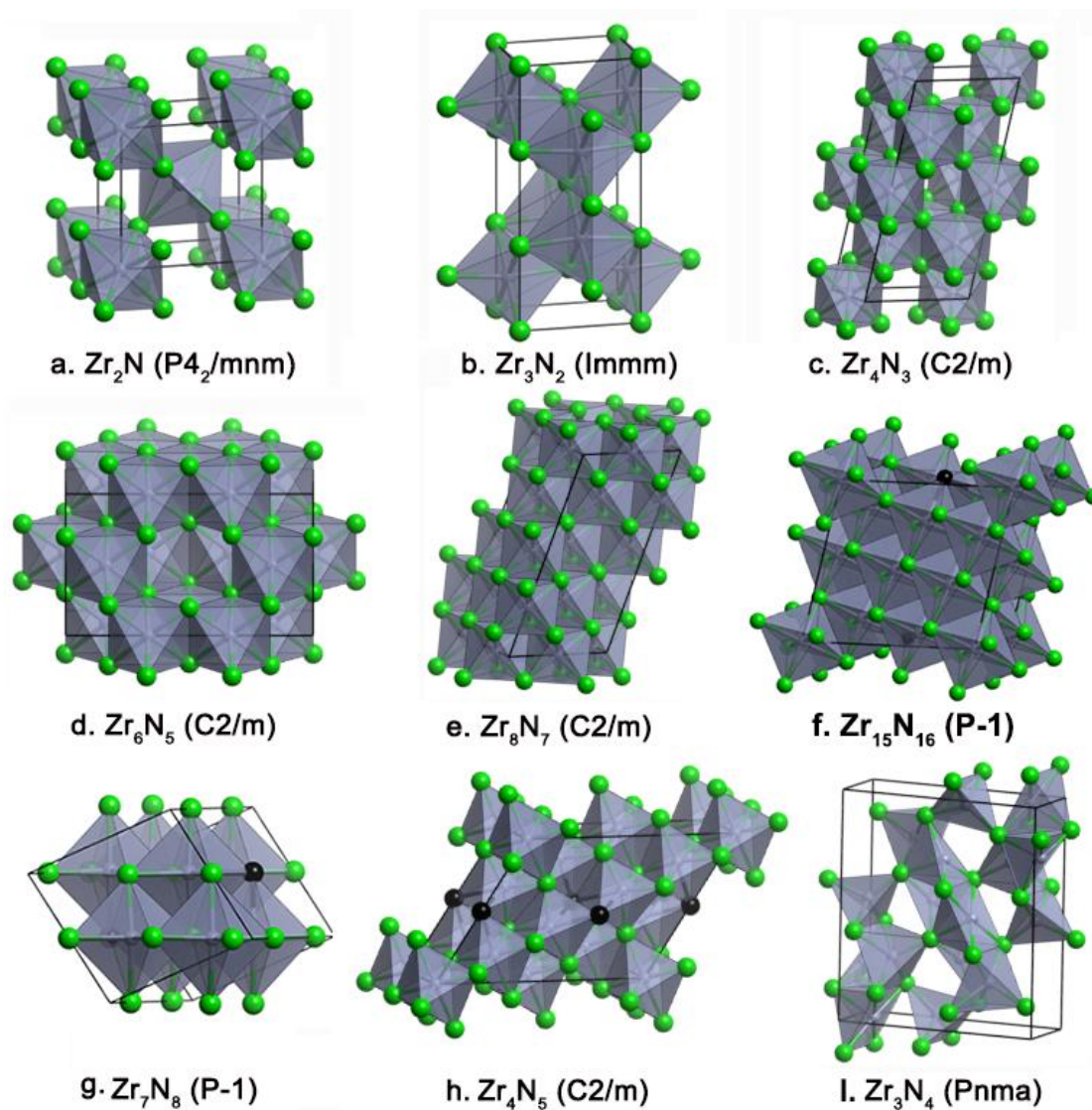


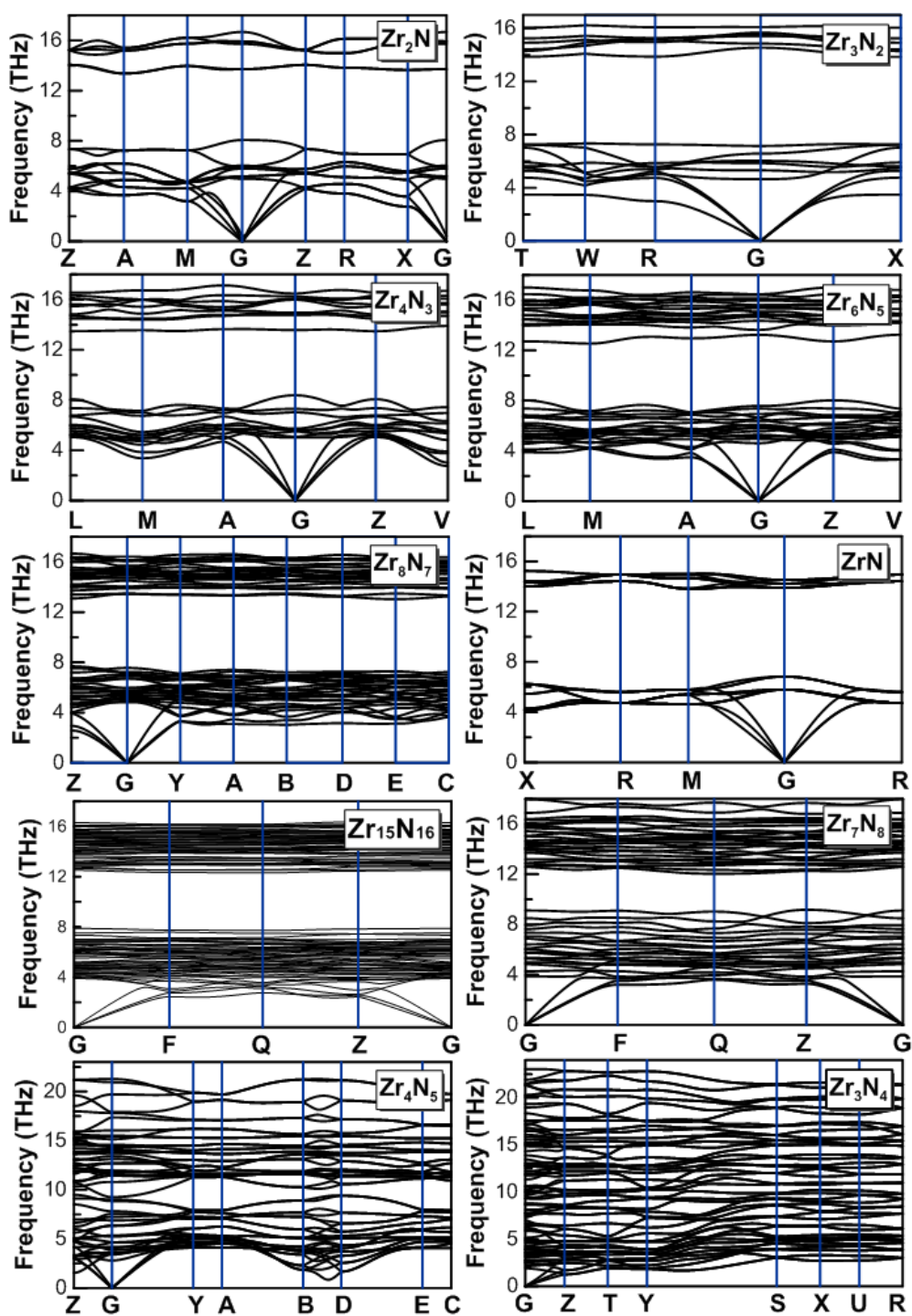
Figure S2 Pressure-driven phase transition in  $Zr_3N_4$ .

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**Figure S3** Crystal structures of the zirconium nitrides. The green spheres denote Zr atoms, while the black ones represent Zr vacancies. Nitrogen atoms are in the centers of the grey octahedra.

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**Figure S4** Calculated phonon dispersion curves of  $Zr_xN_y$  compounds at ambient pressure. All of the dispersion curves have no imaginary phonon frequencies, indicating that they are all dynamically stable.

## Supplementary Information

### Vibrations and thermodynamics of $Zr_xN_y$ compounds at ambient pressure

#### A. Free energies of $Zr_xN_y$ solids

The free energy of a solid, can be expressed as:

$$G = E_0 + pV + F_{vib} + F_{conf} \quad (1)$$

where  $E_0$  is the internal lattice energy,  $p$  is the external pressure,  $V$  is the volume,  $F_{vib}$  is the vibrational free energy, and  $F_{conf}$  denotes configurational entropy contribution. Vibrational contributions represent the phonon energy, which can be treated in the harmonic approximation, by summing over the normal modes frequencies,

$$F_{vib} = \frac{1}{2} \sum_{i=1}^n h\nu_i \left[ \frac{1}{2} + \frac{1}{\exp(-h\nu_i / kT) - 1} \right] \quad (2)$$

where  $k$  is Boltzmann constant,  $h$  is Planck constant.

$$F_{conf} = TS_{conf} = T \cdot k \ln \frac{N!}{n!(N-n)!} \quad (3)$$

where  $S_{conf}$  is the configuration entropy,  $k$  is Boltzmann's constant

#### B. Free energies of $N_2$ gas

The free energy of diatomic  $N_2$  gas molecule can be expressed as:

$$G(N_2) = U(N_2) + \frac{7}{2} k_B T - TS(N_2) \quad (4)$$

where the entropy  $S$  can be divided into three parts, namely, translation entropy  $S_{trans}$ , rotational entropy  $S_{rot}$ , vibrational entropy  $S_{vib}$  and electronic entropy  $S_{elec}$ . Again,  $S_{elec}$  is neglected, while the other terms can be defined as,

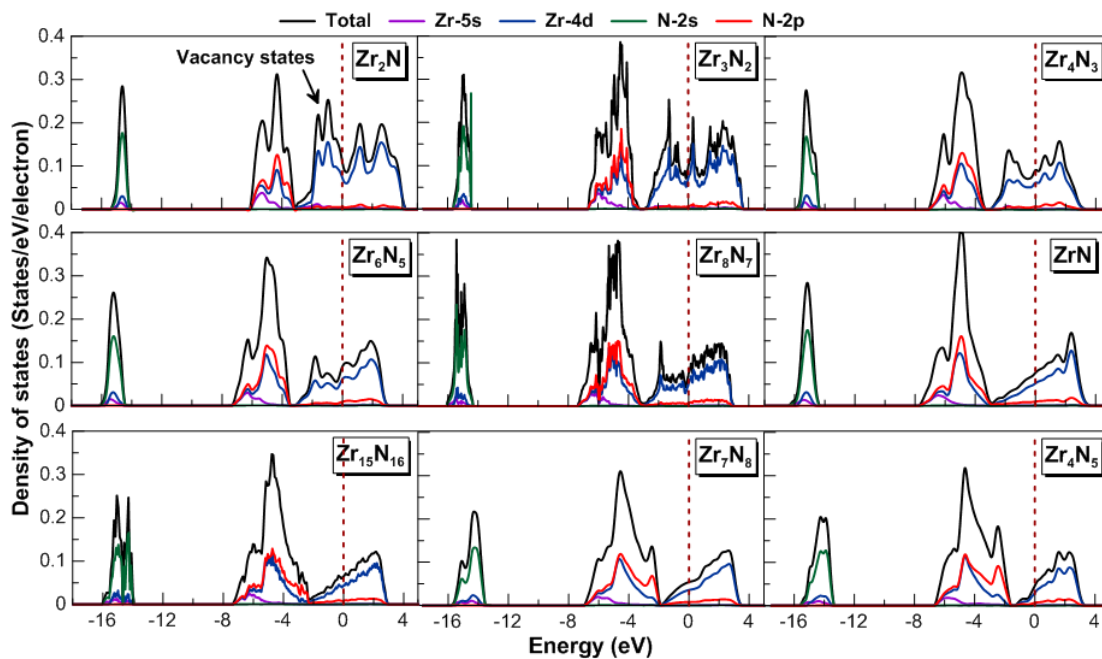
$$S_{trans} = \frac{5}{2} Nk + Nk \ln \left[ \frac{V}{N} \left( \frac{2\pi mkT}{h^2} \right)^{3/2} \right] \quad (5)$$

$$S_{rot} = Nk + Nk \ln \frac{T}{2\Theta_{rot}} \quad (6)$$

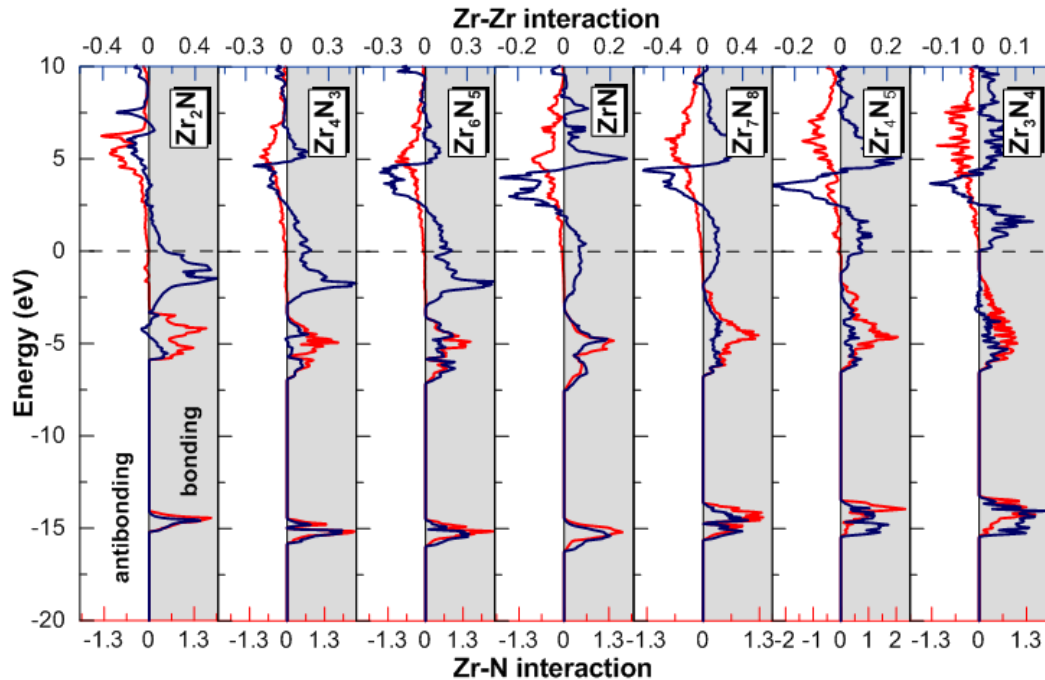
$$S_{vib} = \frac{Nk\Theta_{vib} \left( \frac{1}{2} + \frac{1}{\exp(\Theta_{vib}/T) - 1} \right)}{T} + Nk \ln \left( \frac{\exp(-\Theta_{vib}/T)}{1 - \exp(-\Theta_{vib}/T)} \right) \quad (7)$$

Here,  $\Theta_{rot}$  and  $\Theta_{vib}$  are the characteristic rotational and vibrational temperatures, respectively. To be consistent with the free energy calculation of solids, here we calculated the entropy of  $N_2$  within the same functional.

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**Figure S5** Calculated electronic density of states of zirconium nitrides at ambient conditions.



**Figure S6** Calculated crystal orbital Hamilton population (-COHP) of zirconium nitrides at ambient conditions.

## Supplementary Information

**Table S1 Lattice parameters (Å, °), space groups, atomic coordinates, enthalpies and zero-point energies (ZPE, eV/atom) of Zr<sub>x</sub>N<sub>y</sub> compounds at 0 GPa.**

Phase	Space group	Lattice parameter	Wyckoff position	H ZPE
Zr <sub>2</sub> N	<i>P4<sub>2</sub>/mnm</i>	a=5.426, c=3.319	Zr (0.704, 0.296, 1) N (0.5, 0.5, 0.5)	-9.681 0.052
Zr <sub>3</sub> N <sub>2</sub>	<i>Immm</i>	a=9.921, b=3.292, c=4.542	Zr (0.319, 0, 1); (0, 0, 1) N (0.337, 0, 0.5)	-9.872 0.057
Zr <sub>4</sub> N <sub>3</sub>	<i>C2/m</i>	a=10.836, b=3.267 c=5.645, β=100.7	Zr (0.127, 0, 0.883); (0.132, 0.5, 0.368) N (0.751, 0.5, 0.249); (0, 0.5, 0)	-9.962 0.061
Zr <sub>6</sub> N <sub>5</sub>	<i>C2/m</i>	a=5.659, b=9.763 c=5.654, β=109.5	Zr (0.756, 0.673, 0.253); (0.759, 0.5, 0.742) N (0.5, 0.165, 0); (0.5, 0.334, 0.5); (0, 0.5, 0.5)	-10.037 0.062
Zr <sub>8</sub> N <sub>7</sub>	<i>C2/m</i>	a=11.302, b=6.534 c=5.639, β=109.8	Zr (0.372, 0.256, 0.243); (0.628, 0, 0.241); (0.875, 0, 0.751) N (0.249, 0, 0.999); (0.25, 0.25, 0.5); (0, 0, 0.5); (0, 0.249, 1)	-10.064 0.062
ZrN	<i>Fm-3m</i>	a=4.608 4.58 <sup>[1,2]</sup> , 4.57 <sup>[3]</sup>	Zr (0, 0, 0) N (0.5, 0.5, 0.5)	-10.132 0.061
Zr <sub>15</sub> N <sub>16</sub>	<i>P-1</i>	a=5.640, b=8.590 c=8.584, α=106.3 β=90.0, γ=77.3	Zr (0.189, 0.439, 0.126); (0.685, 0.938, 0.133); (0.622, 0.635, 0.257); (0.144, 0.124, 0.256); (0.569, 0.306, 0.378); (0.06, 0.82, 0.378); (0, 0.5, 0.5); (0.246, 0.749, 0.994); N (0.370, 0.871, 0.241); (0.442, 0.188, 0.123); (0.311, 0.563, 0.375); (0.174, 0.935, 0.63); (0.246, 0.254, 0.5); (0, 1, 0); (0.5, 0.5, 0); (0.062, 0.312, 0.874); (0.122, 0.624, 0.747)	-10.092 0.063
Zr <sub>7</sub> N <sub>8</sub>	<i>P-1</i>	a=6.499, b=5.617 c=5.601, α=80.3 β=73.1, γ=90.0	Zr (0.867, 0.749, 0.520); (0.258, 0.999, 0.996); (0.377, 0.736, 0.507); (0, 0.5, 0) N (0.065, 0.877, 0.751); (0.563, 0.891, 0.742); (0.305, 0.365, 0.741); (0.174, 0.630, 0.262)	-10.046 0.065
Zr <sub>4</sub> N <sub>5</sub>	<i>C2/m</i>	a=10.820, b=3.239 c=8.158, β=122.2	Zr (0.771, 0, 0.684); (0.913, 0.5, 0.115) N (0.121, 0.5, 0.379); (0.302, 0.5, 0.189); (0, 0, 0)	-9.969 0.064
Zr <sub>3</sub> N <sub>4</sub>	<i>Pnma</i>	a=9.875, b=3.297 c=10.904 9.79, 3.30, 10.88 <sup>[2]</sup> 9.73, 3.28, 10.82 <sup>[4]</sup>	Zr (0.772, 0.75, 0.648); (0.892, 0.75, 0.907); (0.930, 0.75, 0.377) N (0.103, 0.75, 0.970); (0.494, 0.75, 0.775); (0.643, 0.75, 0.018); (0.764, 0.25, 0.787)	-9.893 0.068
Zr <sub>3</sub> N <sub>4</sub>	<i>I-43d</i>	a=6.784 6.740 <sup>[5]</sup> , 6.786 <sup>[6]</sup>	Zr (0.375, 0, 0.25) N (0.063, 0.063, 0.063)	-9.875 0.063

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### The directional dependence of the Young's modulus

$$\begin{aligned}
 \frac{1}{E} = & l_1^4 s_{11} + 2l_1^2 l_2^2 s_{12} + 2l_1^2 l_3^2 s_{13} + 2l_1^2 l_2 l_3 s_{14} + 2l_1^3 l_3 s_{15} + 2l_1^3 l_2 s_{16} \\
 & + l_2^4 s_{22} + 2l_2^2 l_3^2 s_{23} + 2l_2^3 l_3 s_{24} + 2l_1 l_2^2 l_3 s_{25} + 2l_1 l_2^3 s_{26} \\
 & + l_3^4 s_{33} + 2l_2 l_3^3 s_{34} + 2l_1 l_3^3 s_{35} + 2l_1 l_2 l_3^2 s_{36} \\
 & + l_2^2 l_3^2 s_{44} + 2l_1 l_2 l_3^2 s_{45} + 2l_1 l_2^2 l_3 s_{46} \\
 & + l_1^2 l_3^2 s_{55} + 2l_1^2 l_2 l_3 s_{56} \\
 & + l_1^2 l_2^2 s_{66}
 \end{aligned} \tag{8}$$

where  $S_{ij}$  are the elastic compliance tensor components constants and  $l_1$ ,  $l_2$  and  $l_3$  are the direction cosines.

### Using Šimůnek model to calculate the crystal hardness

The Vickers hardness was further estimated by Šimůnek's method for the vacancy-ordered  $Zr_xN_y$  phases within the following forms:

$$H = \frac{C}{\Omega} n \left[ \prod_{i,j=1}^n b_{ij} s_{ij} \right]^{1/n} e^{-\sigma f_e} \tag{9}$$

$$f_e = 1 - \left[ \frac{k \left( \prod_{i=1}^k e_i \right)^{1/k}}{\sum_{i=1}^k e_i} \right]^2 \tag{10}$$

$$s_{ij} = \frac{\sqrt{e_i e_j}}{n_i n_j d_{ij}} \tag{11}$$

where  $n_i$  and  $n_j$  are coordination numbers of atoms  $i$  and  $j$ , respectively;  $d_{ij}$  is the interatomic distance of atoms  $i$  and  $j$ ; the reference energy  $e_i = Z_i / R_i$ , which represents the potential of the individual atom  $i$  to attract crystal valence charge;  $Z_i$  is the valence electron number of the atom  $i$ , and  $R_i$  is the ionic radius of the atom, here  $R(N) = 0.99^{[7]}$  and  $R(Zr) = 1.21$ ;  $k$  corresponds to the number of different atoms in the system;  $S_{ij}$  is the strength of the individual bond between atoms  $i$  and  $j$ ;  $b_{ij}$  is the number of interatomic bonds between atoms  $i$  and  $j$  in the unit cell;  $\Omega$  is the volume of the cell; The semi-empirical constant  $C$  equals to 1450;  $e^{-\sigma f_e}$  modifies the effect of the difference between reference energies ( $e_i$ ,  $e_j$ ) and  $\sigma$  equals to 2.8.



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### References

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