## (Supplementary Materials)

## Ideal strength of two-dimensional stanene may reach or exceed Griffith strength estimate

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Table S1. Mechanical properties of stanene computed by adopting six different functionals.

|  |  | PBEsol (RRKJ) | PBEsol (PAW) | PW-LDA | HSE06 | PBE0 | B3LYP |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Atoms/supercell |  | 6 | 6 | 6 | 2 | 6 | 2 |
| $a_{0}(\AA)$ |  | 4.587 | 4.580 | 4.484 | 4.665 | 4.484 | 4.483 |
| $\boldsymbol{\delta}(\AA)$ |  | 0.837 | 0.848 | 0.769 | 0.860 | 0.748 | 0.727 |
| Bond length ( $\AA$ ) |  | 2.773 | 2.776 | 2.701 | 2.827 | 2.695 | 2.689 |
| Bond angle |  | $111.346^{\circ}$ | $111.128^{\circ}$ | $112.238^{\circ}$ | $111.178^{\circ}$ | $112.603^{\circ}$ | $112.967^{\circ}$ |
| UTS ( $\mathrm{N} / \mathrm{m}$ ) | Armchair | 3.656 | 3.718 | 4.076 | 2.919 | 3.821 | 3.622 |
|  | Zigzag | 4.084 | 4.097 | 4.520 | 3.380 | 4.224 | 3.930 |
|  | Biaxial | 3.318 | 3.361 | 3.670 | 3.318 | 3.333 | 3.186 |
| UTS (GPa) | Armchair | 11.216 | 11.405 | 12.502 | 8.953 | 11.720 | 11.110 |
|  | Zigzag | 12.527 | 12.567 | 13.866 | 10.369 | 12.957 | 12.054 |
|  | Biaxial | 10.178 | 10.311 | 11.258 | 10.178 | 10.223 | 9.773 |
| $\varepsilon_{\text {UTS }}^{\text {True }}$ | Armchair | 0.215 | 0.199 | 0.199 | 0.199 | 0.191 | 0.182 |
|  | Zigzag | 0.315 | 0.285 | 0.300 | 0.270 | 0.247 | 0.262 |
|  | Biaxial | 0.191 | 0.182 | 0.166 | 0.191 | 0.157 | 0.157 |
| $\varepsilon_{\text {UTS }}^{\text {Eng }}$ | Armchair <br> Zigzag <br> Biaxial | 0.240* | 0.220 | 0.220 | 0.220 | 0.210 | 0.200 |
|  |  | 0.370 | 0.330 | 0.350 | 0.310 | 0.280 | 0.300 |
|  |  | 0.210 | 0.200 | 0.181 | 0.210 | 0.170 | 0.170 |
| Poisson ratio $\boldsymbol{v}$ |  | 0.456 | 0.369 | 0.465 | 0.425 | 0.474 | 0.408 |
| E (N/m) | Armchair <br> Zigzag | $24.448{ }^{\ddagger}$ | 26.600 | 29.865 | 21.842 | 26.470 | 30.880 |
|  |  | 23.897 | 26.685 | 29.637 | 21.637 | 27.670 | 30.620 |
| E/UTS ratio | Armchair | 6.687 | 7.154 | $\underline{7.328}$ | 7.483 | $\underline{6.928}$ | $\underline{8.526}$ |
|  | Zigzag | 5.852 | 6.514 | 6.556 | 6.401 | 6.550 | 7.792 |

*Consider elongation of a specimen from an initial length $L_{0}$ to a final length $L$ and denote the stretch ratio $\left(\frac{L_{0}}{L}\right)$ as $\lambda$. The engineering, true, and Lagrangian strain can be expressed, respectively, as $\varepsilon^{T r u e}=\ln \lambda, \varepsilon^{\text {Eng }}=\lambda-1$, and $\eta^{L a g}=\frac{1}{2}\left(\lambda^{2}-1\right)$.
${ }^{\ddagger}$ Our Young's Modulus computed by suing the PEBsol functional implemented in Quantum ESPRESSO matches extremely well with another group's result ( $24.46117 \mathrm{~N} / \mathrm{m}$ ) with different DFT software ${ }^{1}$.

Table S2. Non-zero second and higher-order elastic constants (in Voigt notation and unit N/m) computed by adopting the six exchange-correlational (XC) functionals.

| XC functional (method) | Second-order | Third-order | Fourth-order | Fifth-order |
| :---: | :---: | :---: | :---: | :---: |
| PBEsol (RRKJ) | $\mathrm{C}_{11}=30.87$ | $\mathrm{C}_{111}=-191.79$ | $\mathrm{C}_{1111}=186.06$ | $\mathrm{C}_{111111}=5213.8$ |
|  | $\mathrm{C}_{12}=14.08$ | $\mathrm{C}_{112}=-136.76$ | $\mathrm{C}_{1112}=2053.8$ | $\mathrm{C}_{11112}=-18823$ |
|  | $\mathrm{C}_{22}=30.41$ | $\mathrm{C}_{222}=-195.37$ | $\mathrm{C}_{1122}=-1822.5$ | $\mathrm{C}_{11122}=18452$ |
|  |  |  | $\mathrm{C}_{2222}=100.94$ | $\mathrm{C}_{12222}=916.99$ |
|  |  |  |  | $\mathrm{C}_{22222}=8260.8$ |
| PBEsol (PAW) | $\mathrm{C}_{11}=29.80$ | $\mathrm{C}_{111}=-135.66$ | $\mathrm{C}_{1111}=-1099.6$ | $\mathrm{C}_{11111}=17937$ |
|  | $\mathrm{C}_{12}=9.765$ | $\mathrm{C}_{112}=-32.13$ | $\mathrm{C}_{1112}=63.32$ | $\mathrm{C}_{11112}=-1863.4$ |
|  | $\mathrm{C}_{22}=30.62$ | $\mathrm{C}_{222}=-182.80$ | $\mathrm{C}_{1122}=1794.1$ | $\mathrm{C}_{11122}=-2513.2$ |
|  |  |  | $\mathrm{C}_{2222}=-463.62$ | $\mathrm{C}_{12222}=474.7$ |
|  |  |  |  | $\mathrm{C}_{22222}=15743$ |
| PW-LDA (MT) | $\mathrm{C}_{11}=38.55$ | $\mathrm{C}_{111}=-347.24$ | $\mathrm{C}_{1111}=2686.2$ | $\mathrm{C}_{11111}=-15198$ |
|  | $\mathrm{C}_{12}=18.297$ | $\mathrm{C}_{112}=-232.50$ | $\mathrm{C}_{1112}=3035.3$ | $\mathrm{C}_{11112}=-20797$ |
|  | $\mathrm{C}_{22}=38.37$ | $\mathrm{C}_{222}=-377.62$ | $\mathrm{C}_{1122}=-3357.6$ | $\mathrm{C}_{11122}=12033$ |
|  |  |  | $\mathrm{C}_{2222}=3084.1$ | $\mathrm{C}_{12222}=-16596$ |
|  |  |  |  | $\mathrm{C}_{22222}=-14809$ |
| HSE06 (GHHT) | $\mathrm{C}_{11}=26.64$ | $\mathrm{C}_{111}=-186.64$ | $\mathrm{C}_{1111}=597.15$ | $\mathrm{C}_{11111}=-173.32$ |
|  | $\mathrm{C}_{12}=11.306$ | $\mathrm{C}_{112}=-78.37$ | $\mathrm{C}_{1112}=484.47$ | $\mathrm{C}_{11112}=-2206.2$ |
|  | $\mathrm{C}_{22}=26.47$ | $\mathrm{C}_{222}=-226.93$ | $\mathrm{C}_{1122}=172.39$ | $\mathrm{C}_{11122}=-3804.1$ |
|  |  |  | $\mathrm{C}_{2222}=1590$ | $\mathrm{C}_{12222}=-7560.6$ |
|  |  |  |  | $\mathrm{C}_{22222}=-7188.7$ |
| PBE0 (MT) | $\mathrm{C}_{11}=34.02$ | $\mathrm{C}_{111}=-278.59$ | $\mathrm{C}_{1111}=1175.8$ | $\mathrm{C}_{111111}=-653.09$ |
|  | $\mathrm{C}_{12}=16.027$ | $\mathrm{C}_{112}=-251.68$ | $\mathrm{C}_{1112}=3337.6$ | $\mathrm{C}_{11112}=-29272$ |
|  | $\mathrm{C}_{22}=34.87$ | $\mathrm{C}_{222}=-287.28$ | $\mathrm{C}_{1122}=-4882.2$ | $\mathrm{C}_{11122}=28258$ |
|  |  |  | $\mathrm{C}_{2222}=608.02$ | $\mathrm{C}_{12222}=-1721.6$ |
|  |  |  |  | $\mathrm{C}_{22222}=9730.1$ |
| B3LYP (MT) | $\mathrm{C}_{11}=37.05$ | $\mathrm{C}_{111}=-305.08$ | $\mathrm{C}_{1111}=1411.8$ | $\mathrm{C}_{111111}=-4319.5$ |
|  | $\mathrm{C}_{12}=15.119$ | $\mathrm{C}_{112}=-123.33$ | $\mathrm{C}_{1112}=812.42$ | $\mathrm{C}_{11112}=-1967.7$ |
|  | $\mathrm{C}_{22}=36.83$ | $\mathrm{C}_{222}=-331.66$ | $\mathrm{C}_{1122}=1283$ | $\mathrm{C}_{11122}=-8760.9$ |
|  |  |  | $\mathrm{C}_{2222}=1710.9$ | $\mathrm{C}_{12222}=-13248$ |
|  |  |  |  | $\mathrm{C}_{22222}=-1519.4$ |

Table S3. Indifferent results obtained through 2-atom supercell simulation (for phonon) and 6-atom supercell simulation (for stress-strain response) with the implementation of LDA.

|  | Phonon calculation (DFPT) | Stress-strain calculation |
| :---: | :---: | :---: |
| Functional | LDA | LDA |
| k-grid | $21 \times 21 \times 1$ | $11 \times 11 \times 3$ |
| Atoms/supercell | 2 | 6 |
| $\boldsymbol{a}_{\mathbf{0}}(\AA)$ | 4.483 | 4.484 |
| $\boldsymbol{\delta}(\AA)$ | 0.771 | 0.769 |
| Bond length $(\AA \AA)$ | 2.701 | 2.701 |
| Bond angle | $113.187^{\circ}$ | $112.238^{\circ}$ |
| Armchair | 4.107 | 4.076 |
| Zigzag | 4.508 | 4.520 |
| UTS (N/m) | 3.690 | 3.670 |

Figure S1. Demonstration of the negligible effect of spin-orbit coupling in mechanical property calculations in stanene ( 2 atoms/supercell, LDA functional incorporated).


Figure S2. The sensitivity of Young's modulus fitting on a true stress-true strain curve. The two fitted curves are almost identical but yield slightly different Efitted. The choice of DFT data point (A vs. B) as the choice of onset point of nonlinearity has an appreciable influence on determining elastic constants. To eliminate this ambiguity, E was determined from elastic constants fitted to PK2 stress Lagrangian strain curve.


Figure S3. (a) Phonon dispersion and density of states for stanene at the undeformed state. (b) The reciprocal space and first Brillouin zone with high symmetrical points associated with the phonon calculations. The present dispersion shows excellent agreement with the previous phonon calculations for stanene using either DFPT ${ }^{2,3}$ or the small displacement method ${ }^{4}$.


Figure S4. The computed stress-strain response of stanene by using the PBEsol (RRKJ), PBEsol (PAW), PW-LDA, HSE06, PBE0, and B3LYP functional, respectively.


Movie S1. (Online) Charge density in the horizontal middle plane passing through the buckled stanene at different strain values under uniaxial armchair, zigzag, and equibiaxial tension is shown.

## References

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