

The Growth and Fluorescence of Phthalocyanine Monolayers, Thin Films and Multilayers on Hexagonal Boron Nitride

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ELECTRONIC SUPPLEMENTARY INFORMATION

1. Computational details

The geometric parameters of the species under study were optimized using the range-separated hybrid ω B97X-D functional with empirical dispersion correction^[1] in combination with the correlation-consistent cc-pVDZ basis set^[2] employing the (75,302) integration grid. The maximum overlap method (MOM)^[3] was used for the structural optimization of compounds in the S_1 excited state. The *h*BN surface was modelled by a flake consisting of 65 boron atoms and 65 nitrogen atoms, the flake's edges were capped with hydrogen atoms. The B–N bond lengths were calculated to be of 1.44 Å in average. Atomic positions of the surface were optimized and kept fixed in the subsequent calculations. The adsorption energy of H₂Pc molecule was calculated using the ω B97X-D/cc-pVTZ method with the structure optimized with the smaller cc-pVDZ basis set. The most energetically preferred adsorption site was used for calculating excited state properties. Excitation energies corresponding to absorption ($S_1 \leftarrow S_0$ transition) and fluorescence ($S_0 \leftarrow S_1$ transition) were determined with time-dependent density functional theory (TD-DFT) at the same level of theory using the (MOM-) ω B97X-D/cc-pVDZ optimized structures of the S_0 and S_1 states, respectively.

0–0 energies were estimated using zero-point vibrational energies computed at the (MOM-) ω B97X-D/cc-pVDZ level of theory within harmonic approximation. To save computational time, vibrational frequencies of adsorbed molecules were calculated without taking into account *h*BN surface. Atomic partial charges were calculated within ChelpG formalism.^[4] All calculations were performed with the Q-Chem software package.^[5]

Results and discussion

The equilibrium structures of the H₂Pc molecule adsorbed on the *h*BN surface are shown in Figure S1. The molecule-surface separation was calculated to be *ca.* 3.3 Å. The adsorption energy computed at the ω B97X-D/cc-pVTZ level is 3.35 eV (Table S1). To calculate fluorescence energies, structural optimization of the first singlet excited state, which is the first dipole-allowed excited state for all systems under study, was performed using MOM. It is worth noting that MOM provides structural parameters which are comparable to those calculated by TD-DFT but at lower computational cost.^[6]

Figure S2 shows that the difference between the S_0 and S_1 relaxed structures is rather small ($\Delta r = 0.00$ – 0.02 Å). Upon interaction with the surface bond lengths of H₂Pc are only slightly changed and adsorbed molecule maintains its planarity. The most intense band in the fluorescence spectra of isolated and adsorbed H₂Pc arises from the electronic transition between the highest occupied (HOMO) and the lowest unoccupied (LUMO) molecular orbitals which have π and π^* character, respectively (Figure S3). The corresponding transition density is presented in Figure S4 (see Krueger et al.^[7] for a discussion of transition density).

Calculated absorption ($S_1 \leftarrow S_0$) and fluorescence ($S_0 \leftarrow S_1$) energies with the corresponding transition dipole moments and oscillator strengths for compounds under study are collated in Table S2. Computed vertical transition energies agree well with the available experimental data ($\Delta\Delta E = 0.01$ – 0.08 eV; Table S3). 0–0 energies that include relaxation of the final electronic state and zero-point energy correction were also calculated (Table S3).

When going from the gas phase to adsorbed on the *h*BN surface the calculated position of maximum in the fluorescence spectrum is red-shifted by 0.11 eV. This corresponds to the non-resonant shift discussed in the main text. In order to explain the origin of this contribution to the observed red shift in fluorescence spectra after adsorption of H₂Pc molecule, an energy decomposition of the shift into the contributions from molecular distortion, electrostatic and orbital interactions was performed according to methodology suggested by Hochheim and Bredow.^[8] The orbital interactions between the molecule and the surface atoms were found to be responsible for the red shift ($\Delta E_{\text{chem}} = 0.11$ eV; Table S4).

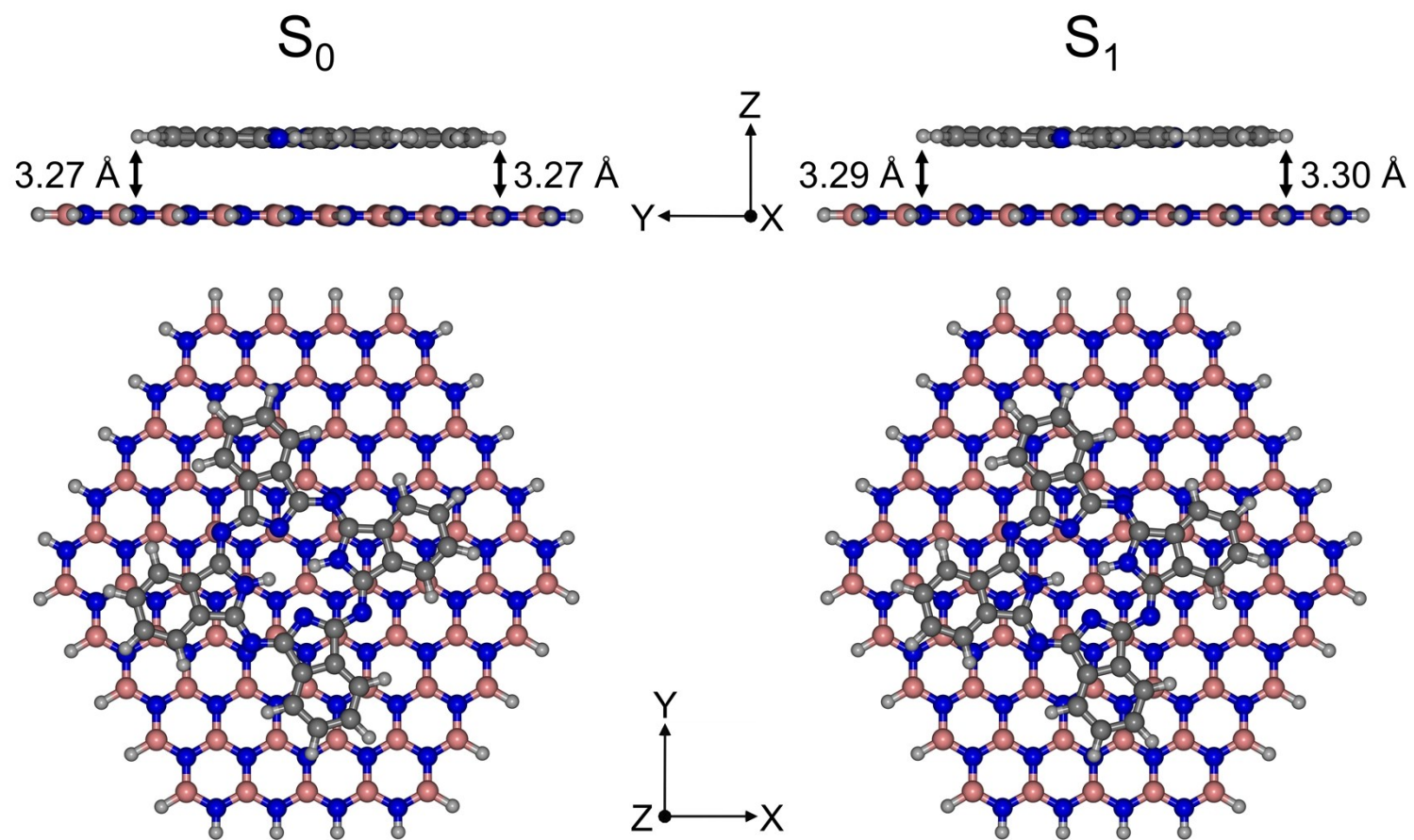
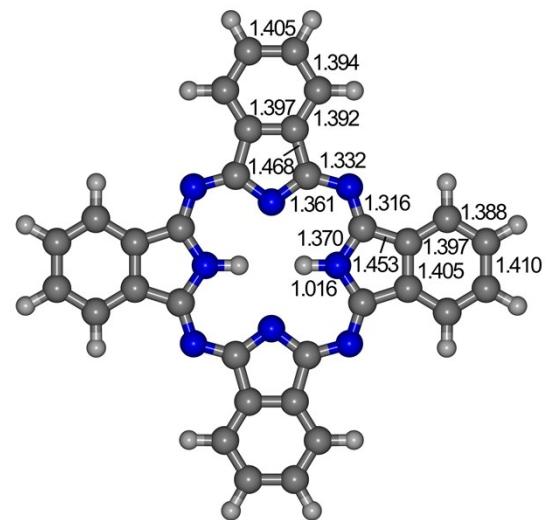
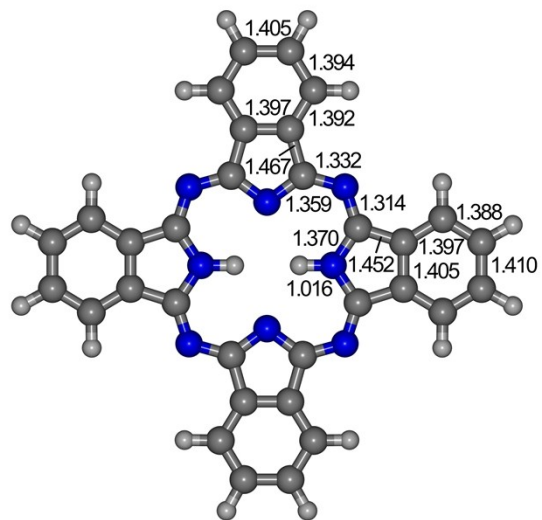


Figure S1. Structure of the H_2Pc/hBN system in S_0 and S_1 electronic states calculated at the (MOM-) ω B97X-D/cc-pVDZ level of theory.

H₂Pc

H₂Pc/hBN

S₀



S₁

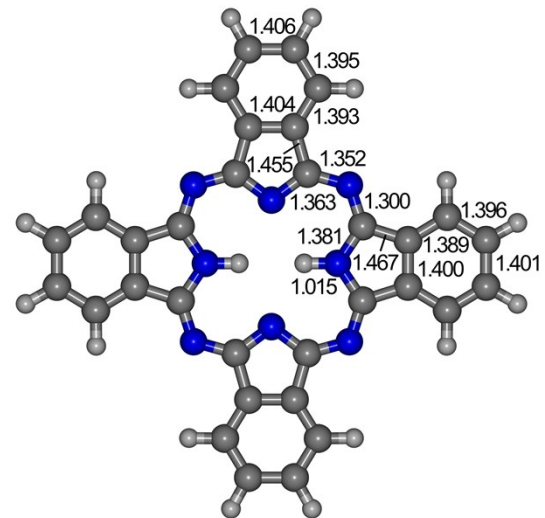
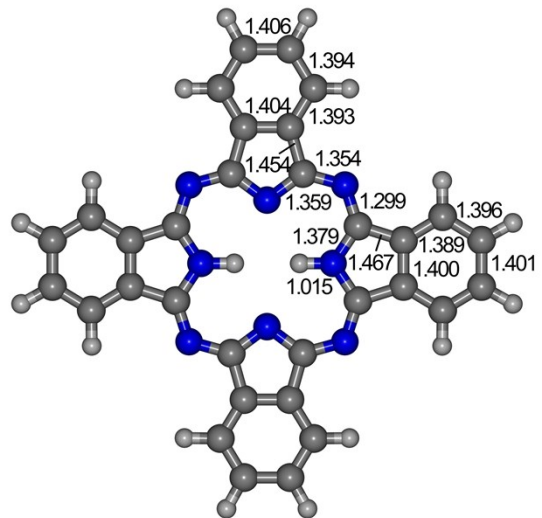


Figure S2. Selected bond lengths (Å) of H₂Pc in the gas phase and adsorbed on *h*BN in S₀ and S₁ electronic states calculated at the (MOM-)ωB97X-D/cc-pVDZ level of theory.

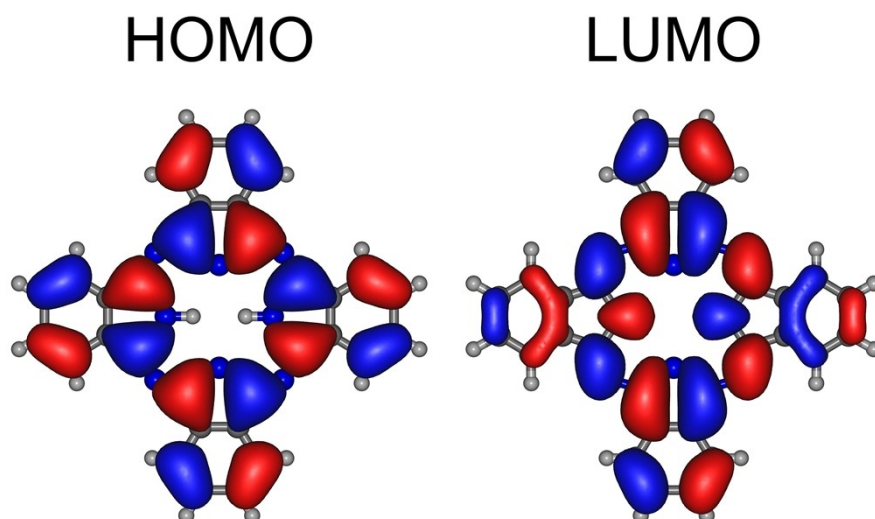


Figure S3. Highest occupied (HOMO) and lowest unoccupied (LUMO) molecular orbitals of H₂Pc in its equilibrium S₁ geometry on the *h*BN surface calculated at the ω B97X-D/cc-pVDZ level of theory. Isosurfaces with ± 0.01 a.u. isovalues are shown.

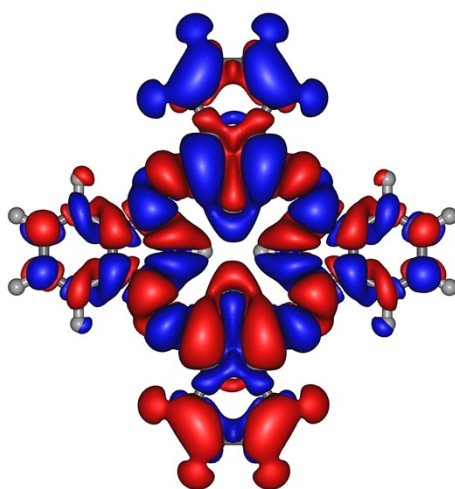


Figure S4. Transition density (S₀←S₁ fluorescence process) calculated for H₂Pc molecule in its equilibrium S₁ geometry on the *h*BN surface at the TD- ω B97X-D/cc-pVDZ level of theory. An isosurface with ± 0.0003 a.u. isovalues is shown.

Table S1. Energy of adsorption of H₂Pc molecule on the hBN surface (ΔE_{ads}) calculated at the ω B97X-D/cc-pVTZ// ω B97X-D/cc-pVDZ and ω B97X-D/cc-pVDZ// ω B97X-D/cc-pVDZ (in parentheses) theoretical levels.

System	ΔE_{ads} , eV
H ₂ Pc/hBN	3.35 (3.68)

Table S2. Details of absorption ($S_1 \leftarrow S_0$) and fluorescence ($S_0 \leftarrow S_1$) processes for H₂Pc, H₂Pc/hBN (in parentheses) and H₂Pc at H₂Pc/hBN geometry (in brackets) calculated at the TD- ω B97X-D/cc-pVDZ//((MOM) ω B97X-D/cc-pVDZ level of theory. Energies of frontier molecular orbitals are also given.

	Absorption	Fluorescence
Transition energy, eV / nm	1.96 / 632 (1.88) / (660) [1.96] / [632]	1.87 / 663 (1.76) / (703) [1.87] / [662]
Transition dipole moment, D	(-7.89, 0.00, 0.00) ((-6.34, -3.07, 0.06)) [(-7.62, -2.02, 0.04)]	(0.00, -8.93, 0.00) ((-1.18, 7.92, 0.02)) [(-1.60, 8.77, -0.01)]
Oscillator strength, au	0.46 (0.35) [0.46]	0.57 (0.43) [0.56]
$E_{\text{HOMO}} / E_{\text{LUMO}}$, eV	-6.31 {-6.41} ^a / -1.66 (-6.31 / -1.71) [-6.31 / -1.66]	-6.26 / -1.88 (-6.26 / -1.93) [-6.26 / -1.88]
$\Delta E_{\text{HOMO-LUMO}}$, eV	4.65{4.03} ^b (4.60) [4.65]	4.38 (4.33) [4.38]

^a -IP (ionization potential) experimental value. [9]

^b Value calculated by GW-BSE method. [9]

Table S3. Absorption ($S_1 \leftarrow S_0$), fluorescence ($S_0 \leftarrow S_1$), and 0–0 energies of H₂Pc and H₂Pc/hBN calculated at the TD- ω B97X-D/cc-pVDZ//((MOM) ω B97X-D/cc-pVDZ level of theory.

System	Absorption, eV / nm	Fluorescence, eV / nm	ΔE^{0-0} , eV / nm
H ₂ Pc	1.96 / 632 ^a [1.88 / 661] ^d [1.83 / 677] ^e	1.87 / 663 ^b [1.83 / 677] ^e	1.88 / 661 ^c
H ₂ Pc/hBN	1.88 / 660 ^a	1.76 / 703 ^b [1.75 / 708] ^f	1.75 / 706 ^c

^a Vertical absorption energy.

^b Vertical fluorescence energy.

^c Zero-point vibrational energy corrected 0–0 transition energy.

^d Experimental value (free jet). ^[10]

^e Experimental value (Ar matrix). ^[10]

^f Experimental value. This work.

Table S4. Energy decomposition of fluorescence shift for H₂Pc/hBN calculated at the TD- ω B97X-D/cc-pVDZ//((MOM) ω B97X-D/cc-pVDZ level of theory. All values are given in eV.

System	E_{gas}	E_{dist}	E_{pc}	E_{ads}	ΔE_{dist}	ΔE_{elec}	ΔE_{chem}	ΔE_{tot}
H ₂ Pc/hBN	1.87	1.87	1.87	1.76	0.00	0.00	0.11	0.11

2. Experimental Methods

Substrates are prepared by mechanically exfoliating hBN flakes from mm-scale crystals using the scotch tape method; the flakes are transferred to thermally oxidised silicon wafers, with an oxide thickness of 300 nm. The flakes are then immersed in toluene for approximately 12 hours, annealed in H₂:Ar (5% : 95%) at 400 °C for 8 hours, subjected to brief flame annealing prior to the introduction to the vacuum system.

Fluorescence spectroscopy is carried out using a Horiba LabRAM HR spectrometer, equipped with a 660 nm excitation laser. Laser powers in the range of 1-50 μ W are used to reduce photo-bleaching and damage to the sample. The sample morphology is determined using AFM, carried out under ambient conditions in tapping mode using the Asylum Research Cypher S with Multit75Al-G silicon cantilevers from Budget Sensors.

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