

SUPPORTING INFORMATION

Interactions between H-bonded $[\text{Cu}^{\text{II}}_3(\mu_3\text{-OH})]$ triangles; A combined magnetic susceptibility and EPR study.

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Details for the simulation of the solide powder EPR spectra

In order to account for broadening effects we assume one nuclear spin per trimer with $I_a = I_b = 3/2$ with a hyperfine interaction $A_{||,i} = 280$ MHz. It was found that better agreement between theory and experiment can be achieved by allowing for a slight rotation of the \mathbf{g}_a and \mathbf{g}_b along the \mathbf{r}_{ab} vector to an angle $\xi = 0.5^\circ$ and introducing an ordering of the molecules according to the orientation distribution [see reference 26] $P(\theta) = \exp(-U(\theta))$ with $U(\theta) = -\lambda(3 \cos 2\theta - 1)/2$, where θ is the angle between the molecular z axis and the static magnetic field. The ordering parameter was fixed to $\lambda = -10$ (<http://www.easyspin.org/>). The simulation of the parallel mode X-band spectrum does not account for several resonances in the 250 - 500 mT range which are remnants from the allowed $\Delta M_S = \pm 1$ transitions. The little difference of the obtained parameters at the two frequencies are attributed to instrumental conditions

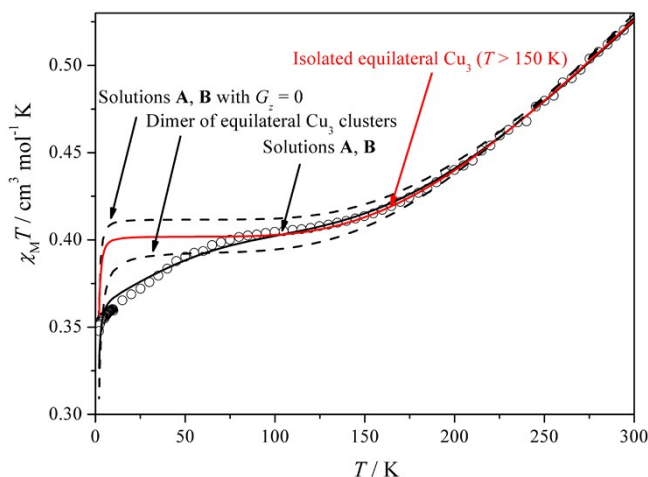


Figure S1. $\chi_M T$ vs T experimental data and calculated curves for **1** according to the models described in the text. The solid black line represents fit according to solution **B**. In the dimer of equilateral Cu_3 cluster model $\chi_M T$ is considered per trimeric unit.

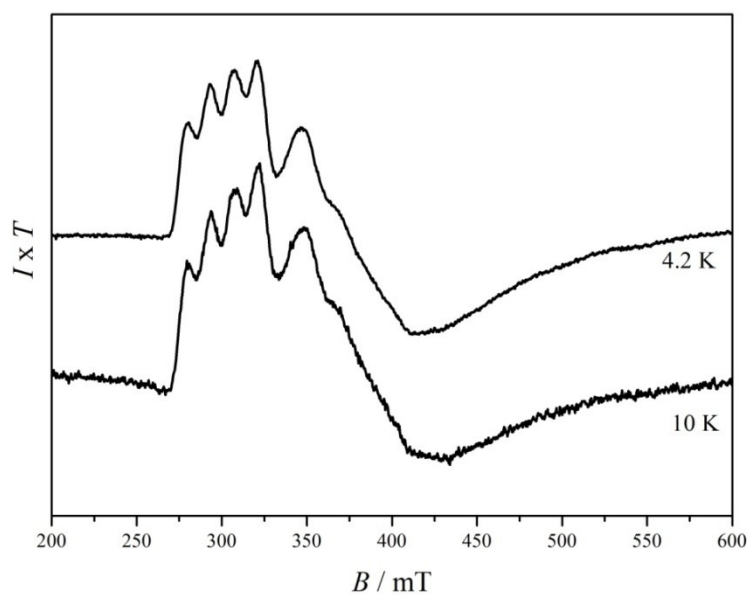


Figure S2. X-band EPR spectra of a frozen solution in THF of **1** at 4.2 and 10 K. The signals, scaled as $I \times T$, are super-imposable indicating that they arise from a ground state. EPR conditions: microwave frequency, 9.41 GHz; microwave power, 0.7 μW ; modulation amplitude, 10 G_{pp} .

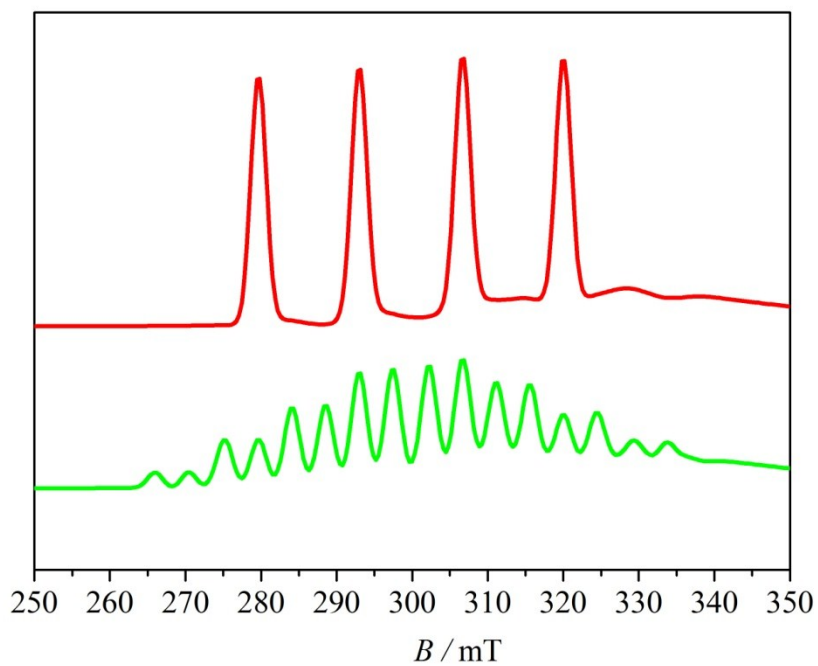


Figure S3. Hyperfine patterns from an axial $S=1/2$ system in the $g_{||}$ ($>g_{\perp}$) region with three $I=3/2$ nuclei according to model I ($A_1 = -a/3$, $A_2 = A_3 = +2a/3$; green line) and model II ($A_1 = a$, $A_2 = A_3 = 0$; red line). $g_{||} = 2.25$; $a_{||} = 425$ MHz. For the simulations an intrinsic line-width of $\sigma_L = 0.8$ mT was assumed.

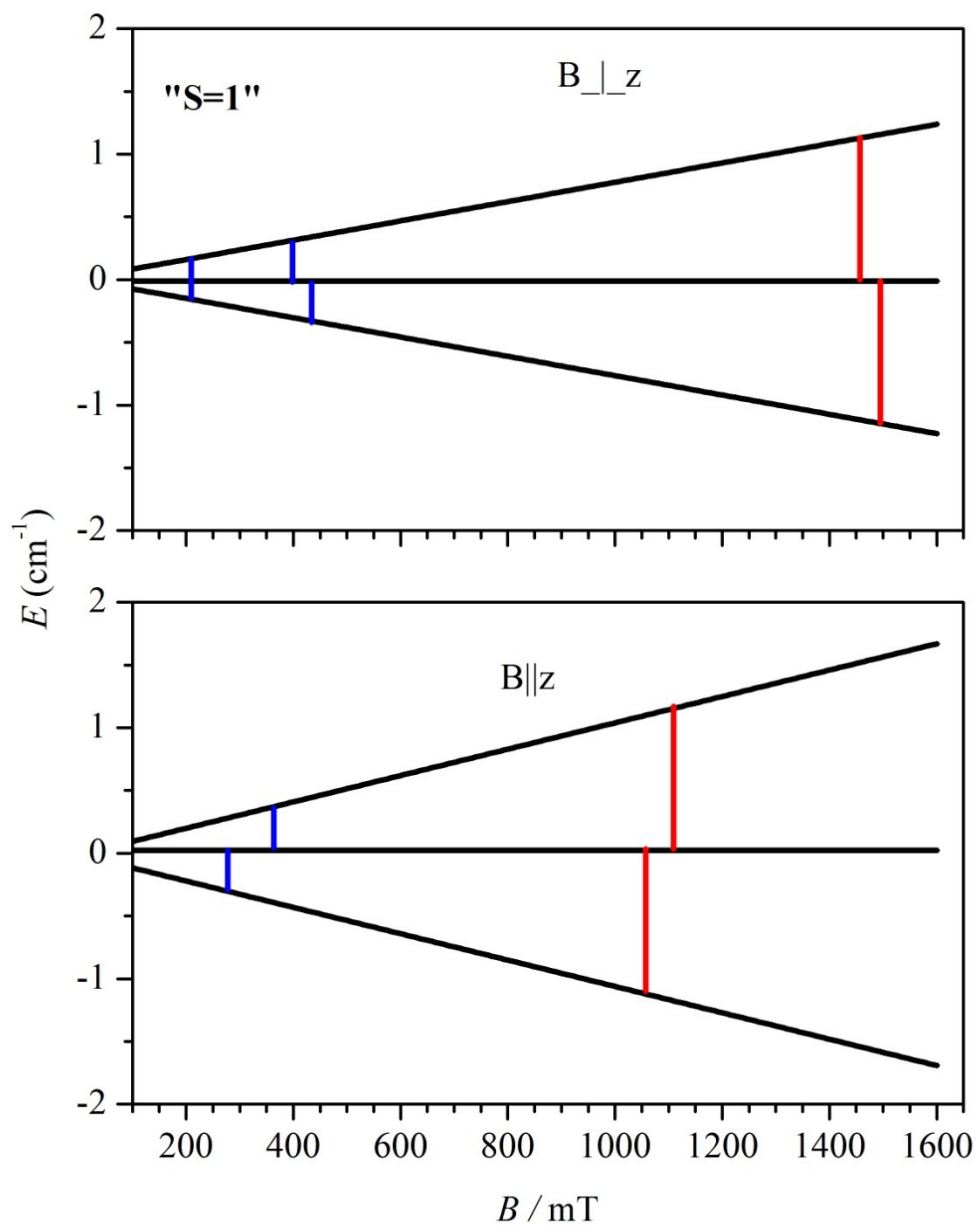


Figure S4. Dependence of the energy levels on the magnetic field for an effective $S = 1$ system characterized by the spin Hamiltonian $H_{zfs} = DS_z^2 + \beta \mathbf{B} \mathbf{g} \mathbf{S}$ for two orientations of the magnetic field relative to the z-axis of the \mathbf{D} -tensor. $D = -0.033$ cm^{-1} , $\mathbf{g} = [1.65, 1.65, 2.25]$. The blue and red vertical bars indicate the EPR transitions at X- and Q-band respectively.

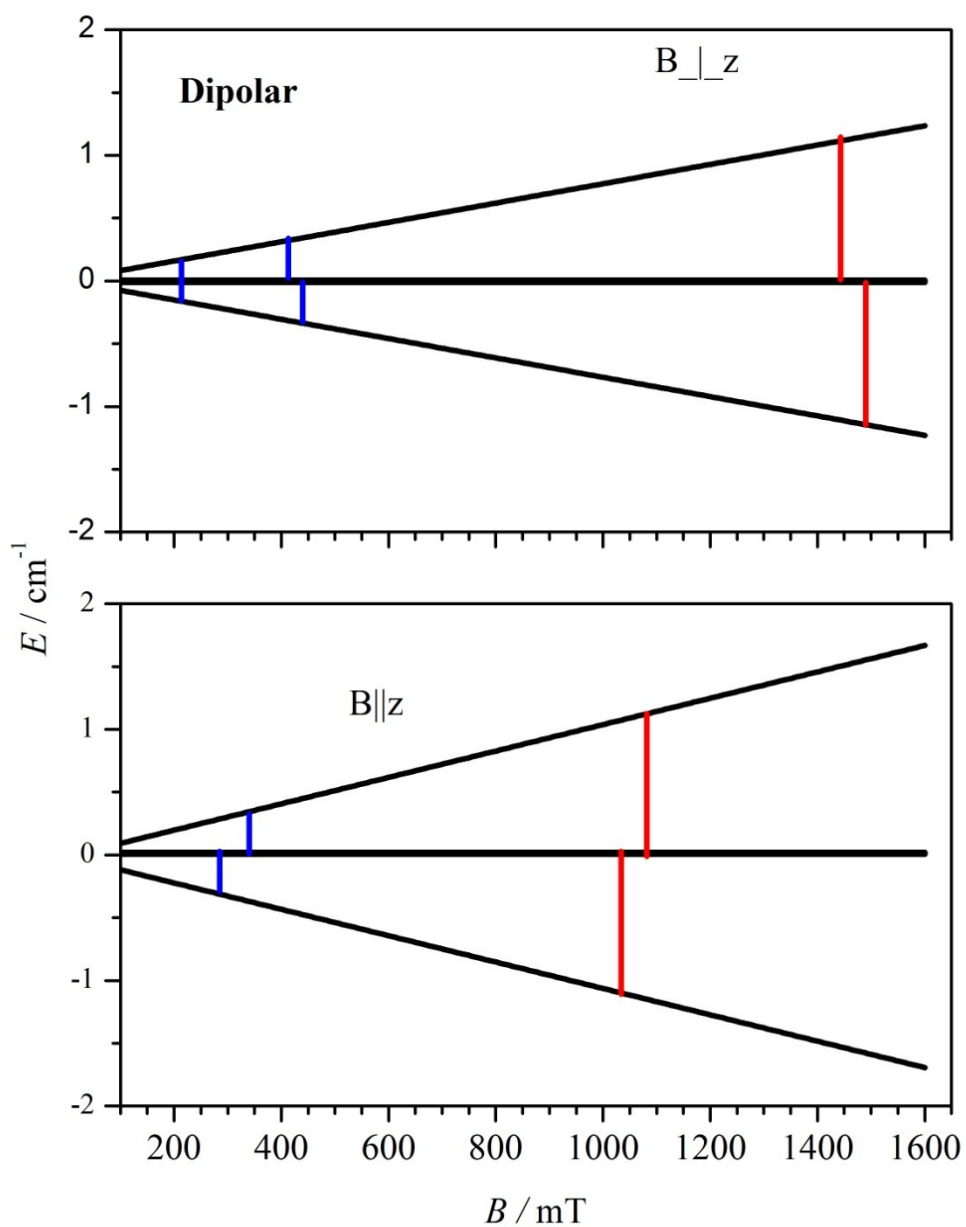


Figure S5. Dependence of the energy levels on the magnetic field for two spins with $S_a = S_b = 1/2$ in the point dipolar approximation. $\mathbf{g}_{a,b} = [1.65, 1.65, 2.25]$, $r_{ab} = 4.4 \text{ \AA}$. The two \mathbf{g} -tensors are assumed co-linear. The blue and red vertical bars indicate the EPR transitions at X- and Q-band respectively.