Supplementary Information

| Calculations details | .2 |
|--|----|
| Magnetic Interactions | .2 |
| Magnetic Topology | .3 |
| Effect of spin-orbit coupling | .4 |
| Ag ₂ ZnZr ₂ F ₁₄ reference system | .5 |
| Coupling constant calculations for LP, HP1, and HP2 | .7 |
| Structures in VASP format | 0 |

Calculations details

In the present study we have used structures of AgF_2 polymorphs derived from DFT+U calculations performed in ref. 6. The supercells of the LP, HP1, and HP2 polymorphs used in this study are given, for selected pressures, at the end of this document.

Magnetic Interactions

Our propose here is to illustrate with a simple analytical model the origin of the strong angular dependence found in the main text, and the sign and order of magnitude of the interaction in the different configurations. We consider an Ag-F-Ag bridge. The exchange interaction can be computed using perturbation theory in the *pd* hopping matrix elements. We take the same setting as in Ref. 4 with one $d^{x^2-y^2}$ in each Ag and two p orbitals in the bridging F: p_{\parallel} , oriented parallel to the Ag-Ag direction and p_z which is oriented perpendicular to the Ag-Ag direction but parallel to the triangle formed by the Ag-F-Ag complex. Since most parameters are not known we take only a minimal set of parameters different from zero to illustrate the main trends. The dependence on the angle α of the bridge is determined by the hopping matrix elements $t_{\parallel d}$ between p_{\parallel} and $d_{x^2-y^2}$ and t_{zd} between p_z and $d_{x^2-y^2}$ which can be parameterized in terms of Slater-Koster²⁵ integrals as $t_{\parallel d} = t_{pd} \sin(\alpha/2)$ and $t_{zd} = t_{pd} \cos(\alpha/2)$. Here t_{pd} is the hopping matrix element for a straight bond. In addition there is a dependence of the hopping matrix element on the Ag-F distance which can be parameterized as $t_{pd} \sim 1/R^4$. Due to the small changes in bond length with pressure, the correction due to this effect is much less relevant and will be neglected in the present simplified treatment. Therefore, all the computations are done keeping the t_{pd} the same for all bond lengths. The magnetic interaction is computed as the energy difference between the singlet and the triplet in the Ag-F-Ag system and reads, $I = I^{(2)} + I^{(4,SE)} + I^{(4,HR)},$

where $\int^{(2)} \propto t \|_{d}^{2}$ is a ferromagnetic contribution due to direct exchange, $K_{\|d} > 0$ between the $d_{x^2 - y^2}$ and the $p_{\|}$ orbital,

$$J^{(2)} = t_{pd}^{2} \sin^{2}\left(\frac{\alpha}{2}\right) \left[\frac{1}{(\Delta - K_{\parallel d})} - \frac{1}{(\Delta - K_{\parallel d})}\right],$$

 $J^{(4,SE)} \propto \left(t_{\parallel d}^2 - t_{zd}^2\right)^2$ is the superexchange antiferromagnetic contribution,

$$J^{(4,SE)} = -t_{pd}^{4} \cos^{2} \alpha \frac{1}{\Delta^{2}} \left[\frac{4}{U_{d}} + \frac{8}{(2\Delta + U_{p})} \right],$$

and $J^{(4,HR)} \propto t \frac{2}{\|d|^2} dt_{zd}^2$ is a ferromagnetic contribution due to Hund's rule exchange interaction $J_H > 0$ on fluorine,

$$J^{(4,HR)} = t_{pd}^{4} \sin^{2} \alpha \frac{8}{\Delta^{2}} \left[\frac{1}{(2\Delta + U_{p} - J_{H})} - \frac{1}{(2\Delta + U_{p})} \right]$$

Similar expressions were presented in the supplementary information to Ref. 4 except that $J^{(4,HR)}$ contribution was neglected. In Fig. 3a of the main text we compare these perturbative expressions with the results obtained with the full DFT solution. As a reference parameter set we took, $K \parallel d = 0.07 \text{ eV}$, $J_H = 0.7 \text{ eV}$ for Hund's interaction among p orbitals, $\Delta = 2.7 \text{ eV}$ for the charge transfer energy, $U_d = 9.4 \text{ eV}$ ($U_p = 4 \text{ eV}$) for the Hubbard interaction on Ag (F) and $t_{pd} = 1.15 \text{ eV}$.

We see from Fig. 3a of the main paper that the superexchange contribution, $J^{(4,SE)}$, is in general dominant except close to $\alpha = 90^{\circ}$ where it vanishes. Around this angle the sign of the magnetic interaction gets reversed due to the direct exchange contributions and in agreement with Goodenough-Kanamori-Anderson rules. $J^{(4,HR)}$ appears at fourth order so it is a much smaller contribution than $J^{(2)}$ and vanishes in the case of a straight bond where only p is relevant. For $J^{(2)}$, we only took into account the direct exchange interaction with the p orbital. A more accurate treatment would consider the direct exchange with both orbitals and the angular dependence of the K's matrix elements.

The parameters used are similar but not equal to the ones of Ref. 4 which were optimized only for zero pressure. There is considerable freedom for the choice of parameters. For example, practically an equally good fit of the DFT data is obtained with $U_d = 6 \text{ eV}$ and $t_{pd} = 1.1 \text{ eV}$. A more realistic computation would take into account also four-ring exchange processes that are expected to be important in covalent materials. Also, an accurate parameterization would require to take into account the effects of other orbitals which, however, would jeopardize our intention to get a simple understanding of the trends. With these caveats, we see that the general trends of the DFT computation are reproduced and the main microscopic matrix elements giving rise to the interactions are identified.

Magnetic Topology

In order to discuss the possible magnetic ground state of the system it is useful to identify the motifs formed by the more relevant superexchange paths identified.

For the LP phase, layer directions are equivalent from a magnetic point of view so in a first approximation the system can be described by the two-dimensional Heisenberg antiferromagnet within layers (Supplementary Fig. 1a) with weak coupling among layers. The ground state has robust 2D antiferromagnetic order inside the layers which becomes long-range three-dimensional order below the Néel temperature. Neglecting spin-orbit coupling the magnetic excitation spectrum is gapless due to acoustic spin-wave modes.



Supplementary Figure 1: Magnetic lattices corresponding to the AgF₂ polymorphs shown in Fig. 1 of the main text. The structure can be seen as formed by sheets in LP phase(a) zig-zag chains in HP1 (b), and ladders in HP2 (c).

For the HP1 phase, two non-equivalent paths appear within the layers forming a zig-zag pattern (Supplementary Fig. 1b). Lowering the temperature the system is expected to show 1D quasi-long range order along the J'_{2D} zig-zag chains at high temperatures which becomes more 2D as the temperature is lower and finally 3D long-range antiferromagnetic order below the Néel temperature. Close to the critical pressure to the HP2 phase the J''_{2D} coupling vanishes and the system becomes effectively 1D at all temperatures. The ground state is gapless unless a dimerization and a further lowering of symmetry (not considered here) occurs.

For the HP2 phase, the J''_t chains can be seen as the legs of a ladder. We conventionally take J'_t as being the rungs of the ladder by a fictitious displacement of one leg respect to the other. With this setting J''_t forms "diagonal" rungs (Fig. 3c). The ground state consists of OD singlets defined on the rungs of the ladder²² and remains gapped up to zero temperature unless weaker couplings among ladders drive 3D antiferromagnetic order at low temperature. Given the large difference in values of the interactions we think the OD-singlet ground state scenario is more likely.

Effect of spin-orbit coupling

The effect of spin orbit coupling was analyzed using the GGA+U method (U=5eV) as implemented in VASP. We checked that the ordering of the phases is not altered respect to SCAN meta-GGA without spin-orbit coupling.

Supplementary Figure 1 shows the magnetic structure of the LP polymorph. The figure setting reproduces Fig. 3 of Ref. 2. In excellent agreement with this reference we find that magnetic moments are approximately directed in the crystallographic **a** direction.



Supplementary Figure 2: Magnetic structure of LP polymorph to be compared with the experimental result in Ref 2. Dark Ag atoms are at z=1/2 while light gray are at z=0. There is a trivial inversion of the magnetic moments in the *ab* plane respect to Fig. 3 of Ref. 2 related to the choice of origin of the unit cell.

Ag₂ZnZr₂F₁₄ reference system

 $Ag_2ZnZr_2F_{14}$ was chosen as a reference system due to the presence of planar $[AgF_4]^{2-}$ units connected via a single F atom bridge forming $Ag_2F_7^{3-}$ dimers with straight Ag–F–Ag bridges in its crystal structure (Supplementary Figure 3).²³

We have evaluated the superexchange coupling constants within the dimers (*J*, Fig. 3a), as well as along two inter-dimer routes (J_{i1} , Fig 3b, J_{i2} , Fig 3c). The Ag···Ag distances along the *J*, J_{i1} , and J_{i2} routes are 4.03 Å, 3.97 Å, and 5.60 Å, respectively. Apart from these contacts there are no other Ag···Ag distances shorter than 5.8 Å.



Supplementary Figure 3. The Ag/F bonding framework of $Ag_2ZnZr_2F_{14}$ (a) featuring the main intradimer superexchange pathway, as measured by J, together with the depiction of the inter-dimer J_{11} (b) and J_{12} (c) superexchange routes.



Supplementary Figure 4. Inter-dimer coupling constants in Ag₂ZnZr₂F₁₄.

The J_{i1} and J routes form a honeycomb-like lattice (Fig. 4a), while J_{i2} connects the dimers into a square lattice (Fig. 4b). The SE routes not taken into account in this analysis are J_{i3} (Fig 4c, Ag...Ag distances of 5.82 Å) which connect the dimers into chains, and J_{i4} (Fig 4d, Ag...Ag distances of 6.64 Å) through which ladders are formed. As shown below already the J_{i1} and J_{i2} constants are small (absolute values smaller than 0.5 meV), and one should expect J_{i3} and J_{i4} to be negligibly small.

Supplementary Table 1. Magnetic states of Ag₂ZnZr₂F₁₄ and their energy. Spin up/down sites are indicated with a +/– sign, site labelling follows that of Supplementary Figure 3; E_{nm} denotes the part of the total-energy of the system which is independent of the spin state.

| Site: | 1 | 2 | 3 | 4 | Energy per f.u. | Energy (eV per f.u.) |
|-------|---|----------|----------|-------------|---|--------------------------------|
| F1 | + | + | + | + | $-0.125J - 0.25J_{i1} - 0.5J_{i2} + E_{nm}$ | -120.8945 |
| F2 | + | + | _ | _ | $-0.125J + 0.25J_{i1} + 0.5J_{i2} + E_{nm}$ | -120.8945 |
| A1 | + | _ | + | _ | $0.125J + 0.25J_{i1} - 0.5J_{i2} + E_{nm}$ | -120.9725 |
| A2 | + | - | _ | + | $0.125J - 0.25J_{i1} + 0.5J_{i2} + E_{nm}$ | -120.9729 |
| | | | Su | per- | exchange constants | Super-exchange constants (meV) |
| | | J | = 2 | EE_{A1} | $+ 2E_{A2} - 2E_{F1} - 2E_{F2}$ | J =- 312.7 |
| | | | J_{i1} | $= E_{\mu}$ | $E_{A1} - E_{A2} - E_{F1} + E_{F2}$ | $J_{i1} = 0.4$ |
| | | J_{i2} | = 0. | .5(– | $E_{A1} + E_{A2} - E_{F1} + E_{F2} \big)$ | $J_{i2} = -0.2$ |

The coupling constants were extracted with the broken symmetry method through equations given in Supplementary Table 1. The same method was for applied in an earlier study of the SE coupling in $Ag_2ZnZr_2F_{14}$.²⁰ They used spin states with the following energies:

$$E_{AF1} = (+4J_1 - 8J_2 + 8J_3)(1/4)$$
$$E_{AF2} = (-4J_1 + 8J_2)(1/4)$$
$$E_{AF3} = (+4J_1 + 8J_2)(1/4)$$
$$E_{AF3} = (+4J_1 - 8J_2 - 8J_3)(1/4)$$

where J_1 is the intra-dimer coupling constant and J_2 and J_3 are inter-dimer couplings. From the above one can see that:

 $J_1 = 1/2(E_{AF3} - E_{AF2})$

From Table 4 of ref. 20 one can easily see that the AF2 states lies 1174.81 to 638.6 meV higher than AF3 for the *U* parameter in the DFT+U method equal from 2 to 6 eV. This gives intra-dimer coupling constants in the range from -587.4 to -319.3 meV. The latter value, which corresponds to a more realistic *U* value in the case of silver(II) compounds, is close to that computed in the present work.

Coupling constant calculations for LP, HP1, and HP2

For the LP polymorph (Fig. 5a) we have evaluated the intra-sheet coupling constant (J_{2D}) and two intersheet ones (J_{11}, J_{12}) as shown in Fig. 5b.



Supplementary Figure 5. The Ag/F bonding framework of LP-AgF₂ (a), together with the depiction of the intra-sheet coupling constant (J_{2D}), as well as the two inter-sheet SE routes (b). The 1x1x2 supercell used in the calculations is depicted in (c).

Supplementary Table 2. Magnetic states of LP-AgF₂ and their energy. Spin up/down sites are indicated with a +/– sign, site labelling follows that of Supplementary Figure 5c; E_{nm} denotes the part of the total-energy of the system which is independent of the spin state.

| Site: | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | Energy per f.u. | | |
|---|--------------------------|---|---|---|---|---|---|---|---|--|--|
| F4 | + | + | + | + | _ | _ | + | + | E _{nm} | | |
| A1 | + | + | + | + | _ | _ | _ | _ | $0.5J_{2D} - 0.5J_{i1} + 0.5J_{i2} + E_{nm}$ | | |
| A2 | + | + | _ | _ | _ | _ | + | + | $-0.5J_{2D} + 0.5J_{i1} + 0.5J_{i2} + E_{nm}$ | | |
| A4 | + | + | _ | _ | + | + | _ | _ | $0.5J_{2D} + 0.5J_{i1} - 0.5J_{i2} + E_{nm}$ | | |
| | Super-exchange constants | | | | | | | | | | |
| $J_{2D} = E_{A1} + E_{A4} - 2E_{F4} \qquad \qquad J_{i1} = E_{A2} + E_{A4} - 2E_{F4} J_{i2} = E_{A1} + E_{A2} - 2E_{F4}$ | | | | | | | | | | | |

In case of the HP1 polymorph the basic cell is analogous to that of LP (Fig. 5a/5c), but due to the alternation in the Ag-F-Ag angle there are two intra-sheet couplings $(J'_{2D} \text{ and } J''_{2D})$. The inter-sheet SE also split into pairs $(J'_{i1}/J''_{i1} \text{ and } J'_{i2}/J''_{i2})$, but given their small value we have evaluated only their mean, that is $J^{mean}_{i1} = \frac{1}{2}(J'_{i1} + J''_{i1})$ and $J^{mean}_{i2} = \frac{1}{2}(J'_{i2} + J''_{i2})$. The corresponding equations are given in Table 3.

Supplementary Table 3. Magnetic states of HP1-AgF₂ and their energy. Spin up/down sites are indicated with a +/- sign, site labelling follows that of Supplementary Figure 5c, except for the F7 state which is depicted in Supplementary Figure 6; E_{nm} denotes the part of the total-energy of the system which is independent of the spin state.

| Site: | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | Energy per f.u. |
|------------------------------|---|---|---|-------|------|---|---|---|--|
| F3 | + | + | _ | + | _ | + | + | + | $-0.25J'_{2D} - 0.25J''_{2D} + E_{nm}$ |
| F4 | + | + | + | + | _ | _ | + | + | E_{nm} |
| F7 | | | | see l | Fig. | 6 | | | $0.25J'_{2D} - 0.25J''_{2D} + E_{nm}$ |
| A1 | + | + | + | + | _ | _ | _ | _ | $0.25J'_{2D} + 0.25J''_{2D} - 0.5J''_{i1} + 0.5J''_{i2} + E_{nm}$ |
| A2 | + | + | _ | _ | _ | _ | + | + | $-0.25J'_{2D} - 0.25J''_{2D} + 0.5J^{mean}_{i1} + 0.5J^{mean}_{i2} + E_{nm}$ |
| A3 | + | _ | + | + | _ | _ | _ | + | $0.25J'_{2D} + 0.25J''_{2D} + E_{nm}$ |
| A4 | + | + | _ | - | + | + | _ | _ | $0.25J'_{2D} + 0.25J''_{2D} + 0.5J^{mean}_{i1} - 0.5J^{mean}_{i2} + E_{nm}$ |
| Super-exchange constants | | | | | | | | | |
| J" _{2D} = | $J'_{2D} = 2E_{A3} - 2E_{F7}$ $J''_{2D} = 2E_{F7} - 2E_{F3}$ $J''_{i1} = E_{A2} + E_{A4} - 2E_{F4}$ $J''_{i2} = E_{A1} + E_{A2} - 2E_{F4}$ | | | | | | | | |
| $J_{2D} = 2E_{F7} - 2E_{F3}$ | | | | | | | | | |

Supplementary Figure 6. The 2x1x2 supercell of HP1 used in the calculations of the F7 magnetic state (see Table 3). Spin up/down sites are marked with blue/red balls.

For the HP2 structure apart from the intra-tube SE routes (Supplementary Figure 7a) we considered three inter-tube couplings (J_{i1-i3} , Supplementary Figure7b-d).



Supplementary Figure 7. The 1x1x2 supercell of HP2 used in the calculations depicting the intra-tube SE routes (a), as well as the inter-tube ones (b-d).

Supplementary Table 4. Magnetic states of HP2-AgF₂ and their energy. Spin up/down sites are indicated with a +/– sign, site labelling follows that of Supplementary Figure 7a.

| Site: | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 | 11 | 12 | 13 | 14 | 15 | 16 |
|-------|--------------------------|---|-----|---|---|----|---|---|---|------|----|------|-----|------------|------|----|
| A1 | + | + | _ | _ | + | + | - | _ | + | + | _ | _ | + | + | - | _ |
| A2 | + | + | + | + | + | + | + | + | _ | _ | _ | _ | _ | _ | _ | _ |
| A3 | + | _ | _ | + | _ | + | + | _ | + | _ | _ | + | _ | + | + | _ |
| F1 | + | + | + | + | + | + | + | + | + | + | + | + | + | + | + | + |
| F2 | + | + | + | + | _ | _ | _ | _ | _ | _ | _ | _ | _ | _ | _ | _ |
| F3 | + | + | _ | _ | _ | _ | + | + | + | + | _ | _ | _ | _ | + | + |
| F4 | + | + | _ | _ | _ | _ | + | + | + | + | _ | _ | _ | _ | + | + |
| | Super-exchange constants | | | | | | | | | | | | | | | |
| | 1 | _ | 1 (| Г | 1 | 2E | | | | 0.05 | | 1 [] | 1 4 | <u>с</u> т | 0.05 | |

 $\begin{aligned} J'_{t} &= -1.6E_{F4} - 1.2E_{F3} - 1.6E_{F2} - 0.8E_{F1} + 4E_{A3} + 4E_{A2} + 0.8E_{A1} \\ J''_{t} &= -1.0666E_{F4} + 0.5333E_{F3} + 1.6E_{F2} - 0.5333E_{F1} + 0.9333E_{A2} - 1.4666E_{A1} \\ J'''_{t} &= 1.6E_{F4} - 2.8E_{F3} + 1.6E_{F2} + 0.8E_{F1} - 4E_{A3} + 3.6E_{A2} - 0.8E_{A1} \\ J_{i1} &= 2.1333E_{F4} + 0.9333E_{F3} - 3.2E_{F2} - 0.9333E_{F1} + 0.1333E_{A2} + 0.9333E_{A1} \\ J_{i2} &= 1.6E_{F4} + 0.8E_{F3} + 1.6E_{F2} - 1.2E_{F1} - 0.4E_{A2} - 0.8E_{A1} \\ J_{i3} &= 0.1333E_{F4} - 0.0666E_{F3} + 0.8E_{F2} + 0.0666E_{F1} + 0.1333E_{A2} - 0.0666E_{A1} \end{aligned}$

Structures in VASP format

Below we give the structural information (in VASP format) for the supercells of LP (at 4 GPa), HP1 (at 12 GPa), and HP2 (at 30 GPa) which were used in the calculations. The lattice constants and the fractional coordinates are taken unchanged from our previous work (Ref.6 in the main paper) – in that work they have been taken from DFT+U calculations which were then used to propose various polymorphic forms to be tested in Rietveld fits to experimental data.

| LP @ 4 GPa | | |
|----------------|---------------|---------------|
| 1.0 | | |
| 4.9556121826 | 0.0000000000 | 0.000000000 |
| 0.0000000000 | 5.4461851120 | 0.0000000000 |
| | 0.00000000000 | 11.3111810684 |
| Ag F | | |
| 8 10 Direct | | |
| Direct | 0.00000000 | 0.00000000 |
| 0.00000000 | 0.000000000 | 0.000000000 |
| 0.00000000 | 0.000000000 | 0.500000000 |
| 0.50000000 | 0.000000000 | 0.250000000 |
| 0.50000000 | 0.000000000 | 0.750000000 |
| 0.00000000 | 0.500000000 | 0.250000000 |
| 0.00000000 | 0.500000000 | 0.750000000 |
| 0.50000000 | 0.500000000 | 0.000000000 |
| 0.50000000 | 0.500000000 | 0.500000000 |
| 0.169777229 | 0.315485209 | 0.433581322 |
| 0.169777229 | 0.315485209 | 0.933581352 |
| 0.830222785 | 0.684514761 | 0.066418663 |
| 0.830222785 | 0.684514761 | 0.566418648 |
| 0.669777215 | 0.184514791 | 0.066418663 |
| 0.669777215 | 0.184514791 | 0.566418648 |
| 0.330222785 | 0.815485239 | 0.433581322 |
| 0.330222785 | 0.815485239 | 0.933581352 |
| 0.330222785 | 0.684514761 | 0.183581337 |
| 0.330222785 | 0.684514761 | 0.683581352 |
| 0.669777215 | 0.315485209 | 0.316418678 |
| 0.669777215 | 0.315485209 | 0.816418648 |
| 0.830222785 | 0.815485239 | 0.316418678 |
| 0.830222785 | 0.815485239 | 0.816418648 |
| 0.169///229 | 0.184514791 | 0.183581337 |
| 0.169777229 | 0.184514791 | 0.683581352 |
| | | |
| HP1 @ 12 GPa | | |
| 1.0 | | |
| 4.4479866028 | 0.0000000000 | 0.0000000000 |
| 0 000000000 | 5 50613/15100 | 0 000000000 |

| 4.4479866028 | 0.000000000000 | 0.00000000000 |
|--------------|----------------|---------------|
| 0.0000000000 | 5.5061345100 | 0.0000000000 |
| 0.0000000000 | 0.0000000000 | 11.1615753174 |
| Ag F | | |
| 8 16 | | |
| Direct | | |
| 0.999999046 | 0.999995649 | 0.000001085 |
| 0.000000000 | 0.999995649 | 0.500001073 |
| 0.463440925 | 0.086044364 | 0.250001073 |
| 0.463440925 | 0.086044364 | 0.750001073 |
| 0.999999046 | 0.586044371 | 0.250001073 |
| 0.000000000 | 0.586044371 | 0.750001073 |

| 0.463440925 | 0.499995649 | 0.000001085 |
|-------------|-------------|-------------|
| 0.463440925 | 0.499995649 | 0.500001073 |
| 0.344793469 | 0.767811835 | 0.166536614 |
| 0.344793469 | 0.767811835 | 0.666536629 |
| 0.608406901 | 0.422753751 | 0.299424320 |
| 0.608406901 | 0.422753751 | 0.799424291 |
| 0.855033040 | 0.922753751 | 0.299424320 |
| 0.855033040 | 0.922753751 | 0.799424291 |
| 0.118646532 | 0.267811835 | 0.166536614 |
| 0.118646532 | 0.267811835 | 0.666536629 |
| 0.118646532 | 0.318228185 | 0.416536599 |
| 0.118646532 | 0.318228185 | 0.916536570 |
| 0.855033040 | 0.663286269 | 0.049424332 |
| 0.855033040 | 0.663286269 | 0.549424350 |
| 0.608406901 | 0.163286254 | 0.049424332 |
| 0.608406901 | 0.163286254 | 0.549424350 |
| 0.344793469 | 0.818228126 | 0.416536599 |
| 0.344793469 | 0.818228126 | 0.916536570 |

HP2 @ 30 GPa

1.0

| 5.0722217560 | 0.0000000000 | 0.0000000000 |
|--------------|--------------|---------------|
| 0.0000000000 | 7.8814549446 | 0.0000000000 |
| 0.0000000000 | 0.0000000000 | 11.4406347275 |
| Ag F | | |
| 16 32 | | |
| Direct | | |
| 0.203389764 | 0.872389853 | 0.469119519 |
| 0.203389764 | 0.872389853 | 0.969119549 |
| 0.796610236 | 0.127610177 | 0.030880490 |
| 0.796610236 | 0.127610177 | 0.530880511 |
| 0.296610236 | 0.627610147 | 0.219119504 |
| 0.296610236 | 0.627610147 | 0.719119489 |
| 0.703389764 | 0.372389853 | 0.280880481 |
| 0.703389764 | 0.372389853 | 0.780880451 |
| 0.296610236 | 0.372389853 | 0.469119519 |
| 0.296610236 | 0.372389853 | 0.969119549 |
| 0.703389764 | 0.627610147 | 0.030880490 |
| 0.703389764 | 0.627610147 | 0.530880511 |
| 0.203389764 | 0.127610177 | 0.219119504 |
| 0.203389764 | 0.127610177 | 0.719119489 |
| 0.796610236 | 0.872389853 | 0.280880481 |
| 0.796610236 | 0.872389853 | 0.780880451 |
| 0.500000000 | 0.358051956 | 0.125000000 |
| 0.500000000 | 0.358051956 | 0.625000000 |
| 0.500000000 | 0.805217743 | 0.125000000 |
| 0.500000000 | 0.805217743 | 0.625000000 |
| 0.179253072 | 0.413714439 | 0.300836891 |
| 0.179253072 | 0.413714439 | 0.800836921 |
| 0.320746928 | 0.913714468 | 0.300836891 |
| 0.320746928 | 0.913714468 | 0.800836921 |
| 0.179253072 | 0.586285532 | 0.050836906 |
| 0.179253072 | 0.586285532 | 0.550836921 |
| 0.320746928 | 0.086285546 | 0.050836906 |
| 0.320746928 | 0.086285546 | 0.550836921 |
| 0.500000000 | 0.641948044 | 0.375000000 |

| 0.500000000 | 0.641948044 | 0.875000000 |
|-------------|-------------|-------------|
| 0.500000000 | 0.194782287 | 0.375000000 |
| 0.500000000 | 0.194782287 | 0.875000000 |
| 0.000000000 | 0.694782257 | 0.375000000 |
| 0.000000000 | 0.694782257 | 0.875000000 |
| 0.000000000 | 0.141948074 | 0.375000000 |
| 0.000000000 | 0.141948074 | 0.875000000 |
| 0.820746899 | 0.586285532 | 0.199163094 |
| 0.820746899 | 0.586285532 | 0.699163079 |
| 0.679253101 | 0.086285546 | 0.199163094 |
| 0.679253101 | 0.086285546 | 0.699163079 |
| 0.820746899 | 0.413714439 | 0.449163109 |
| 0.820746899 | 0.413714439 | 0.949163079 |
| 0.679253101 | 0.913714468 | 0.449163109 |
| 0.679253101 | 0.913714468 | 0.949163079 |
| 0.000000000 | 0.305217743 | 0.125000000 |
| 0.000000000 | 0.305217743 | 0.625000000 |
| 0.000000000 | 0.858051956 | 0.125000000 |
| 0.000000000 | 0.858051956 | 0.625000000 |