Supporting Information Structure and stability of CaH₂ surfaces: on the possibility of electron-rich surfaces in metal hydrides for catalysis

Phuong Vu Ong,*,† Lewis E. Johnson,† Hideo Hosono,*,‡ and Peter V. Sushko†

†Physical Sciences Division, Physical & Computational Sciences Directorate, Pacific Northwest National Laboratory, Richland, WA 99352, U.S.A.

‡Materials Research Center of Element Strategy, Tokyo Institute of Technology, 4259 Nagatsuta, Midori-ku, Yokohama 226-8053, Japan

E-mail: phuong-vu.ong@pnnl.gov; hosono@msl.titech.ac.jp

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A. Structure of CaH₂

Figure S1 (A), (B) and (C) display different choices of CaH_2 unit cells where *c*-axis parallel to [001], [110], and [111] crystallographic orientations, respectively. These unit-cell structures make it convenient to create surfaces by axial scanning method. Slabs with all possible terminations can then be generated by axial scanning and cleaving the bulk crystal along [001], [110] and [111] crystallographic axes.

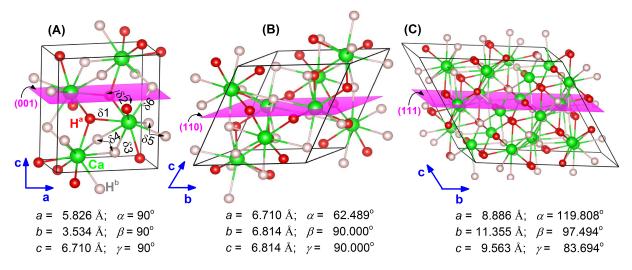


Figure S1: (A), (B) and (C) Different choices of cell of CaH₂ where *c*-axis parallel to [001], [110], and [111] crystallographic orientations, respectively. The shaded area indicate (001), (110), and (111) lattice planes, respectively. Sticks connecting a Ca to its nearest neighbor H atoms. Two sets of symmetrically inequivalent hydrogens are denoted by H^a and H^b. Ca-H^a bond lengths: $\delta 1 = 2.265$ Å, $\delta 2 = 2.225$ Å (two bonds), and $\delta 3 = 2.267$ Å; Ca-H^b bond lengths: $\delta 4 = 2.568$ Å (two bonds), $\delta 5 = 2.461$ Å (two bonds), and $\delta 6 = 2.391$ Å.

B. Doubly-occupied vacancy band on a CaH_2 surface

To demonstrate that a surface vacancy can accommodate an additional electron transferred from a nearby H^- , we constructed (111) slabs with a neutral hydrogen vacancy and a Ne substitution for a surface hydrogen. As Ne is an inert gas element, the extra electron (otherwise transferred to the substituted H to form H^-) will be transferred to the vacancy site,

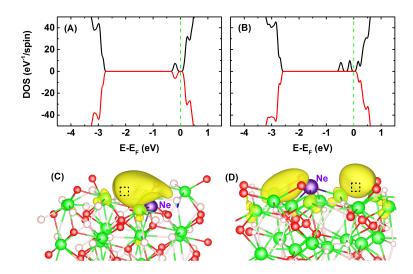


Figure S2: Total density of states (DOS) for a (111) slab containing a neutral hydrogen vacancy and a Ne substitution for a surface hydrogen nearby (A) and further from (B) the vacancy. (C) and (D) surfaces of the constant charge density for electrons occupying the gap states in (A) and (B), respectively. The dotted squares indicate positions of hydrogen vacancies.

form a new confined electron at a different site, or become delocalized. We found that when the Ne ion is close enough to the hydrogen vacancy, the extra electron is localized in the vacancy site and spin-paired with the vacancy-induced electron in the same vacancy band [Figure S2 (A) and (C)]. On the other hand, if the Ne is far enough from the vacancy, the extra electron is localized nearby the Ne and has the same spin as the vacancy-induced one [Figure S2 (B) and (D)]. The same conclusions have been reached when He, instead of Ne, was used as a substitute for the surface hydrogen.

C. Equation for Gibbs free energy of CaH₂ surface and boundary conditions

Gibbs free energy of surface – Consider a CaH₂ slab with N_{Ca}^s and N_{H}^s being the numbers of Ca and H, respectively. The Gibbs free energy of the surface is defined as,

$$\Omega^s = G^s_{slab} - N^s_{\rm Ca} \mu_{\rm Ca} - N^s_{\rm H} \mu_{\rm H} \tag{1}$$

where G_{slab}^s is Gibbs free energy of the CaH₂ slab and μ_{Ca} and μ_{H} are chemical potentials of Ca and H ions in the CaH₂ slab, respectively. If the surface is in equilibrium with bulk CaH_2 reservoir, the sum of the chemical potentials is constrained to be equal to chemical potential μ_{CaH_2} of the bulk CaH_2

$$\mu_{\rm CaH_2} = N^b_{\rm Ca} \mu_{\rm Ca} + N^b_{\rm H} \mu_{\rm H} \tag{2}$$

where N_{Ca}^b and N_{H}^b are the numbers of Ca and H, respectively, in the bulk. Therefore, using equation (2) we can eliminate μ_{Ca} in equation (1). Then the expression of Ω becomes

$$\Omega^s = G^s_{slab} - \frac{N^s_{\text{Ca}}}{N^b_{\text{Ca}}} \mu_{\text{CaH}_2} - \left(N^s_{\text{H}} - N^s_{\text{Ca}} \frac{N^b_{\text{H}}}{N^b_{\text{Ca}}}\right) \mu_{\text{H}}$$
(3)

 Put

$$\Gamma^s_{\mathrm{Ca,H}} = N^s_{\mathrm{H}} - N^s_{\mathrm{Ca}} \frac{N^b_{\mathrm{H}}}{N^b_{\mathrm{Ca}}}$$

$$\tag{4}$$

A positive or negative value of $\Gamma_{Ca,H}^s$ indicates an excess or deficit, respectively, of H ions with respect to Ca ones in the CaH₂ slab. Therefore, a surface corresponding to positive, negative, or zero $\Gamma_{Ca,H}^s$ is said to be hydrogen-rich, hydrogen-poor, or stoichiometric, respectively. Substituting $\Gamma_{Ca,H}^v$ into the equation (3), we obtain

$$\Omega^s = G^s_{slab} - \frac{N^s_{Ca}}{N^b_{Ca}} \mu_{CaH_2} - \Gamma^s_{Ca,H} \mu_H$$
(5)

where we define

$$\mu_{\rm H} = \frac{1}{2} E({\rm H}_2) + \Delta \mu_{\rm H} \tag{6}$$

where $\Delta \mu_{\rm H}$ is variation of chemical potential of hydrogen atom with respect to half an energy of the molecular hydrogen $E({\rm H}_2)$. Substituting equation (6) into (5), we obtain

$$\Omega^s = \phi^s - \Gamma^s_{\mathrm{H,Ca}} \Delta \mu_{\mathrm{H}} \tag{7}$$

where

$$\phi^{s} = G^{s}_{slab} - \frac{N^{s}_{Ca}}{N^{b}_{Ca}} \mu_{CaH_{2}} - \frac{1}{2} \Gamma^{s}_{Ca,H} E(H_{2})$$
(8)

The chemical potential μ_{CaH_2} is nothing but the molar Gibbs free energy of the bulk CaH₂

$$\mu_{\rm CaH_2} = g^b_{\rm CaH_2} \tag{9}$$

Therefore, equation (8) can be rewritten,

$$\phi^{s} = G^{s}_{slab} - \frac{N^{s}_{Ca}}{N^{b}_{Ca}} g^{b}_{CaH_{2}} - \frac{1}{2} \Gamma^{s}_{Ca,H} E(H_{2})$$
(10)

In general, we have

$$g_{\rm CaH_2}^b = E_{\rm CaH_2}^{static} + E_{\rm CaH_2}^{phonon} - TS_{\rm CaH_2} + pV_{\rm CaH_2}$$
(11)

where $E_{\text{CaH}_2}^{static}$ and $E_{\text{CaH}_2}^{phonon}$ are, respectively, static and vibrational components of the crystal energy; S_{CaH_2} is entropy; and V_{CaH_2} is volume; T and p are applied temperature and pressure of the system. Throughout the present work, we approximate the Gibbs free energy to the total energy obtained from *ab initio* calculation, *i.e.*,

$$g^b_{\rm CaH_2} \approx E^{tot}_{\rm CaH_2} \tag{12}$$

Similarly,

$$G_{slab}^s \approx E_{slab}^{tot} \tag{13}$$

Our estimation indicates that total contribution of the vibrational, entropic and mechanical terms is negligible. The above approximations result in an error of the order of 10^{-2} J/m^2 in CaH₂ surface energy.

Boundary conditions – To make the CaH_2 phase stable, precipitation of Ca metal and condensation of hydrogen ions into hydrogen gas should be prevented. These conditions can

be expressed through relations of chemical potentials of the atomic constituents of CaH_2 , *i.e.*, μ_{Ca} and μ_{H} , and the Gibbs free energies (*i.e.*, molar chemical potentials) of the undesired precipitated substances:

(i)
$$2\mu_{\rm H} \le g_{\rm H_2}^b$$
 (14)

(ii)
$$\mu_{Ca} \le g^b_{Ca}$$
 (15)

From equation (6) and the approximations in equations (12) and (13), the condition (i) is equivalent to

(i)
$$\Delta \mu_{\rm H} \le 0$$
 (16)

The upper bound $\Delta \mu_{\rm H} = 0$ is called hydrogen-rich limit.

Similarly, we have

(ii)
$$\Delta \mu_{\rm Ca} \le 0$$
 (17)

From equation (2), we have

$$g_{\rm CaH_2}^b = N_{\rm Ca}^b (g_{\rm Ca}^b + \Delta \mu_{\rm Ca}) + N_{\rm H}^b (\frac{1}{2}E({\rm H}_2) + \Delta \mu_{\rm H})$$
(18)

or

$$\Delta g^b_{\mathrm{CaH}_2} = N^b_{\mathrm{Ca}} \Delta \mu_{\mathrm{Ca}} + N^b_{\mathrm{H}} \Delta \mu_{\mathrm{H}} \tag{19}$$

where

$$\Delta g^{b}_{\rm CaH_2} = g^{b}_{\rm CaH_2} - N^{b}_{\rm Ca} g^{b}_{\rm Ca} - \frac{1}{2} N^{b}_{\rm H} E({\rm H}_2)$$
(20)

From the approximation in (12)

$$\Delta g^b_{\text{CaH}_2} \approx E^{tot}_{\text{CaH}_2} - N^b_{\text{Ca}} E^{tot}_{\text{Ca}} - \frac{1}{2} N^b_{\text{H}} E(\text{H}_2)$$
(21)

where E_{Ca}^{tot} is total energy per Ca atom.

From equations (17) and (19), we have

(ii)
$$\Delta \mu_{\rm H} \ge \frac{1}{N_{\rm H}^b} \Delta g_{\rm CaH_2}^b$$
 (22)

The lower bound $\Delta \mu_{\rm H} = \Delta g^b_{\rm CaH_2}/N^b_{\rm H}$ is called the hydrogen-poor limit. From *ab initio* calculations, $E^{tot}_{\rm Ca} = -2.000$ eV and $E({\rm H}_2) = -6.770$ eV. For a 8×8 bulk supercell, in which $N^b_{\rm Ca} = 32$ and $N^b_{\rm H} = 64$, total energy calculation gave $E^{tot}_{\rm CaH_2} = -334.402$ eV. Therefore, the hydrogen-poor limit is $\Delta \mu_{\rm H} = -0.840$ eV. From (16) and the calculation above, the range of $\Delta \mu_{\rm H}$ is

$$-0.840 \le \Delta \mu_{\rm H} \le 0 \quad (\text{in eV}) \tag{23}$$

D. Temperature and pressure dependence of chemical potential of hydrogen in CaH_2 slabs

In this part, we present the calculation of variation of chemical potential of hydrogen, $\Delta \mu_{\rm H}$, in the CaH₂ slab as a function of temperature and pressure.

We start from the Maxwell's relation

$$\left(\frac{\partial\mu_{\rm H}}{\partial p_{\rm H_2}}\right)_T = \frac{k_B T}{2} \frac{1}{p_{\rm H_2}} \tag{24}$$

Its solution takes the form

$$\mu_{\rm H}(T, p_{\rm H_2}) = \mu_{\rm H}(T, p_{\rm H_2}^{\rm o}) + \frac{k_B T}{2} ln \frac{p_{\rm H_2}}{p_{\rm H_2}^{\rm o}}$$
(25)

where standard state pressure $p_{\rm H_2}^{\rm o} = 0.1$ MPa and k_B is Boltzmann's constant. We rewrite equation (25) as

$$\mu_{\rm H}(T, p_{\rm H_2}) = \mu_{\rm H}(0{\rm K}, p_{\rm H_2}^{\rm o}) + \left[\mu_{\rm H}(T, p_{\rm H_2}^{\rm o}) - \mu_{\rm H}(0{\rm K}, p_{\rm H_2}^{\rm o})\right] + \frac{k_B T}{2} ln \frac{p_{\rm H_2}}{p_{\rm H_2}^{\rm o}}$$
(26)

Oxygen in the CaH_2 slab is in equilibrium with H_2 gas. Therefore,

$$\mu_{\rm H}(T, p_{\rm H_2}) = \frac{1}{2} \mu_{\rm H_2}(T, p_{\rm H_2}) = \frac{1}{2} G^m_{\rm H_2}(T, p_{\rm H_2})$$
(27)

where $G_{H_2}^m(T, p_{H_2})$ is molar Gibbs free energy of H_2 gas. Then the terms enclosed in [..] in equation (26) can be rewritten

$$\mu_{\rm H}(T, p_{\rm H_2}^{\rm o}) - \mu_{\rm H}(0{\rm K}, p_{\rm H_2}^{\rm o}) = \frac{1}{2} \left[G_{\rm H_2}^m(T, p_{\rm H_2}^{\rm o}) - G_{\rm H_2}^m(0{\rm K}, p_{\rm H_2}^{\rm o}) \right]$$
$$\equiv \frac{1}{2} \Delta G_{\rm H_2}^m(T, p_{\rm H_2}^{\rm o})$$
(28)

Using the relation G = H - TS, where H and S are enthalpy and entropy, respectively, we obtain

$$\Delta G_{\mathrm{H}_{2}}^{m}(T, p_{\mathrm{H}_{2}}^{\mathrm{o}}) = \begin{bmatrix} H_{\mathrm{H}_{2}}^{m}(T, p_{\mathrm{H}_{2}}^{\mathrm{o}}) - H_{\mathrm{H}_{2}}^{m}(T_{R}, p_{\mathrm{H}_{2}}^{\mathrm{o}}) \end{bmatrix} \\ - \begin{bmatrix} H_{\mathrm{H}_{2}}^{m}(0\mathrm{K}, p_{\mathrm{H}_{2}}^{\mathrm{o}}) - H_{\mathrm{H}_{2}}^{m}(T_{R}, p_{\mathrm{H}_{2}}^{\mathrm{o}}) \end{bmatrix} \\ - TS_{\mathrm{H}_{2}}^{m}(T, p_{\mathrm{H}_{2}}^{\mathrm{o}}) \tag{29}$$

Equation (29) can then be separated as

$$A = \left[H_{\rm H_2}^m(T, p_{\rm H_2}^{\rm o}) - H_{\rm H_2}^m(T_R, p_{\rm H_2}^{\rm o}) \right]$$
(30)

$$B = \left[H_{\rm H_2}^m(0\mathrm{K}, p_{\rm H_2}^{\rm o}) - H_{\rm H_2}^m(T_R, p_{\rm H_2}^{\rm o}) \right]$$
(31)

$$C = TS_{H_2}^m(T, p_{H_2}^o)$$
(32)

The A, B, and C values can be determined from thermodynamic tables.¹

From the equation (6), (26), (28), and the approximation $\frac{1}{2}E(H_2) = \mu_H(0K, p_{H_2}^o)$ we obtain the equation determining the variation of chemical potential of oxygen, $\Delta \mu_H$, in

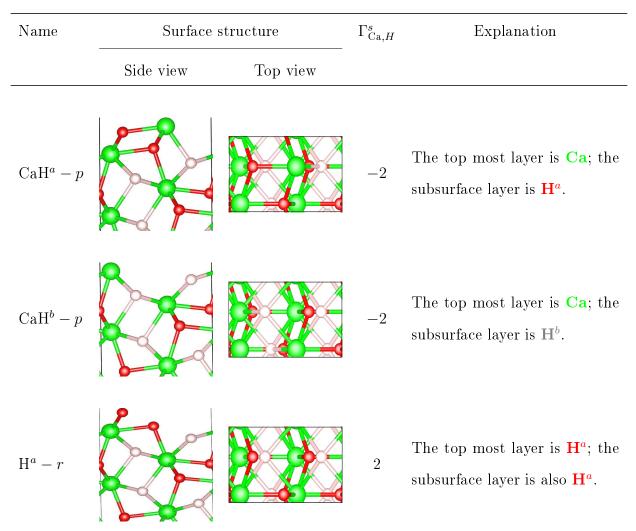
 CaH_2 slab with temperate and oxygen pressure.

$$\Delta \mu_{\rm H}(T, p_{\rm H_2}) = \frac{1}{2} \left[\Delta G^m_{\rm H_2}(T, p^{\rm o}_{\rm H_2}) + k_B T ln \frac{p_{\rm H_2}}{p^{\rm o}_{\rm H_2}} \right]$$
(33)

E. Surface nomenclature

In this section, we explain surface nomenclature and show atomic structures of (001), (110) and (111) surfaces. Symmetric slabs with all possible terminations were generated by axial scanning and cleaving the bulk crystal along [001], [110] and [111] crystallographic axes.

Table S1: Surface nomenclature for (001)CaH₂. Positive, negative, or zero $\Gamma_{Ca,H}^{s}$ indicates a hydrogen-rich (-r), hydrogen-poor (-p), or stoichiometric (-s) surface, respectively.



\dots continued (001)CaH₂

Name	Surface structure		$\Gamma^s_{\mathrm{Ca},H}$	Explanation
	Side view	Top view		
$\mathrm{H}^a - s$			0	The top most layer is \mathbf{H}^{a} ; the subsurface layer is \mathbf{Ca} .
$\mathbf{H}^{b}-r$			2	The top most layer is \mathbf{H}^b ; the subsurface layer is also \mathbf{H}^b .
$\mathbf{H}^{b}-s$			0	The top most layer is \mathbf{H}^b ; the subsurface layer is \mathbf{Ca} .

Table S2: Surface nomenclature for (110)CaH₂. Positive, negative, or zero $\Gamma_{Ca,H}^{s}$ indicates a hydrogen-rich (-r), hydrogen-poor (-p), or stoichiometric (-s) surface, respectively.

Name	Surface s	tructure	$\Gamma^s_{\mathrm{Ca},H}$	Explanation
	Side view	Top view		
Ca-p			-4	The top most layer is Ca .
Ca-s			0	The top most layer is Ca of lower concentration.
$\mathbf{H}^{a}-p$			-2	The top most layer is H ^a ; the subsurface layer is Ca .
$\mathrm{H}^a - r$			4	The top most layer is \mathbf{H}^{a} ; the subsurface layer is \mathbf{H}^{b} .

\ldots continued (110)CaH₂

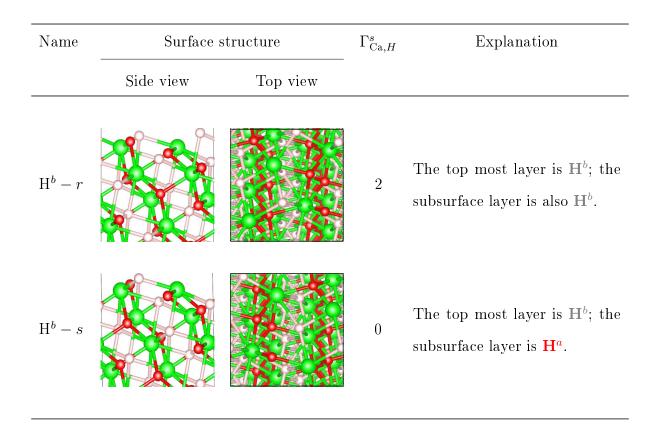


Table S3: Surface nomenclature for (111)CaH₂. Positive, negative, or zero $\Gamma_{Ca,H}^{s}$ indicates a hydrogen-rich (-r), hydrogen-poor (-p), or stoichiometric (-s) surface, respectively.

Name	Surface structure		$\Gamma^s_{\mathrm{Ca},H}$	Explanation
	Side view	Top view		
$CaH^a - p$			-4	The top most layer is Ca ; the subsurface layer is H^a .

\dots continued (111)CaH₂

Name	Surface structure		$\Gamma^s_{\mathrm{Ca},H}$	Explanation
	Side view	Top view		
$\operatorname{CaH}^b - p$			-12	The top most layer is Ca ; th subsurface layer is H^b .
$\mathbf{H}^{a}-p$			-8	The top most layer is H ^a ; th subsurface layer is Ca .
$\mathrm{H}^{a}-s$			0	The top most layer is \mathbf{H}^{a} ; the subsurface layers are \mathbf{H}^{b} , \mathbf{H}^{a} and \mathbf{Ca} , in that order.
$\alpha - \mathbf{H}^a - r$			12	The top most layer is \mathbf{H}^{a} ; the subsurface layers are \mathbf{H}^{b} , \mathbf{H}^{a} and \mathbf{H}^{a} , in that order.

\dots continued (111)CaH₂

Name	Surface s	structure	$\Gamma^s_{{\rm Ca},H}$	Explanation
	Side view	Top view		
$\beta - \mathbf{H}^a - r$			4	The top most layer is \mathbf{H}^{a} ; the subsurface layers are \mathbf{H}^{a} , \mathbf{H}^{a} and \mathbf{H}^{a} , in that order.
$\mathbf{H}^{b} - p$			-4	The top most layer is \mathbf{H}^b ; the subsurface layers are \mathbf{H}^a and \mathbf{Ca} .
$\alpha - \mathbf{H}^b - r$			4	The top most layer is \mathbf{H}^b ; the subsurface layer is \mathbf{Ca} .
$\beta - \mathbf{H}^b - r$			8	The top most layer is \mathbf{H}^b ; the next two subsurface layers as \mathbf{H}^a .

References

JANAF Thermochemical Tables, 3rd edition, by M. W. Chase, Jr., C. A. Davies, I. R. Downey, Jr., D. J. Frurip, R. A. McDonald, and A. N. Syverud, published by Am. Chem. Soc. and Am. Inst. Phys. for NBS (now NIST), Washington, DC (1986).