

## Supplementary Information

### An integrated microwell array platform for cell lasing analysis

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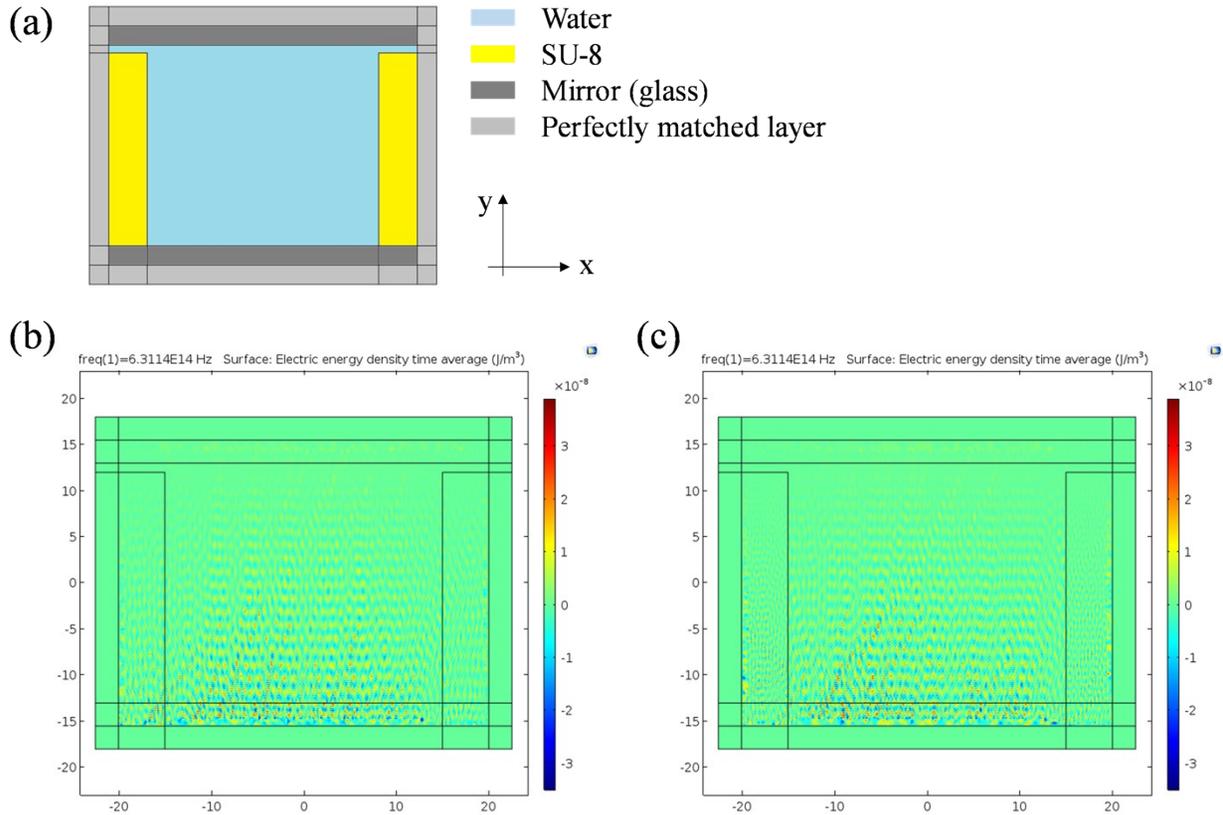
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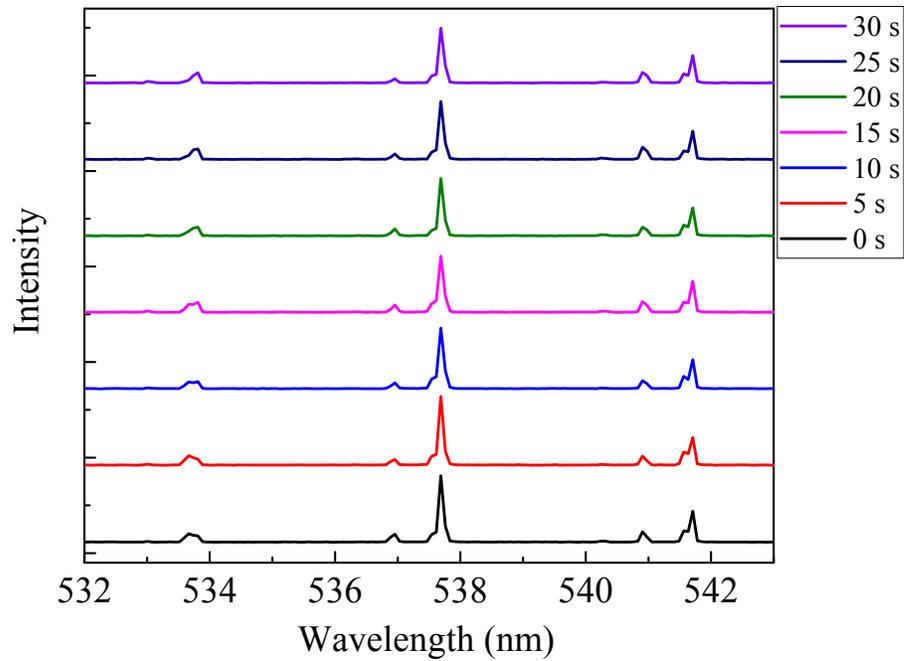
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Figure S1. Simulation of the scattered field of the incident Gaussian beam



(a) The simulation geometry.  $n_{\text{water}}=1.33$ .  $n_{\text{mirror}}=1.45$ . A Gaussian beam along y-axis is set as the background field with an FWHM of  $60 \mu\text{m}$  and the focal plane at the top interface between the mirror and water. (b) The distribution of the scattered field when there is no microwell structure (the SU-8 domain is set to be the same as the water domain,  $n_{\text{SU-8}}=1.33$ ). (c) The distribution of the scattered field when the microwell is present ( $n_{\text{SU-8}}=1.60$ ). The integration of the scattered field energy density over the water domain is  $2.45\text{e-}19$  and  $2.60\text{e-}19 \text{ J/m}^3$  without and with the SU-8 microwell structure, respectively. The discrepancy is less than 10%, thus we conclude that the microwell structure does not have any significant effect on the excitation profile.

Figure S2. Lasing stability of a cell captured in a microwell



Lasing stability of a cell captured in a microwell. The cell was continuously excited for 30 seconds with a pump intensity of  $100 \mu\text{J}/\text{mm}^2$  per pulse. The cell underwent up to 600 excitation pulses during the test, about twice the number of pulses that a cell might receive when it was scanned for 6 times in our work. As shown in Figure S2, no wavelength shift is observed in those lasing peaks.

### The thermal effect on the microwell-integrated cell lasing array

Given that there is environmental temperature fluctuation during the experiment, the resulting lasing wavelength shift is examined as follows. The dependence of the lasing wavelength  $\lambda$  on

temperature can be written as  $\frac{\Delta\lambda}{\lambda} = \left( \frac{1}{n_{eff}} \frac{dn_{eff}}{dT} + \frac{1}{L} \frac{dL}{dT} \right) \Delta T$ , where  $n_{eff}$  is the effective refractive

index,  $\frac{dn_{eff}}{dT}$  the thermal-optic coefficient,  $L$  the cavity length and  $\frac{dL}{dT}$  the linear thermal expansion coefficient. A cell laser cavity consists of mainly water, DNAs and proteins, with an effective thermal-optic coefficient estimated to be  $-1 \sim -4 \times 10^{-4}/^\circ\text{C}$ .<sup>1-3</sup> Since the SU-8 layer acts as the spacer for the laser cavity, the linear thermal expansion coefficient of SU-8 ( $52 \times 10^{-6}/^\circ\text{C}$ , MicroChem

Corp.) is used to estimate  $\frac{1}{L} \frac{dL}{dT}$ . Thus, for a temperature drift of  $1^\circ\text{C}$ ,  $\left| \frac{\Delta\lambda}{\lambda} \right| < 0.05\%$  and the corresponding lasing wavelength shift  $\Delta\lambda$  around 540 nm is about 0.2 nm.

We also examine the thermal effect of the pump laser. At  $120 \mu\text{J}/\text{mm}^2$  pump fluency, the total energy impinged on an area of  $2 \times 10^{-4} \text{mm}^2$  (the area of a cell) is  $2.4 \times 10^{-8} \text{J}$  per excitation pulse. Typically, less than 10% of the total energy can be absorbed (absorption cross section of dyes at excitation wavelength  $\sigma_a \sim 1 \times 10^{-16} \text{cm}^2$ , dye concentration  $C \sim 1 \text{mM}$ , gain medium length  $L \sim 10 \mu\text{m}$ , thus absorbance  $A = \sigma_a CL = 0.06$ ) and only a fraction that non-radiatively dissipates turns into heat (assume to be 40%, since quantum yield of SYTO9=0.6). The resulting fluctuation in

local temperature is estimated to be around  $0.3^\circ\text{C}$  ( $\Delta T = \frac{E_{heat}}{C_p m}$ ,  $E_{heat}$  the energy that turns into heat,  $C_p$  the specific heat (water:  $4.184 \text{J}/(\text{g}\cdot^\circ\text{C})$ ),  $m$  the mass). The lasing wavelength might have a drift of 0.06 nm accordingly. However, no build-up effect is expected under continuous pulse excitation due to the small duty cycle (5 ns/50 ms) of the 20 Hz OPO laser.

## References

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