

## Low resistivity $\text{HfN}_x$ grown by plasma-assisted ALD with external rf substrate biasing

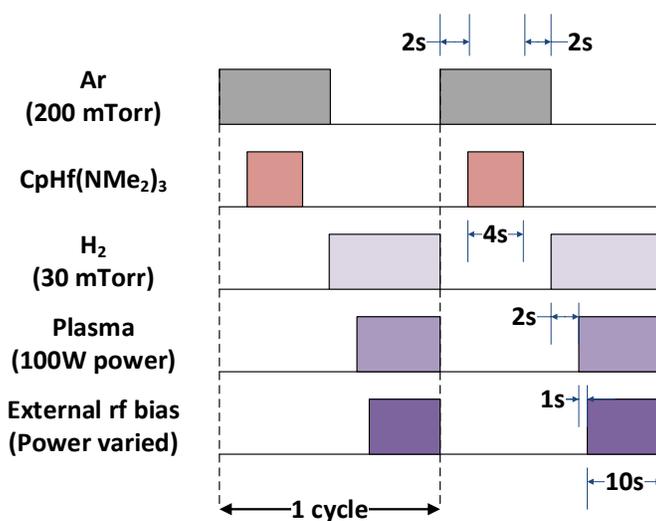
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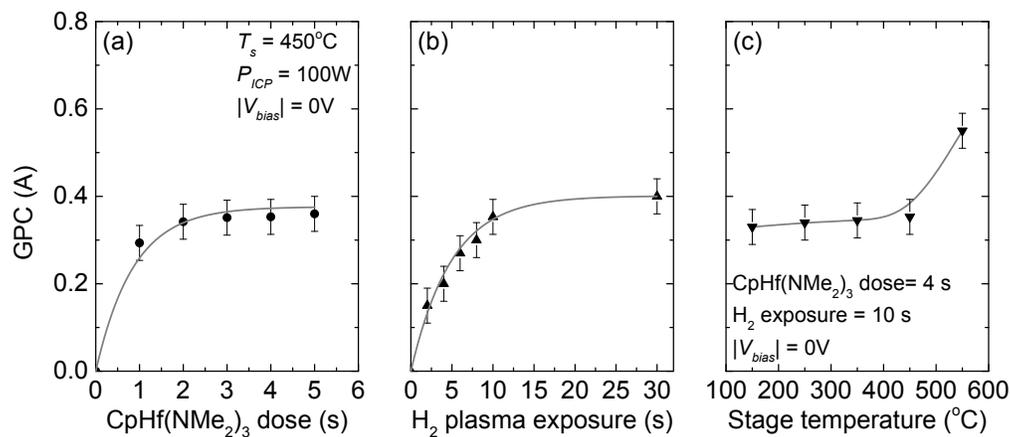
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### $\text{HfN}_x$ ALD process

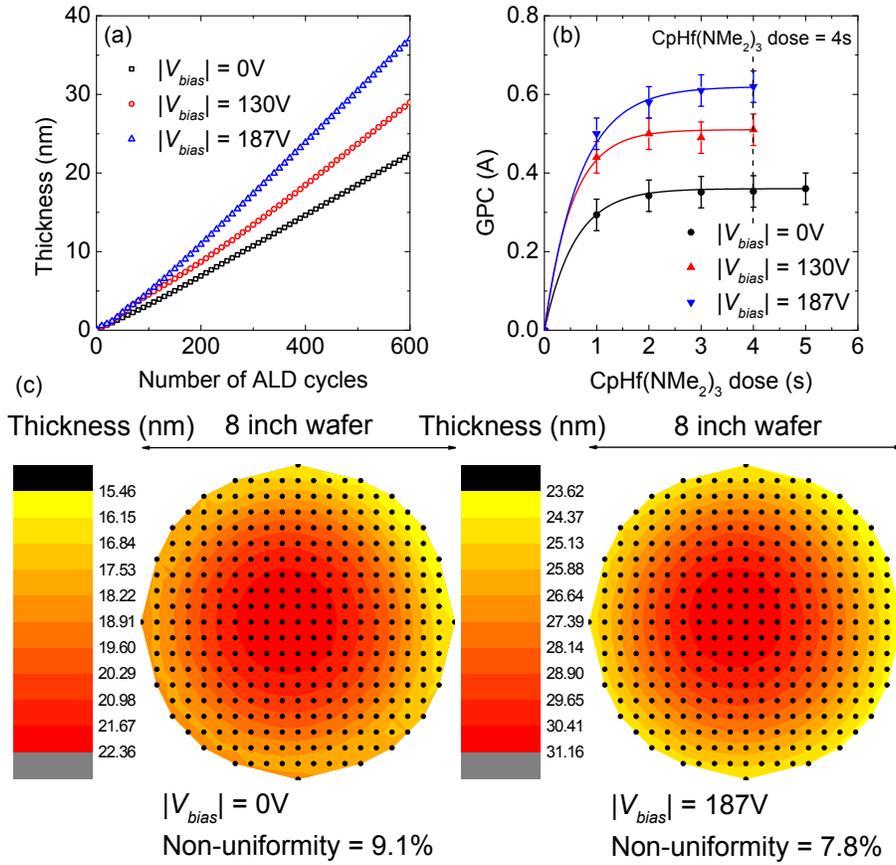


**Figure S1.** Time sequence for the complete ALD cycles of  $\text{HfN}_x$  using  $\text{CpHf(NMe}_2)_3$  and  $\text{H}_2$  plasma depicting the application of an external rf substrate bias during the  $\text{H}_2$  plasma step. A plasma stabilization time of 1s was used and can be neglected. Time intervals are not drawn to scale.



**Figure S2.** Growth per cycle (GPC) (Å) for HfN<sub>x</sub> films prepared at  $|V_{bias}| = 0V$  as a function of (a) CpHf(NMe<sub>2</sub>)<sub>3</sub> dose time and (b) H<sub>2</sub> plasma exposure time. The CpHf(NMe<sub>2</sub>)<sub>3</sub> dose, H<sub>2</sub> plasma exposure and purge times under saturation conditions were 4 s, 10 s and 2 s respectively. (c) GPC(Å) for the HfN<sub>x</sub> films prepared under saturation conditions as a function of the stage temperature showing an abrupt increase above 450°C. Lines serve as a guide to the eye.

Figure S3a shows a linear thickness increase with number of ALD cycles for HfN<sub>x</sub> films prepared at various values of  $|V_{bias}|$ . Furthermore, the ALD process was found to saturate using a CpHf(NMe<sub>2</sub>)<sub>3</sub> dose time of 4 s at all values of  $|V_{bias}|$  investigated (Figure S3b). Therefore, the application of external rf substrate bias does not affect the self-saturation behavior of the ALD process. The HfN<sub>x</sub> films were also prepared at  $|V_{bias}| = 0V$  and at  $|V_{bias}| = 187V$  on 8-inch Si wafers and the thickness non-uniformity was evaluated (Figure S3c). The thickness non-uniformity was deduced by taking the ratio between standard deviation and average film thickness (1-sigma). A similar thickness non-uniformity of 9.1% and 7.8% was obtained for the HfN<sub>x</sub> films prepared with grounded electrode and at  $|V_{bias}| = 187V$  respectively.



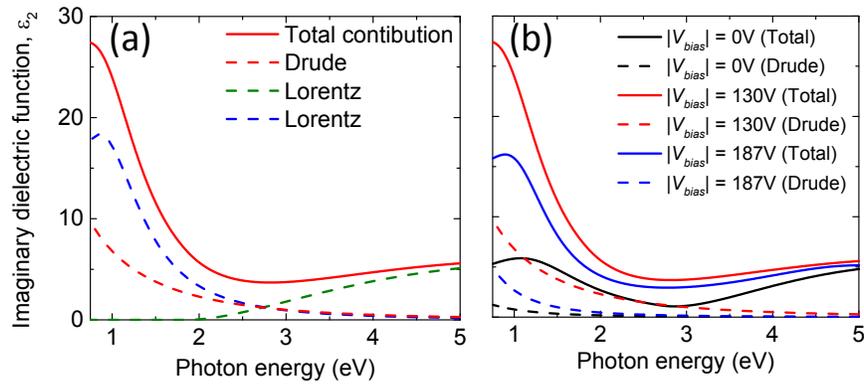
**Figure S3.** (a) HfN<sub>x</sub> film thickness as a function of number of ALD cycles prepared for various values of  $|V_{bias}|$ ; (b) GPC(Å) as a function of CpHf(NMe<sub>2</sub>)<sub>3</sub> dose time for HfN<sub>x</sub> films prepared at various values of  $|V_{bias}|$ ; (c) thickness uniformity maps on 8-inch Si wafer for HfN<sub>x</sub> films prepared with grounded electrode ( $|V_{bias}| = 0V$ ) and at  $|V_{bias}| = 187V$  showing a similar thickness non-uniformity in both cases. The black points on the wafer are the actual measurement points with 5 mm edge exclusion.

### Modeling of HfN<sub>x</sub> dielectric functions using spectroscopic ellipsometry

The dielectric functions can be modelled using one Drude and two Lorentz oscillators:

$$\varepsilon(E) = \varepsilon_1(E) + i\varepsilon_2(E) = \varepsilon_\infty - \frac{\overset{Drude}{E_p^2}}{E^2 - i\Gamma_D E} + \sum_{j=1}^2 \frac{\overset{Lorentz}{S_j E_{0j}^2}}{E_{0j}^2 - E_j^2 + iE_j \Gamma_j} \quad (1)$$

where,  $\varepsilon_\infty$  represents transitions at higher energy which are not accounted in Lorentz oscillators,  $E_p$  is plasma energy and  $\Gamma_D$  is the damping factor for Drude oscillator. The Lorentz oscillators are centered at  $E_o$  which corresponds to the resonance frequency, while  $S$  indicates the strength of the oscillators and  $\Gamma$  is the damping factor for the Lorentz oscillators.



**Figure S4.** Imaginary dielectric function  $\varepsilon_2$  for the  $\delta$ -HfN<sub>x</sub> films prepared at (a)  $|V_{bias}| = 130V$  obtained *via* fitting of the ellipsometry data with parameterization including one Drude and two Lorentz oscillators and (b) various values  $|V_{bias}|$ . The individual Drude oscillators for the films are also given revealing an increase in Drude absorption by applying an external rf substrate bias.

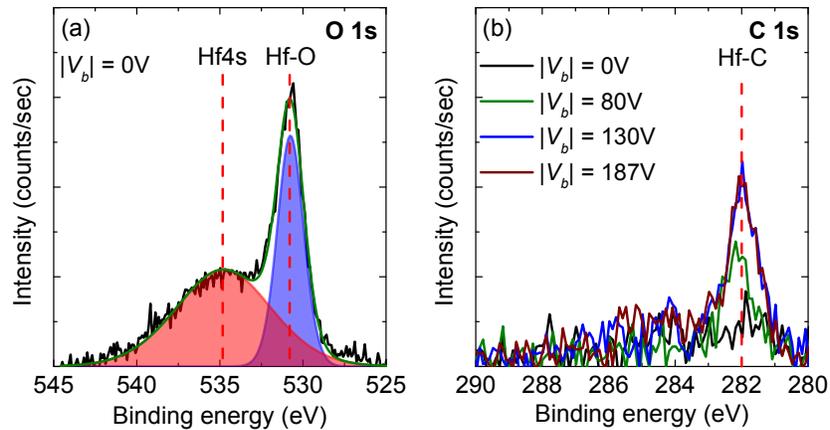
**Table S1.** SE fitting parameter values for one Drude and two Lorentz oscillators for the  $\delta$ -HfN film prepared at  $|V_{bias}| = 130V$ .

$ V_{bias} $ (V)	130	
$\varepsilon_{\infty}$	4.15	
$E_p$ (eV)	1.7	
Lorentz oscillator 1	$S$ (eV)	19.4
	$\Gamma$ (eV)	1.2
	$E_o$ (eV)	1.1
Lorentz oscillator 2	$S$ (eV)	5.4
	$\Gamma$ (eV)	4.3
	$E_o$ (eV)	5.7

Knoops *et al.* described that the insights into grain boundary scattering can be gained by probing opto-electronic properties using SE.<sup>1</sup> In Figure S4, the imaginary part of the dielectric function  $\varepsilon_2$  and the corresponding Drude absorption for the HfN<sub>x</sub> films grown at various values of  $|V_{bias}|$  are presented. An optical resistivity ( $\rho_{op}$ ) of  $(5.9 \pm 0.1) \cdot 10^{-3} \Omega\text{cm}$  was deduced from the magnitude of Drude absorption for the HfN<sub>x</sub> film grown with grounded electrode ( $|V_{bias}| = 0V$ ). An increase in  $|V_{bias}|$  to 130V led to an increase in the magnitude of Drude absorption, signifying an increase in the in-grain conduction electron density. An  $\rho_{op}$  of  $(9.0 \pm 0.2) \cdot 10^{-4} \Omega\text{cm}$  was deduced for the HfN<sub>x</sub> film grown at  $|V_{bias}| = 130V$ . A further increase in the  $|V_{bias}|$  to 187V led to a decrease in the magnitude of Drude absorption and an increase in  $\rho_{op}$  to  $(1.7 \pm 0.1) \cdot 10^{-3} \Omega\text{cm}$ . In addition, the

interaction distance ( $d_{int}$ ) of the incident light with the  $\text{HfN}_x$  films was calculated using the HfN effective mass of 0.88,<sup>2</sup> the fermi velocity ( $v_e = \hbar(3\pi^2N)^{1/3}/m^*$ ) and a photon energy of 0.75 eV (lower limit of SE) in the same manner as described by Knoops et al.<sup>1</sup> An  $d_{int}$  of 2.3 nm, 3.3 nm and 2.9 nm were deduced for  $\text{HfN}_x$  films grown at  $|V_{bias}|$  of 0V, 130V and 187V respectively. In view of the small interaction distance as compared to the lateral grain size of  $\sim 20\text{-}25$  nm of the  $\text{HfN}_x$  films (as will be explained in microstructural characterization section), it is plausible to use the difference between electrical and optical resistivity in order to judge the extent of scattering at the grain boundaries as a function of  $|V_{bias}|$ .

### Chemical composition analyses of $\text{HfN}_x$ films using X-ray photoelectron spectroscopy



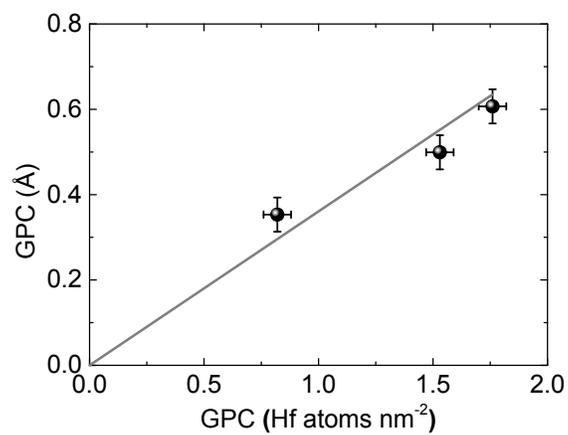
**Figure S5.** (a) De-convoluted O 1s XPS spectrum for  $\text{HfN}_x$  film prepared at  $|V_{bias}| = 0\text{V}$  showing the presence of Hf-O bonds and (b) C 1s XPS spectra for  $\text{HfN}_x$  films prepared at various values of  $|V_{bias}|$  showing an increase in the peak intensity of Hf-C bonds with an increase in  $|V_{bias}|$ .

**Table S2.** Corresponding peak assignment, binding energies and full width half maximum for the Hf, N, O and C spectral lines used to deconvolute the peaks, measured by XPS.

Spectral line	Peak designation	Binding energy (eV)	FWHM (eV)	Reference
Hf 4f <sub>7/2</sub>	Hf <sup>4+</sup>	16.0	1.6	3,4
Hf 4f <sub>7/2</sub>	Hf <sup>3+</sup>	14.9	1.6	5
N 1s	Hf <sup>3+</sup> N	397.5	1.5	6
N 1s	Hf <sup>4+</sup> N	396.5	1.6	7

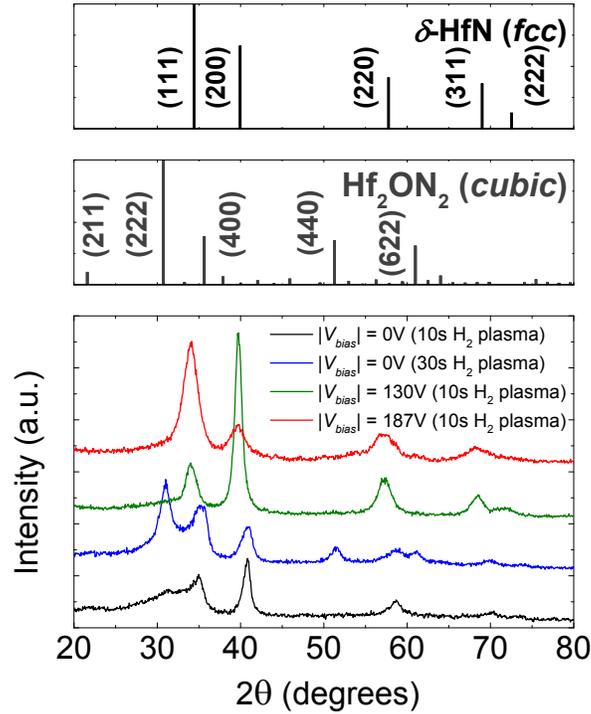
O 1s	HfO <sub>x</sub>	530.8	1.9	8
C 1s	HfC <sub>x</sub>	282	1.2	9

### Variation of GPC (Å) with GPC (Hf atoms nm<sup>-2</sup>)



**Figure S6.** GPC in terms of thickness as a function of GPC in terms of Hf atoms deposited per nm<sup>2</sup> illustrating a proportional dependence.

## GIXRD patterns for HfN<sub>x</sub> films as a function of $|V_{bias}|$

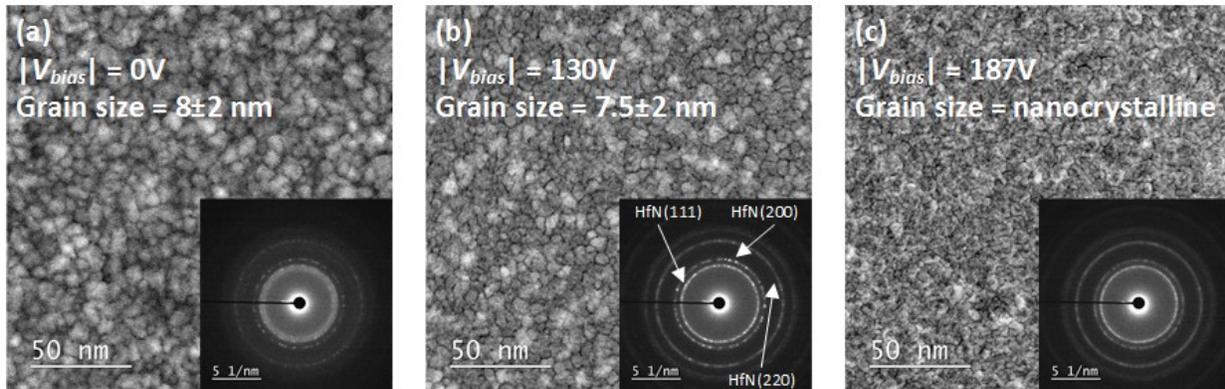


**Figure S7.** Grazing incidence X-ray diffractograms for  $\sim 75$  nm thick HfN<sub>x</sub> films prepared at various values of  $|V_{bias}|$  referenced with powder *fcc*  $\delta$ -HfN and *cubic* Hf<sub>2</sub>ON<sub>2</sub> XRD patterns.

## HAADF-STEM images for HfN<sub>x</sub> films during nucleation phase

The development of microstructure was studied during the nucleation phase of film growth by preparing  $\sim 15$  nm of HfN<sub>x</sub> films on Si<sub>3</sub>N<sub>4</sub> TEM windows that are coated with  $\sim 5$  nm ALD SiO<sub>2</sub>. Figure S8 shows the HAADF-STEM images of the HfN<sub>x</sub> layers deposited at various values of  $|V_{bias}|$ . The HfN<sub>x</sub> film grown at  $|V_{bias}| = 130V$  exhibits a lateral grain size of  $7.5 \pm 2.0$  nm whereas the film grown at  $|V_{bias}| = 187V$  was found to be nanocrystalline in nature. In addition, the selected area electron diffraction patterns (SAED) were acquired from  $1.3 \mu\text{m}$  diameter areas for the corresponding HfN<sub>x</sub> films (insets in Figure S8). Interestingly, no significant difference in crystallographic texture was observed in the nucleation layer for both the HfN<sub>x</sub> films. However, the HfN<sub>x</sub> film grown at  $|V_{bias}| = 130V$  exhibits discontinuous diffraction rings signifying the formation of relatively *large* crystallites whereas a more continuous diffraction pattern was

obtained for the HfN<sub>x</sub> film grown at  $|V_{bias}| = 187V$ . This result implies that at  $|V_{bias}| = 187V$ , relatively *small* lateral crystallites are formed.



**Figure S8.** Top-view high-angle annular dark-field scanning transmission electron microscope (HAADF-STEM) images for ~15 nm thick HfN<sub>x</sub> films prepared (a)  $|V_{bias}| = 0V$  (b)  $|V_{bias}| = 130V$  and (c)  $|V_{bias}| = 187V$ . Inset depicts the corresponding selected area electron diffraction (SAED) patterns.

## References

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