

## Supplementary Information

### for Derivation from Equation (new 2) to Equation (new 6)

Equation (10, new 6) is the variant of Equation (6, new 2) and has been derived from H. L. Ong's work<sup>1,2</sup>.

In equation (10, new 6),

$$f_{\theta} = \frac{1}{2} \chi_a H^2 \sin^2 \theta_m - 2\chi_a H^2 \int_0^d \sin^2 \theta dz + W_2 \sin^2 \theta_0 \quad (10, \text{new } 6)$$

Here we had omitted the complex derivation and caused your confusion. In short, term  $\theta_m$  contains the elastic constants and we just give the compact expression.

We have already know Eq. (6, new 2) as total free energy function under external magnetic field per area on substrate<sup>3, 22</sup>:

$$f = \int_0^d [(K_{11} \sin^2 \theta + K_{33} \cos^2 \theta) \left(\frac{d\theta}{dz}\right)^2 - \chi_a H^2 \sin^2 \theta] dz + \frac{1}{2} [\delta(z) + \delta(dz)] \sum_{n=1}^{\infty} W_{2n} \sin^{2n} \theta \quad (6)$$

$K_{11}$  and  $K_{33}$  are splay and bend elastic constants,  $\chi_a$  is the magnetic anisotropy of NLC, we can define  $\theta = \theta_m$  when  $z = \frac{d}{2}$ .

By using variational method, based on  $\frac{\partial f}{\partial \theta} - \frac{d}{dz} \left( \frac{\partial f}{\partial \left(\frac{\partial \theta}{\partial z}\right)} \right) = 0$ , we can get the system

equilibrium state's Euler equation:

$$(K_{11} \sin^2 \theta + K_{33} \cos^2 \theta) \frac{d^2 \theta}{dz^2} - (K_{33} - K_{11}) \cos \theta \sin \theta \left(\frac{d\theta}{dz}\right)^2 - 2\chi_a H^2 \sin \theta \cos \theta = 0 \quad (7)$$

$$\text{And boundary condition: } (K_{11} \sin^2 \theta + K_{33} \cos^2 \theta) \frac{d\theta}{dz} \Big|_{z=0,d} = -2W_2 \sin \theta_0 \cos \theta_0 \quad (8)$$

By doing further Dimensionality reduction of Eq. (7)

$$\text{We can have } (K_{11} \sin^2 \theta + K_{33} \cos^2 \theta) \left(\frac{d\theta}{dz}\right)^2 = \chi_a H^2 (\cos^2 \theta_m - \cos^2 \theta) \quad (18)$$

$$\text{And } \left(\frac{d\theta}{dz}\right)^2 = \frac{\chi_a H^2 (\cos^2 \theta_m - \cos^2 \theta)}{(K_{11} \sin^2 \theta + K_{33} \cos^2 \theta)} \quad (19)$$

Because LC cell is symmetrical around  $z = d/2$ , So we only need to consider  $0 \leq z \leq \frac{d}{2}$ ,

$$\frac{d\theta}{dz} = -\sqrt{\frac{\chi_a H^2}{K_{11}}} \sqrt{\frac{\cos^2 \theta_m - \cos^2 \theta}{1 + k \cos^2 \theta}} \quad (20)$$

$$\text{In which } k = \frac{K_{33} - K_{11}}{K_{11}}$$

Based on Eq. (18), we have  $H \sqrt{\frac{\chi_a}{K_{11}}} \int_0^{d/2} dz = - \int_{\theta_0}^{\theta_m} \sqrt{\frac{1+k \cos^2 \theta}{\cos^2 \theta_m - \cos^2 \theta}} d\theta$  (21)

And we can get  $H = -\frac{2}{l} \sqrt{\frac{K_{11}}{\chi_a}} \int_{\theta_0}^{\theta_m} \sqrt{\frac{1+k \cos^2 \theta}{\cos^2 \theta_m - \cos^2 \theta}} d\theta$  (22)

Then, based on Eq. (6) and Eq. (22)

$$f_\theta = \frac{1}{2} \chi_a H^2 \sin^2 \theta_m - 2 \chi_a H^2 \int_0^d \sin^2 \theta dz + W_2 \sin^2 \theta_0 \quad (10, \text{new } 6)$$

In which  $\theta = \theta_m \sin(\pi z / d) + O(\theta_m^3)$  ,  $\theta_m \approx [(H / H_0)^2 - 1] / 2B$  ,  $B = (1 - k - 9u / 4) / 4$  ,

$u = 1 - \chi_\perp / \chi_\parallel$  ,  $H_0 = (\pi / d) (K_{33} / \chi_a)^{1/2}$ .  $\theta_0$  is the director orientation of the LC at the substrates deviates from the easy direction which provides resistance to hold back LC's tendency to align at magnetic field's direction.

So equation (10, new 6) is dependent on the elastic splay or bend constants. It is just author's simplification that causes careful reader's confusion. Based on derivation above, we can know the first two part of equation (10, new 6) precisely and it is very useful for our combination of theory and experiment results.

#### Reference

1. Optically induced Fredericksz transition and bistability in a nematic liquid crystal, Hiap Liew Qng, PHYSICAL REVIEW A, VOLUME 28, NUMBER 4 OCTOBER 1983,
2. "Optical-field-enhanced and static-field-induced first-order Fredericksz transitions in a planar parallel nematic liquid crystal", Phys. Rev. A 33, 3550 (1986)