Supplementary Information

The structural and photophysical properties of multibranched derivatives with curved conjugated aromatic

cores

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| M1-H | | M1-CH3 | | M1-OH | | M1-NH2 | | M1-NO2 | |
|---------|--------|---------|--------|---------|--------|---------|--------|---------|--------|
| 10.0868 | 0.0007 | 8.4932 | 0 | 8.7734 | 0 | 8.0831 | 0.1686 | 6.3775 | 0.0022 |
| 10.2599 | 0.0015 | 8.4932 | 0 | 8.7734 | 0 | 8.5577 | 0.0007 | 6.567 | 0.0034 |
| 12.1128 | 0.0144 | 10.1245 | 0.078 | 11.9306 | 0 | 10.013 | 0.5615 | 8.2071 | 0.1791 |
| 13.1581 | 0.0313 | 10.1245 | 0.078 | 11.9306 | 0 | 10.9035 | 0.1082 | 9.3072 | 0.0392 |
| 14.5848 | 0.0206 | 10.385 | 0 | 12.1798 | 2.7339 | 12.3429 | 0.0216 | 9.6912 | 0.4701 |
| 15.1134 | 0.4128 | 10.385 | 0 | 12.5176 | 0.1508 | 12.6623 | 0.2304 | 10.1984 | 0.5408 |
| 15.4958 | 0.0046 | 10.5442 | 0.1394 | 12.5176 | 0.1508 | 13.2841 | 1.5145 | 10.42 | 0.716 |
| 15.9939 | 0.0256 | 13.1814 | 0.4597 | 15.5351 | 1.6263 | 13.9395 | 0.641 | 10.8752 | 0.4278 |
| 16.8296 | 0.0977 | 14.2989 | 0 | 16.2261 | 0.3563 | 14.7425 | 1.0614 | 12.0297 | 0.0503 |
| 17.3075 | 0.0671 | 14.2989 | 0 | 16.2261 | 0.3563 | 18.1896 | 0.3497 | 14.8405 | 0.0928 |

Table S1 the lowest ten vibrational frequencies and corresponding intensity of M1-R derivatives.

Table S2 the lowest ten vibrational frequencies and corresponding intensity of M2-R derivatives.

| М2-Н | | M2-CH3 | | M2-OH | | M2-NH2 | | M2-NO2 | |
|---------|--------|---------|--------|---------|--------|---------|--------|---------|--------|
| 7.442 | 0 | 5.8985 | 0.0014 | 7.6584 | 0.0255 | 5.0003 | 0.0424 | 6.2333 | 0.0028 |
| 7.442 | 0 | 7.4877 | 0.0066 | 8.4298 | 0.0032 | 6.4209 | 0.0144 | 6.4356 | 0.0066 |
| 9.5652 | 0.0011 | 7.9382 | 0.0115 | 8.9073 | 0.1256 | 7.3094 | 0.281 | 7.1215 | 0.0627 |
| 9.5652 | 0.0011 | 8.5686 | 0.0321 | 9.1503 | 0.0354 | 7.6271 | 0.0736 | 7.483 | 0.115 |
| 10.1108 | 0.0384 | 9.1228 | 0.0061 | 9.4056 | 0.1057 | 8.599 | 0.1126 | 7.6055 | 0.0799 |
| 12.1973 | 0 | 10.3329 | 0.0043 | 10.1473 | 0.0012 | 9.8593 | 0.0176 | 8.065 | 0.0305 |
| 12.1973 | 0 | 10.5854 | 0.0002 | 10.3116 | 0.0267 | 9.9541 | 0.0062 | 8.495 | 0.0209 |
| 23.6729 | 0 | 20.6572 | 0.0586 | 19.76 | 0.5788 | 20.2164 | 0.0053 | 15.9205 | 0.0857 |
| 23.6729 | 0 | 23.2807 | 0.1365 | 22.5245 | 0.565 | 22.6189 | 0.1817 | 17.4966 | 0.0169 |
| 25.2561 | 0.069 | 23.9047 | 0.2582 | 23.3248 | 0.3978 | 23.9697 | 0.9541 | 18.213 | 0.0722 |



Table S3 IR spectrum and corresponding normal vibrational mode of the lowest frequencies of M1-R and M2-R derivatives.





The development of range-separated exchange(RS) density functionals has allowed the mitigation of the CT issue with TDDFT. As reported in recent work, TDDFT methods using conventional hybrid functionals surprisingly fail in describing the low-lying excited states. The description of excitonic properties for CT molecular systems requires the full 100% hartree-fock exchange and a reasonable long-range-separated parameter^{1,2.} In this part, we compared the results (as shown in Table S4) from CAM-B3YP, M06-2x and a range-separated exchange density functionals with optimized range-separation parameter ω and contribution of Hartree-Fock exchange. It is well know that the ω values in RS functionals have been shown to be strongly system-dependent. In present calculation, the choice of parameter ω is based on "Golden proportion" method recommended by Baer, Kronik, and their collaborators³

To facilitate functional tuning we employed a RS hybrid with a three-parameter errorfunction separation of the interelectronic distance r_{12} as follows

$$\frac{1}{r_{12}} = \frac{1 - \left[\alpha + \beta erf(\omega r_{12})\right]}{r_{12}} + \frac{\alpha + \beta erf(\omega r_{12})}{r_{12}}$$

The long-range component of the hybrid functional is given by the second term on the righthand side. The switching from DFT-like to (HF-like) exact exchange (eX) is determined by the range-separation parameter γ . The procedure is designed such that ε_H is as close as possible to the negative IP, which is an exact condition in Kohn–Sham (KS) and generalized KS theory, for both the N-electron and N+ 1 electron systems. The procedure produces H–L energy gaps that are optimally close to the fundamental gap IP–EA of the N-electron system.

$$J^{2} = \sum_{i=0}^{1} \left[\mathcal{E}_{H} \left(N + i \right) + IP \left(N + i \right) \right]^{2}$$

Here, N is the electron number of the system, ϵ_{H} is the HOMO energy, and IP is the ionization potential calculated from self-consistentfield (SCF) energy differences.

The optimal ω values are in the range from 0.126 to 0.129 Bohr⁻¹ for M1-R systems, The optimal ω values of M2-R range from 0.217 to 0.221 Bohr⁻¹. They are thus significantly smaller than the default ω value for LC- ω PBE, CAM-B3LYP and M062x). In order to compare the behavior of various functionals, some excited state datum are listed in Table S2. The untuned RS functionals, such as CAM-B3LYP and M062x, produce slightly greater errors compared to experimental results. the allowed maximum absorption wavelengths (λ_{max}) of M1-NH₂ predicted by CAM-B3LYP and M062x consistently are overestimated by percentages 15.0 and 14.3%, respectively. The predicted λ_{max} using optimal LC- ω PBE functionals is agree with the experiment data, only a small deviation less than 4nm. If Optimal LC- ω PBE results are regarded as a benchmark, similar tendency was also observed in the multibranched corannulene with acceptor ends, M1-NO₂.

| | M1-NH ₂ | M1-NO ₂ | |
|-----------------------------------|--------------------|--------------------|--|
| IP | -6.088098 | -7.468587 | |
| EA | 1.174446 | 2.348647 | |
| Gap | 4.889344 | 5.089620 | |
| J2 | 0.004920 | 0.000112 | |
| ω | 0.1262 | 0.1291 | |
| λ_{max} -LC- ω PBE | 416.8 | 405.5 | |
| λ_{max} -CAM-B3LYP | 355.6 | 347.5 | |
| λ_{max} -M062X | 360.4 | 346.9 | |
| λ_{max} -experiment | 420.0 | - | |

Table S4 Comparison of excited properties predicted using various functionals.

(1) B. M. Wong, T. H. Hsieh, Optoelectronic and Excitonic Properties of Oligoacenes: Substantial Improvements from Range-Separated Time-Dependent Density Functional Theory, J. Chem. Theory Comput. 2010, 6,3704–3712

(2) T. Seidler, K. Stadnicka, B. Champagne, Evaluation of the Linear and Second-Order NLO Properties of Molecular Crystals within the Local Field Theory: Electron Correlation Effects, Choice of XC Functional, ZPVA Contributions, and Impact of the Geometry in the Case of 2-Methyl-4-nitroaniline, J. Chem. Theory Comput.2014, 10, 2114–2124.

(3) R. Baer, E. Livshits, U.Salzner, Tuned Range-Separated Hybrids in Density Functional Theory, Annu. Rev. Phys. Chem.2010, 61,85–109



Fig S1 Molecular orbital diagrams of M1-NH2 and M1-NO2calculated at the B3LYP/6-31+G(d,p) level of theory.

Note: The front molecular orbital diagrams of M1-H, M1-CH3, M1-OH have similiar electron density distribution to that of M1-NH2.



Fig S2 Molecular orbital diagrams of M2-NH2 and M2-NO2calculated at the B3LYP/6-31+G(d,p) level of theory.

Note: The front molecular orbital diagrams of M2-H, M2-CH3, M2-OH have similiar electron density distribution to that of M2-NH2.



Fig S3 Orbital correlation diagram from single branched to multibranched corannulene derivative.