Supporting Information

# Organic Field-Effect Transistors Integrated with $\mathrm{Ti}_{2} \mathrm{CT}_{\mathrm{x}}$ Electrodes 

Shen Lai ${ }^{\text {a }}$, Sung Kyu Jang ${ }^{\text {a }}$, Jeong Ho Cho ${ }^{\text {a }}$, and Sungjoo Lee ${ }^{\text {a,b }}{ }^{\text {† }}$
${ }^{\text {a }}$ SKKU Advanced Institute of Nanotechnology (SAINT), Sungkyunkwan University (SKKU), Suwon 440-746, Korea
${ }^{\mathrm{b}}$ School of Information and Communication, Sungkyunkwan University (SKKU), Suwon 440746, Korea

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Figure S1. Optical microscopy (OM) image of organic field-effect transistor (OFET) device structure.

Proper $\mathrm{Ti}_{2} \mathrm{CT}_{\mathrm{x}}$ flakes are identified and metal electrodes are then formed. After pentacene is deposited, the source/drain contacts were formed between pentacene and $\mathrm{Ti}_{2} \mathrm{CT}_{\mathrm{x}}$ as shown in the above OM image.


Figure S2. Energy-dispersive spectrometry (EDS) analysis of a $\mathbf{T i}_{\mathbf{2}} \mathbf{C T}_{\mathbf{x}}$ flake. Scale bars, $1 \mu \mathrm{~m}$.

Figure S 2 a shows a scanning electron microscopy (SEM) image of a $\mathrm{Ti}_{2} \mathrm{CT}_{\mathrm{x}}$ flake on $\mathrm{SiO}_{2}$ substrate. Figure S2b and c show EDS mapping of the white dashed square in Figure S2a, for fluorine (F) and oxygen (O), respectively. Full coverage and uniform distribution of F and O are observed, which indicates that the $\mathrm{Ti}_{2} \mathrm{CT}_{\mathrm{x}}$ surface is fully terminated with surface groups $(-\mathrm{F},-\mathrm{OH}$ and $/ \mathrm{or}-\mathrm{O}) .{ }^{\mathrm{Sl}}$ This is consistent with the X-ray photoelectron spectroscopy (XPS) results in Figure 1 in main text.


Figure S3. Evaluation of the mobility attenuation factor $\boldsymbol{\theta}$. Note that $\boldsymbol{\theta}$ is extracted at the high gate field (strong accumulation region) where the $Y$ function method is valid.

The transfer characteristics ( $\mathrm{I}_{\mathrm{ds}}$ versus $\mathrm{V}_{\mathrm{g}}$ ) in the linear region can be expressed as:
$\mathrm{I}_{\mathrm{ds}}=\mathrm{G}_{\mathrm{m}} \mathrm{V}_{\mathrm{ds}}\left(\mathrm{V}_{\mathrm{g}}-\mathrm{V}_{\mathrm{th}}\right) /\left[1+\theta\left(\mathrm{V}_{\mathrm{g}}-\mathrm{V}_{\mathrm{th}}\right)\right] .^{\mathrm{S} 2}$
Here, $\mathrm{G}_{\mathrm{m}}=(\mathrm{W} / \mathrm{L}) \mu_{0} \mathrm{C}_{\mathrm{i}}$ is the transconductance parameter in which $\mu_{0}$ is low-field mobility; $\mathrm{V}_{\text {th }}$ is the threshold voltage; and $\theta$ is the mobility attenuation factor.

The Y -function is defined as $\mathrm{Y}=\mathrm{I}_{\mathrm{ds}} / \mathrm{g}_{\mathrm{m}}{ }^{0.5},{ }^{\mathrm{S} 3}$ so
$\left.\mathrm{Y}=\left(\mathrm{G}_{\mathrm{m}} \mathrm{V}_{\mathrm{ds}}\right)^{0.5}\left(\mathrm{~V}_{\mathrm{g}}-\mathrm{V}_{\mathrm{th}}\right)=\left[(\mathrm{W} / \mathrm{L}) \mu_{0} \mathrm{C}_{\mathrm{i}} \mathrm{V}_{\mathrm{ds}}\right)\right]^{0.5}\left(\mathrm{~V}_{\mathrm{g}}-\mathrm{V}_{\mathrm{th}}\right)$, then
$\mathrm{dY}=\left[(\mathrm{W} / \mathrm{L}) \mu_{0} \mathrm{C}_{\mathrm{i}} \mathrm{V}_{\mathrm{ds}}\right)^{0.5} \mathrm{dV}_{\mathrm{g}}$
Thus, $\mu_{0}$ can be extracted from the slope of Y versus $\mathrm{V}_{\mathrm{g}}$. (see Figure 3d in main text)
Once $\mu_{0}$ is known, $\mathrm{G}_{\mathrm{m}}$ can be extracted as $\mathrm{G}_{\mathrm{m}}=(\mathrm{W} / \mathrm{L}) \mu_{0} \mathrm{C}_{\mathrm{i}}$

The transconductance, $\mathrm{g}_{\mathrm{m}}$, can be obtained:
$\mathrm{g}_{\mathrm{m}}=\mathrm{dI}_{\mathrm{ds}} / \mathrm{dV}_{\mathrm{g}}=\mathrm{G}_{\mathrm{m}} \mathrm{V}_{\mathrm{ds}} /\left[1+\theta\left(\mathrm{V}_{\mathrm{g}}-\mathrm{V}_{\mathrm{th}}\right)^{2}\right.$, so
$\mathrm{g}_{\mathrm{m}}{ }^{0.5}=\left(\mathrm{G}_{\mathrm{m}} \mathrm{V}_{\mathrm{ds}}\right)^{0.5 /[ }\left[1+\theta\left(\mathrm{V}_{\mathrm{g}}-\mathrm{V}_{\mathrm{th}}\right)\right]$, and
$1 / \mathrm{g}_{\mathrm{m}}^{0.5}=\left[1+\theta\left(\mathrm{V}_{\mathrm{g}}-\mathrm{V}_{\mathrm{th}}\right)\right] /\left(\mathrm{G}_{\mathrm{m}} \mathrm{V}_{\mathrm{ds}}\right)^{0.5}$, then
$\mathrm{d}\left(1 / \mathrm{g}_{\mathrm{m}}{ }^{0.5}\right)=\left[\theta /\left(\mathrm{G}_{\mathrm{m}} \mathrm{V}_{\mathrm{ds}}{ }^{0.5}\right] \mathrm{dV}_{\mathrm{g}}\right.$
So $\theta$ can be extracted from the slope of $1 / \mathrm{gm}^{0.5}$ versus $\mathrm{V}_{\mathrm{g}}$. (see Figure S 3 )
After $\mathrm{G}_{\mathrm{m}}$ and $\theta$ are both known, Rc can be calculated from $\theta=\mathrm{G}_{\mathrm{m}} \times$ Rc. ${ }^{\text {S3 }}$


Figure S4. Part of the transfer curve in Figure 3a and $\mathbf{V}_{\mathrm{g}}$ dependence of mobility.
Figure S4 shows the transfer curve in Figure 3a on a linear scale, with the mobility values in Figure S4 all extracted by the linear-regime equation: $\mu=\left(\mathrm{dI}_{\mathrm{ds}} / \mathrm{dV}_{\mathrm{g}}\right) \times\left[\mathrm{L} /\left(\mathrm{WC}_{\mathrm{i}} \mathrm{V}_{\mathrm{ds}}\right)\right]$. The linear-regime equation is valid when: $\left|\mathrm{V}_{\mathrm{ds}}\right|<\left|\mathrm{V}_{\mathrm{g}}-\mathrm{V}_{\mathrm{th}}\right|{ }^{\mid}{ }^{4} \mathrm{~V}_{\mathrm{ds}}$ was -20 V and $\mathrm{V}_{\mathrm{th}}$ was extracted as -25 V , so the linear-regime
equation is valid when $\mathrm{V}_{\mathrm{g}}<-45 \mathrm{~V}$, shown as the green region in Figure S 4 . The extracted mobility values are valid in this region, and are located in a very narrow range $\left(0.9-1.1 \mathrm{~cm}^{2} \mathrm{~V}^{-1} \mathrm{~s}^{-1}\right)$ where the $\mathrm{V}_{\mathrm{g}}$ dependence of the mobility can barely be observed.


Figure S5. $I_{d s}$ vs. $V_{g}$ curve in Figure 3d in liner scale.

We used the constant-Rc mode to fit the experimental curve with the equation: $\mathrm{I}_{\mathrm{ds}}=\mathrm{V}_{\mathrm{ds}} /\left\{\mathrm{Rc}+1 /\left[\left(\mathrm{V}_{\mathrm{g}}-\right.\right.\right.$ $\left.\left.\left.\mathrm{V}_{\mathrm{th}}\right) \mu_{0} \mathrm{C}_{\mathrm{i}} \mathrm{W} / \mathrm{L}\right]\right\},{ }^{\mathrm{S} 5,6}$ where $\mu_{0}$ is the contact-free mobility. As shown in Figure S 5 , the shape of the experimental curve could be well-described by the constant-Rc mode. The extracted Rc was $2.3 \mathrm{k} \Omega \mathrm{cm}$, which is consistent with the Rc value extracted by the Y -function method ( $3 \mathrm{k} \Omega \mathrm{cm}$ ).

## References

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