Electronic Supplementary Information (ESI) for

The hydration of the heavy-alkalines Rb^+ and Cs^+ through Molecular Dynamics and EXAFS Spectroscopy: surface clusters and eccentricity

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Intermolecular ion-water potential and MCDHO2 water potential

The monoatomic cation, M(I), (M being Rb or Cs) is described by a positive charge, $Z_{\rm M} = 3$, and a mobile negative charge density, $\rho_{\rm M}$ with a total charge, $q_{\rm M} = -2$, joined to the nucleus by a spring of force constant, $k_{\rm M}$ (see Figure S1). The intra-atomic energy is defined by:

$$U_{intra} = \frac{1}{2}k_{\rm M} \cdot r^2 \tag{1}$$

where r is the distance between the nucleus and its associated mobile charge density. In the absence of an external field, the equilibrium position of the oscillator is located on the nucleus and $U_{intra} = 0$.



Figure S1: Schematic representation of MCDHO model for the water molecule and the metal cation.

The incorporation of the MCDHO2 model to the new potential needs the consideration of the following intermolecular terms for the $M(I)-H_2O$ interaction:

• Classical interaction between the mobile charge densities, $q_{\rm O}$ and $q_{\rm M}$, given by a two-exponential function:

$$U_{inter}(q_{\rm O}, q_{\rm M}) = A_{\rm MO} \cdot e^{-\alpha_{\rm MO} \cdot r_{\rm MO}} + B_{\rm MO} \cdot e^{-\beta_{\rm MO} \cdot r_{\rm MO}}$$
(2)

where $r_{\rm MO}$ is the distance between the mobile charge densities. $A_{\rm MO}$, $\alpha_{\rm MO}$, $B_{\rm MO}$ and $\beta_{\rm MO}$ are fitting parameters.

• Classical interaction between the M nucleus, $Z_{\rm M}$, and those of the water molecule, $Z_i (i \equiv O, H)$, is given by a two-exponential function as well:

$$U_{inter}(Z_i, Z_{\rm M}) = C_{{\rm M}i} \cdot e^{-\gamma_{{\rm M}i} \cdot R_i} + D_{{\rm M}i} \cdot e^{-\delta_{{\rm M}i} \cdot R_i}$$
(3)

where R_i is the distance between the M nucleus and each *i*-th nucleus of the water molecule, and C_{Mi} , γ_{Mi} , D_{Mi} and δ_{Mi} are fitting parameters.

• Electrostatic interaction between the water mobile charge density, $q_{\rm O}$ and the M nucleus, $Z_{\rm M}$:

$$U_{inter}(q_{\rm O}, Z_{\rm M}) = \frac{q_{\rm O} Z_{\rm M}}{r'} \left[1 - \left(\frac{r'}{\lambda'} + 1\right) e^{-2r'/\lambda'} \right]$$
(4)

where r' is the distance between the center of ρ_0 and the M nucleus and λ' is the intermolecular screening described in the original MCDHO2 paper (Villa, A.; Hess, B.; Saint-Martin, H. J. Phys. Chem. B **2009**, *113*, 7270-7281).

• Electrostatic interaction between the M mobile charge density, $q_{\rm M}$, and each of the charges on the water molecule nuclei, $Z_i (i \equiv O, H)$:

$$U_{inter}(Z_i, q_{\rm M}) = \frac{Z_i q_{\rm M}}{r_i} \left[1 - \left(\frac{r_i}{\lambda'_{\rm M}} + 1\right) e^{-2r_i/\lambda'_{\rm M}} \right]$$
(5)

where r_i is the distance from the $\rho_{\rm M}$ center to Z_i and $\lambda'_{\rm M}$ is the corresponding intermolecular screening.

Thus, the interaction energy for a cluster with N water molecules is computed by the expression:

$$U = \sum_{S=1}^{N} \left(\sum_{i \in S} \sum_{j \in T} [U_{inter}(Z_i, Z_j) + U_{inter}(q_i, q_j) + U_{inter}(q_i, Z_j) + U_{inter}(q_j, Z_i)] + \sum_{i \in S} \frac{1}{2} k_i \cdot r_{ii}^2 + \frac{1}{2} k_M \cdot r^2 \right)$$
(6)

where S runs over the water molecules and T over the M.

The MCDHO2 water molecules are described by a positive charge in the hydrogens $Z_{\rm H} = 0.62$ and in the oxygen $Z_{\rm O} = 2.00$ and with a negative charge in the mobile charge density $q_{\rm M} = -3.92$ joined to the oxygen by a spring of force constant, k(see Figure S1):

$$U_k = \frac{1}{2}k \cdot r^2 \tag{7}$$

The charge density is modelled by λ :

$$\rho(r) = \frac{q}{\pi \lambda^3} e^{-2\frac{r}{\lambda}} \tag{8}$$

being the charge:

$$q = 4\pi \int_0^\infty \rho(r) r^2 dr \tag{9}$$

The interatomic distance is modelled by a Morse potential:

$$U_{d_{\rm OH}} = D_{\rm OH} \left(e^{-2\gamma (R_{\beta} - r_{\rm e})} - 2e^{-\gamma (R_{\beta} - r_{\rm e})} \right)$$
(10)

The internal angle is defined by a quartic potential:

$$U_{\Theta_{\text{HOH}}} = a_1(\Theta - \Theta_e) + a_2(\Theta - \Theta_e)^2 + a_3(\Theta - \Theta_e)^3 + a_4(\Theta - \Theta_e)^4$$
(11)

Being the internal energy of a water molecule:

$$U_{\text{internal}} = \frac{1}{2}kr_{\text{O}}^{2} + \frac{Z_{\text{H}}^{2}}{R_{1,2}} + \frac{qZ_{\text{H}}}{r_{\beta}} \left[1 - \left(\frac{r_{\beta}}{\lambda} + 1\right)e^{-2r_{\beta}/\lambda} \right] + U_{k} + U_{d_{\text{OH}}} + U_{\Theta_{\text{HOH}}}$$
(12)

The energy of a cluster of N water molecules is defined by a Lennard-Jones potential for the interactions between oxygens, between hydrogens and between oxygen and hydrogen.

$$U_{\text{total}} = \sum_{n=1}^{N} \sum_{m=1}^{n-1} \left(\frac{A}{r_{nm}}\right)^{12} - \left(\frac{B}{r_{nm}}\right)^{6} + \frac{q^{2}}{r_{nm}} + \frac{qZ_{\beta}}{r_{n\beta}} \left[1 - \left(\frac{r_{n\beta}}{\lambda} + 1\right)e^{-2r_{n\beta}/\lambda}\right] + \sum_{\beta \in m}^{q} \left(\frac{A_{\alpha\beta}}{r_{\alpha\beta}}\right)^{12} - \left(\frac{B_{\alpha\beta}}{r_{\beta\alpha}}\right)^{6} + \frac{Z_{\alpha}\beta}{r_{\alpha\beta}}$$
(13)

	Rb	Cs
$k_{\rm M}$	0.217025	0.364708
$\lambda'_{\rm M}$	0.621742	0.995166
A _{MO}	539.102465	449.494649
$\alpha_{\rm MO}$	1.357374	1.413723
$B_{\rm MO}$	-383.448443	-362.094604
$\beta_{\rm MO}$	1.293485	1.371814
$C_{\rm MH}$	1857.808870	1776.523448
$\gamma_{\rm MH}$	0.857618	0.834310
$D_{\rm MH}$	-1859.064561	-1776.616635
$\delta_{\rm MH}$	0.857790	0.834339

Table S1: Fitted parameters of the M⁺-H₂O potentials (in a.u.).

$Z_{\rm H}$	0.62
Zo	2.0
q	-3.24
k	1.00
λ	1.90
D _{OH}	0.42954902
r_e	1.3440633
γ	1.1131102
θ_e	1.927
a_1	0.031621
a_2	0.043914
a_3	-0,012721
a_4	-0.00866
AOM	3.228656
B _{OM}	1.962046
A _{OH}	2.037891
$A_{\rm HH}$	0

Table S2: Parameters of the MCDHO2 water potential (in a.u)

Details of MD simulations

- 1 cation + 1000 water molecules.
- \bullet Cubix box with length chosen to get experimental water density= 0.997 g $\rm cm^{-3}$ at simulation conditions.
- NVT ensemble with T=300K
- PBC and Ewald sum.
- MD simulations have been carried out using the Dynamical shell model (P.J.Mitchell and D. Fincham, *J. Phys. Condens. Matter* **1993**, 5, 1031-1038) to account for the polarizable MCDHO2 model.
- A modified version of the DL-POLY Classic code (W. Smith, T.R. Forester and I.T. Todorov, DL-POLY v. 2.19; STFC Daresbury Laboratory, Daresbury, Warrington WA4 4AD, Cheshire, UK 2012).
- 1ns production time.

Quantum Mechanical Potential Energy Surface

The interaction energies to be fitted are obtained from ab initio calculations obtained at the level indicated in the following table.

Octahydrates



Figure S2: Quantum-mechanical optimized structures for $[{\rm Rb}({\rm H_2O})_n]^+$

Table S3: Quantum mechanical methods^a

M(I)	DFT	$Pseudopotential^b$
\mathbf{Rb}	M062x	ECP28MBW
\mathbf{Cs}	M062x	ECP46MBW

 $^{(a)}$ Basis sets from:

M. Dolg et al. Theor.Chem.Acc. **1989**, 75, 173-194; **1993**, 85, 441-450 J. Yang et al. Theor.Chem.Acc. **2005**, 113, 212-224. ^(b) Pseudopotentials taken from http://www.theochem.uni-stuggart.de



Figure S3: Fitting of the exchangeable HI Rb^+ - H_2O interaction potential to the set of QM structures.



Figure S4: Fitting of the exchangeable HI Cs^+ - H_2O interaction potential to the set of QM structures.

Example of EXAFS inputs used in the FEFF simulation

TITLE_Rb_structure_snaps	shot_NVT		
EDGE K S02 1.0			
CONTROL 1 0 0 0 0 0 PRINT 0 0 0 3 0 0			
COREHOLE RPA EXAFS 16.0 CRITERIA 4.0 2.5 RPATH 6.0 NLEG 4			
RPATH 6.0 TDLDA 1 SCF 6.0 EXCHANGE 0 -5.0 0. POTENTIALS 0 37 Rb 3 3 1 8 0 3 3 2 1 H 2 2			
ATOMS 0.000000 0.000000 2.5040010 -1.0831000 2.6410000 -1.896100 3.2770000 -1.223100 -1.4640000 -0.174700	0 0.000000 0 -0.4564000 0 -0.9654000 0 0.1245000 0 -2.4274000	0 Rb 1 0 2 H 2 H 1 0	0.0000000 2.7661214 3.3914707 3.5000290 2.8400857
END			

Figure S5: Part I of the input used on the EXAFS simulation

```
TITLE_Rb_structure_snapshot_NVT

EDGE K

S02 1.0

CONTROL 0 1 1 1 1 1

PRINT 0 0 0 3 0 0

COREHOLE RPA

EXAF5 16.0

CRITERIA 4.0 2.5

RPATH 6.0

NLEG 4

RPATH 6.0

TDLDA 1

SCF 6.0

EXCHANGE 0 -5.0 0.

POTENTIALS

0 37 Rb 3 3

1 8 0 3 3

ATOMS

0.0000000 0.0000000 0.0000000 0 Rb 0.0000000

2.5040010 -1.0831000 -0.4564000 1 0 2.7661214

-1.4640000 -0.1747000 -2.4274000 1 0 2.8400857

.

END
```

Figure S6: Part II of the input used on the EXAFS simulation

For the EXAFS simulation of the K-edge of Rubidium $\Delta E_0 = -5.0$ eV whereas for the L3-edge of Cesium $\Delta E_0 = -2.0$ eV.

Example of XANES inputs used in the FEFF simulation

TITLE_Rb_structure_snapshot_NVT
EDGE K
CONTROL 1 0 0 0 0 0 PRINT 2 2 0 0 0 0
XANES FMS 6.0 1 AFOLP OPCONS MPSE 2
COREHOLE RPA
TDLDA 1 SCF 6.0 0
EXCHANGE 0 0.0 -2.0 2 POTENTIAL 0 37 Rb 3 3 1 8 0 3 3 2 1 H 2 2
ATOMS 0.0000000 0.0000000 0.0000000 0 Rb 0.0000000 2.5040010 -1.0831000 -0.4564000 1 0 2.7661214 2.6410000 -1.8961000 -0.9654000 2 H 3.3914707 3.2770000 -1.2231000 0.1245000 2 H 3.5000290 -1.4640000 -0.1747000 -2.4274000 1 0 2.8400857
END

Figure S7: Part I of the input used on the XANES simulation

TITLE_Rb_structure_snapshot_NVT EDGE K CONTROL 0 1 1 1 1 1 PRINT 2 2 0 0 0 0 XANES FMS 6.0 1 AFOLP OPCONS MPSE 2 COREHOLE RPA TDLDA 1 SCF 6.0 0 EXCHANGE 0 0.0 -2.0 2 POTENTIALS 0 37 Rb 3 3 1 8 0 3 3 ATOMS 0.0000000 0.0000000 0 Rb 0.0000000 2.5040010 -1.0831000 -0.4564000 1 0 2.7661214 -1.4640000 -0.1747000 -2.4274000 1 0 2.8400857 ... END

Figure S8: Part II of the input used on the XANES simulation

The XANES simulation of the K-edge of Rubidium uses a broadening of -2.0 eV whereas the simulation of the L3-edge of Cesium uses a broadening of 0.9 eV.

Equation and parameters of Ohta's equation for multi-electron excitation removing from EXAFS spectrum.

The damped sine equation of Ohta et. al :

$$f_{\chi}(k) = (Ak^2 + B)exp(-Ck^2)sin(Dk^2 + Ek + F) + G + Hk$$
 (14)

Coefficient	$\rm Rb^+$	Cs^+
А	0.194812	0.202138
В	0.281428	0.103479
\mathbf{C}	0.0661385	0.0804508
D	0.100565	-0.00237742
\mathbf{E}	3.4741	4.82325
\mathbf{F}	2.24634	-3.3454
G	0.05852	-0.000956909
Η	-0.0243885	0.000416375

Table S4: Damped sine coefficients



Figure S9: Left: Simulated EXAFS spectrum of Cs^+ obtained with a different number of snapshots evenly taken from trajectory, (500, 250, 125 and 25), and spectrum derived from the average of the first 25 snapshots of the trajectory. Right: Comparison of the average EXAFS spectrum with five different spectra derived from different snapshots.



Figure S10: Left: Simulated XANES spectrum of Cs^+ obtained with a different number of snapshots evenly taken from trajectory, (500, 250, 125 and 25), and spectrum derived from the average of the first 25 snapshots of the trajectory. Right: Comparison of the average XANES spectrum with five different spectra derived from different snapshots.