

ELECTRONIC SUPPLEMENTARY INFORMATION:

**Ab initio quantum-chemical computations of the absorption cross sections in HgX₂
and HgXY (X, Y=Cl, Br, and I): Molecules of interest on the Earth's atmosphere**

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(Dated: 18 November 2018)

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S.I. BASIC INFORMATION ON THE HgX₂ AND HgXY EQUILIBRIUM STRUCTURES

TABLE S.I: Bond lengths in the optimized structures and corresponding vibrational frequencies of the normal modes of the HgX₂ compounds at selected level of theories (both from the literature and used in this work). In all cases, linear structures of HgX₂ have been found to be the most stable ones.

| Method | R_{HgX} [Å] | ν_{sym} [cm ⁻¹] | ν_{asym} [cm ⁻¹] | ν_{bend} [cm ⁻¹] |
|---|---------------------|---------------------------------|----------------------------------|----------------------------------|
| HgCl ₂ | | | | |
| PBE0/aug-cc-pVTZ-PP ^a | 2.257 | 355.5 | 411.4 | 101.5 |
| PBE0-SO/aug-cc-pVTZ-PP ^a | 2.253 | 362.3 | 414.6 | 101.9 |
| CCSD(T)/CBS(TQ) ^b | 2.2509 | 363.0 | 417.0 | 102.2 |
| CCSD(T)-CV/CBS(TQ) ^b | 2.2450 | 367.0 | 420.3 | 104.2 |
| CCSD(T)-CV-SR/CBS(TQ) ^b | 2.2447 | 366.7 | 420.1 | 104.2 |
| CASPT2(12,10)/ANO-RCC-VQZP ^c | 2.230 | 373.8 | 430.6 | 106.8 |
| Electron Diffraction ^d | 2.240, 2.252, 2.275 | | | |
| UV gas phase ^e | | 365,366 | | 70 |
| Raman and IR gas phase ^f | | 355,358 | 413 | 100 |
| Raman and IR matrix ^g | | 358.4 | 407.5 | 107 |
| HgBr ₂ | | | | |
| PBE0/aug-cc-pVTZ-PP ^a | 2.392 | 218.9 | 289.6 | 69.0 |
| PBE0-SO/aug-cc-pVTZ-PP ^a | 2.389 | 224.6 | 295.4 | 65.7 |
| PBE0/Def2QZVP ^c | 2.392 | 218.8 | 288.3 | 68.6 |
| CCSD/Def2QZVP ^c | 2.382 | 228.2 | 297.2 | 70.1 |
| CCSD(T)/CBS(TQ) ^b | 2.3853 | 224.6 | 295.3 | 68.5 |
| CCSD(T)-CV/CBS(TQ) ^b | 2.3764 | 227.5 | 298.4 | 70.1 |
| CCSD(T)-CV-SO/CBS(TQ) ^b | 2.3770 | 226.9 | 297.8 | 69.6 |
| CASPT2(12,10)/ANO-RCC-VTZP ^c | 2.374 | 228.6 | 276.3 | 70.2 |
| CASPT2(12,10)/ANO-RCC-VQZP ^c | 2.363 | 239.9 | 303.8 | 71.5 |
| Electron Diffraction ^h | 2.374, 2.41, 2.378 | | | |
| UV gas phase ⁱ | | 229,221 | 295 | 41,53 |
| Raman and IR gas phase ^f | | 220,218,221.8 | 293 | 68 |
| Raman and IR matrix ^j | | 225 | 294 | 73 |
| HgI ₂ | | | | |
| PBE0/aug-cc-pVTZ-PP ^a | 2.578 | 156.9 | 231.7 | 52.8 |
| PBE0-SO/aug-cc-pVTZ-PP ^a | 2.578 | 164.2 | 239.9 | 51.4 |
| CCSD(T)/aug-cc-pVTZ-PP ^a | 2.568 | 160.1 | 235.6 | 53.8 |
| CASPT2(12,10)/ANO-RCC-VQZP ^c | 2.527 | 174.9 | 255.7 | 56.8 |
| Electron Diffraction ^k | 2.554, 2.5587 | | | |
| UV gas phase ^l | | 156 | 235 | 33,46 |
| Raman and IR gas phase ^m | | 155,158.4 | 237 | 51 |
| Raman and IR matrix ^j | | 163.5 | 237.5 | 63 |

^a Ref.¹

^b Ref.²

^c This work

^d Ref.³⁻⁵

^e Ref.^{6,7}

^f Ref.⁸⁻¹⁰

^g Ref.^{11,12}

^h Ref.^{3,5,13}

ⁱ Ref.^{6,14}

^j Ref.^{11,12,15}

^k Ref.^{3,16}

^m Ref.^{9,10,17,18}

ⁿ Ref.¹⁵

^l Ref.^{6,19}

TABLE S.II: Bond lengths in the optimized structures and corresponding vibrational frequencies of the normal modes of the HgXY compounds at selected level of theories (both from the literature and used in this work). In all cases, linear structures of HgXY have been found to be the most stable ones.

| Method | R_{HgX} [Å] | R_{HgY} [Å] | ν_{sym} [cm ⁻¹] | ν_{asym} [cm ⁻¹] | ν_{bend} [cm ⁻¹] |
|---|---------------|---------------|---------------------------------|----------------------------------|----------------------------------|
| HgBrCl | | | | | |
| CCSD(T)/CBS(TQ) ^a | 2.3761 | 2.2603 | 257.7 | 390.3 | 85.7 |
| CCSD(T)-CV/CBS(TQ) ^a | 2.3674 | 2.2542 | 260.8 | 393.9 | 87.4 |
| CCSD(T)-CV-SO/CBS(TQ) ^a | 2.3682 | 2.2539 | 260.2 | 393.6 | 87.0 |
| CCSD(T)-CV-SO-SR/CBS(TQ) ^a | 2.3685 | 2.2535 | 260.1 | 393.4 | 87.0 |
| CASPT2(12,10)/ANO-RCC-VQZP ^b | 2.353 | 2.240 | 271.0 | 411.1 | 88.3 |
| Raman and IR gas phase ^c | | | 253 | 385 | 83 |
| Raman and IR matrix ^d | | | 255 | 387.5 | 89,91 |
| HgClI | | | | | |
| CASPT2(12,10)/ANO-RCC-VQZP ^b | 2.251 | 2.507 | 211.8 | 394.3 | 80.1 |
| Raman and IR gas phase ^c | | | 196 | 371 | 74 |
| Raman and IR matrix ^d | | | 201.5 | 378.0 | 85,83.5 |
| HgBrI | | | | | |
| CASPT2(12,10)/ANO-RCC-VQZP ^b | 2.374 | 2.515 | 187.9 | 285.8 | 62.9 |
| Raman and IR gas phase ^c | | | 182 | 266 | 60 |
| Raman and IR matrix ^d | | | 187.6 | 272.0 | 66 |

^a Ref.²

^b This work

^c Ref.^{9,10}

^d Ref.^{11,12}

S.II. ADDITIONAL TABLES

A. Allowed vertical electronic transitions of HgX₂ and HgXY

TABLE S.III: Symmetry-allowed vertical electronic transitions of HgCl₂ computed at the DKH3-SOC-MS-CASPT2(12,10)/ANO-RCC-VQZP level of theory at the ground-state equilibrium geometry, including vertical excitations energies in eV (ΔE) and oscillator strengths (f) for the spin-free (SF) and spin-orbit (SO) states (left and right parts, respectively), the main natural orbital excitations which characterize the SF states (columns 2 and 3, only weights larger than 10% are given, for more information see Table S.VIII), and the main SF states contributing to the SO states (columns 8 and 9, only weights larger than 2% are provided).

| SF state | Weight [%] | Excitation | ΔE [eV] | f [au] | SO state/ Ω state | Weight [%] | SF state | ΔE [eV] | f [au] |
|-----------------|------------|--|-----------------|----------|---------------------------------|------------|-----------------|-----------------|----------|
| $1^1\Sigma_g^+$ | 92 | | 0.00 | | $1^1\Sigma_{g0+}^+$ (1) 0_g^+ | 100 | $1^1\Sigma_g^+$ | 0.00 | |
| $1^3\Pi_u$ | 83 | $\pi_u^{nb} \rightarrow \sigma_g^{+*}$ | 5.85 | | $1^3\Pi_{u1}$ (1) 1_u | 77 | $1^3\Pi_u$ | 5.84 | 0.00517 |
| | | | | | | 22 | $1^1\Pi_u$ | | |
| | | | | | $1^3\Pi_{u0+}$ (1) 0_u^+ | 100 | $1^3\Pi_u$ | 5.92 | 0.00001 |
| $1^1\Pi_u$ | 84 | $\pi_u^{nb} \rightarrow \sigma_g^{+*}$ | 5.92 | 0.02415 | $1^1\Pi_{u1}$ (2) 1_u | 77 | $1^1\Pi_u$ | 5.95 | 0.01880 |
| | | | | | | 22 | $1^3\Pi_u$ | | |
| $1^3\Sigma_u^+$ | 88 | $\sigma_u^{+nb} \rightarrow \sigma_g^{+*}$ | 6.22 | | $1^3\Sigma_{u1}^+$ (3) 1_u | 99 | $1^3\Sigma_u^+$ | 6.24 | 0.00020 |
| $1^1\Sigma_u^+$ | 66 | $\pi_g^{nb} \rightarrow \pi_u$ | 6.96 | 0.07999 | $1^1\Sigma_{u0+}^+$ (2) 0_u^+ | 91 | $1^1\Sigma_u^+$ | 6.95 | 0.06352 |
| | 22 | $\sigma_u^{+nb} \rightarrow \sigma_g^{+*}$ | | | | 8 | $1^3\Sigma_u^-$ | | |
| $2^3\Sigma_u^+$ | 88 | $\pi_g^{nb} \rightarrow \pi_u$ | 7.02 | | $2^3\Sigma_{u1}^+$ (4) 1_u | 74 | $2^3\Sigma_u^+$ | 6.88 | 3.2E-06 |
| | | | | | | 25 | $1^3\Sigma_u^-$ | | |
| $1^3\Delta_u$ | 88 | $\pi_g^{nb} \rightarrow \pi_u$ | 7.18 | | $1^3\Delta_{u1}$ (5) 1_u | 100 | $1^3\Delta_u$ | 7.02 | 0.00003 |
| $1^3\Sigma_u^-$ | 88 | $\pi_g^{nb} \rightarrow \pi_u$ | 7.32 | | $1^3\Sigma_{u0+}^-$ (3) 0_u^+ | 91 | $1^3\Sigma_u^-$ | 7.36 | 0.02594 |
| | | | | | | 8 | $1^1\Sigma_u^+$ | | |
| | | | | | $1^3\Sigma_{u1}^-$ (6) 1_u | 74 | $1^3\Sigma_u^-$ | 7.49 | 6.3E-07 |
| | | | | | | 26 | $2^3\Sigma_u^+$ | | |
| $2^1\Sigma_u^+$ | 72 | $\sigma_u^{+nb} \rightarrow \sigma_g^{+*}$ | 8.61 | 2.03991 | $2^1\Sigma_{u0+}^+$ (4) 0_u^+ | 50 | $2^1\Sigma_u^+$ | 8.53 | 0.96067 |
| | 24 | $\pi_g^{nb} \rightarrow \pi_u$ | | | | 50 | $2^3\Pi_u$ | | |
| | | | | | $3^1\Sigma_{u0+}^+$ (5) 0_u^+ | 50 | $2^1\Sigma_u^+$ | 8.73 | 1.07039 |
| | | | | | | 50 | $2^3\Pi_u$ | | |

TABLE S.IV: Symmetry-allowed vertical electronic transitions of HgI₂ computed at the DKH3-SOC-MS-CASPT2(12,10)/ANO-RCC-VQZP level of theory at the ground-state equilibrium geometry, including vertical excitations energies in eV (ΔE) and oscillator strengths (f) for the spin-free (SF) and spin-orbit (SO) states (left and right parts, respectively), the main natural orbital excitations which characterize the SF states (columns 2 and 3, only weights larger than 10% are given, for more information see Table S.IX), and the main SF states contributing to the SO states (columns 8 and 9, only weights larger than 2% are provided).

| SF state | Weight [%] | Excitation | ΔE [eV] | f [au] | SO state/ Ω state | Weight [%] | SF state | ΔE [eV] | f [au] |
|-----------------|------------|--|-----------------|----------|---------------------------------|------------|-----------------|-----------------|----------|
| $1^1\Sigma_g^+$ | 93 | | 0.00 | | $1^1\Sigma_{g0+}^+$ (1) 0_g^+ | 99 | $1^1\Sigma_g^+$ | 0.00 | |
| $1^3\Pi_u$ | 86 | $\pi_u^{nb} \rightarrow \sigma_g^{+\star}$ | 4.53 | | $1^3\Pi_{u1}$ (1) 1_u | 65 | $1^3\Pi_u$ | 4.35 | 0.00440 |
| | | | | | | 32 | $1^1\Pi_u$ | | |
| | | | | | $1^3\Pi_{u0+}$ (1) 0_u^+ | 96 | $1^3\Pi_u$ | 4.85 | 0.04102 |
| | | | | | | 3 | $1^1\Sigma_u^+$ | | |
| $1^1\Pi_u$ | 87 | $\pi_u^{nb} \rightarrow \sigma_g^{+\star}$ | 4.65 | 0.01306 | ^a (2) 1_u | 48 | $1^3\Sigma_u^+$ | 4.57 | 0.00416 |
| | | | | | | 39 | $1^1\Pi_u$ | | |
| | | | | | | 12 | $1^3\Pi_u$ | | |
| $1^3\Sigma_u^+$ | 83 | $\sigma_u^{+nb} \rightarrow \sigma_g^{+\star}$ | 4.87 | | ^a (3) 1_u | 48 | $1^3\Sigma_u^+$ | 5.29 | 0.00407 |
| | | | | | | 27 | $1^1\Pi_u$ | | |
| | | | | | | 19 | $1^3\Pi_u$ | | |
| $1^1\Sigma_u^+$ | 66 | $\pi_g^{nb} \rightarrow \pi_u$ | 5.73 | 0.08338 | $1^1\Sigma_{u0+}^+$ (2) 0_u^+ | 94 | $1^1\Sigma_u^+$ | 5.82 | 0.05127 |
| | 25 | $\sigma_u^{+nb} \rightarrow \sigma_g^{+\star}$ | | | | 4 | $1^3\Pi_u$ | | |
| | | | | | | 2 | $1^3\Sigma_u^-$ | | |
| $2^3\Sigma_u^+$ | 83 | $\pi_g^{nb} \rightarrow \pi_u$ | 5.83 | | $2^3\Sigma_{u1}^+$ (4) 1_u | 58 | $2^3\Sigma_u^+$ | 5.48 | 0.00027 |
| | | | | | | 35 | $1^3\Sigma_u^-$ | | |
| | | | | | | 3 | $1^3\Sigma_u^+$ | | |
| | | | | | | 3 | $1^3\Pi_u$ | | |
| $1^3\Delta_u$ | 89 | $\pi_g^{nb} \rightarrow \pi_u$ | 5.96 | | $1^3\Delta_{u1}$ (5) 1_u | 100 | $1^3\Delta_u$ | 6.06 | 0.00280 |
| $1^3\Sigma_u^-$ | 89 | $\pi_g^{nb} \rightarrow \pi_u$ | 6.09 | | $1^3\Sigma_{u0+}^-$ (3) 0_u^+ | 97 | $1^3\Sigma_u^-$ | 6.14 | 0.00723 |
| | | | | | | 2 | $1^1\Sigma_u^+$ | | |
| | | | | | $1^3\Sigma_{u1}^-$ (6) 1_u | 61 | $1^3\Sigma_u^-$ | 6.57 | 0.00383 |
| | | | | | | 37 | $2^3\Sigma_u^+$ | | |
| $2^1\Sigma_u^+$ | 69 | $\sigma_u^{+nb} \rightarrow \sigma_g^{+\star}$ | 7.35 | 3.18911 | $2^1\Sigma_{u0+}^+$ (4) 0_u^+ | 100 | $2^1\Sigma_u^+$ | 7.44 | 3.15275 |
| | 26 | $\pi_g^{nb} \rightarrow \pi_u$ | | | | | | | |

^a Strong mixing between $1^3\Sigma_u^+$ and $1^1\Pi_u$ (in favour for $1^3\Sigma_u^+$) prevents assigning the spatial symmetry of this SO state unambiguously.

TABLE S.V: Symmetry-allowed vertical electronic transitions of HgCII computed at the DKH3-SOC-MS-CASPT2(12,10)/ANO-RCC-VQZP level of theory at the ground-state equilibrium geometry, including vertical excitations energies in eV (ΔE) and oscillator strengths (f) for the spin-free (SF) and spin-orbit (SO) states (left and right parts, respectively), the corresponding SF states of $D_{\infty h}$ symmetry, the main natural orbital excitations which characterize the SF states (columns 3 and 4, only weights larger than 10% are given, for more information see Table S.XI), and the main SF states contributing to the SO states (columns 9 and 10, only weights larger than 2% are provided).

| SF state | (D_{2h}) | Weight [%] | Excitation | ΔE [eV] | f [au] | SO state/ Ω state | Weight [%] | SF state | ΔE [eV] | f [au] |
|---------------|---------------------|------------|---|-----------------|----------|------------------------------|------------|---------------|-----------------|----------|
| $1^1\Sigma^+$ | ($1^1\Sigma_g^+$) | 95 | | 0.00 | | $1^1\Sigma_{0+}^+$ (1) 0_+ | 100 | $1^1\Sigma^+$ | 0.00 | |
| $1^3\Pi$ | ($1^3\Pi_g$) | 88 | $\pi_2^{nb} \rightarrow \sigma^{+\ast}$ | 4.45 | | $1^3\Pi_1$ (1) 1 | 76 | $1^3\Pi$ | 4.33 | 0.00086 |
| | | | | | | | 23 | $1^1\Pi$ | | |
| | | | | | | $1^3\Pi_{0+}$ (2) 0_+ | 100 | $1^3\Pi$ | 4.77 | 0.01033 |
| $1^1\Pi$ | ($1^1\Pi_g$) | 80 | $\pi_2^{nb} \rightarrow \sigma^{+\ast}$ | 4.76 | 0.00274 | $1^1\Pi_1$ (2) 1 | 68 | $1^1\Pi$ | 4.85 | 0.00174 |
| | | | | | | | 19 | $1^3\Pi$ | | |
| | | | | | | | 13 | $1^3\Sigma^+$ | | |
| $1^3\Sigma^+$ | ($1^3\Sigma_u^+$) | 69 | $\sigma^{+nb} \rightarrow \sigma^{+\ast}$ | 5.48 | | $1^3\Sigma_1^+$ (4) 1 | 61 | $1^3\Sigma^+$ | 5.62 | 0.00022 |
| | | 16 | $\pi_2^{nb} \rightarrow \pi$ | | | | 14 | $2^3\Pi$ | | |
| | | | | | | | 9 | $2^1\Pi$ | | |
| | | | | | | | 7 | $1^1\Pi$ | | |
| | | | | | | | 3 | $2^3\Sigma^+$ | | |
| $2^3\Pi$ | ($1^3\Pi_u$) | 87 | $\pi_1^{nb} \rightarrow \sigma^{+\ast}$ | 5.56 | | $2^3\Pi_1$ (3) 1 | 77 | $2^3\Pi$ | 5.55 | 0.00036 |
| | | | | | | | 19 | $1^3\Sigma^+$ | | |
| | | | | | | | 2 | $2^1\Pi$ | | |
| | | | | | | $2^3\Pi_{0+}$ (3) 0_+ | 100 | $2^3\Pi$ | 5.66 | 0.00307 |
| $2^1\Pi$ | ($1^1\Pi_u$) | 79 | $\pi_1^{nb} \rightarrow \sigma^{+\ast}$ | 5.82 | 0.01120 | $2^1\Pi_1$ (5) 1 | 75 | $2^1\Pi$ | 5.87 | 0.00871 |
| | | | | | | | 13 | $2^3\Sigma^+$ | | |
| | | | | | | | 8 | $1^3\Sigma^-$ | | |
| $2^3\Sigma^+$ | ($2^3\Sigma_u^+$) | 49 | $\pi_2^{nb} \rightarrow \pi$ | 6.21 | | $2^3\Sigma_1^+$ (6) 1 | 49 | $2^3\Sigma^+$ | 5.89 | 0.00152 |
| | | 23 | $\sigma^{+nb} \rightarrow \sigma^{+\ast}$ | | | | 25 | $1^3\Sigma^-$ | | |
| | | 18 | $\pi_1^{nb} \rightarrow \pi$ | | | | 13 | $2^1\Pi$ | | |
| | | | | | | | 5 | $2^3\Pi$ | | |
| | | | | | | | 5 | $1^3\Sigma^+$ | | |
| $2^1\Sigma^+$ | ($1^1\Sigma_u^+$) | 59 | $\pi_2^{nb} \rightarrow \pi$ | 6.28 | 0.03850 | $2^1\Sigma_{0+}^+$ (4) 0_+ | 97 | $2^1\Sigma^+$ | 6.31 | 0.02758 |
| | | 20 | $\sigma^{+nb} \rightarrow \sigma^{+\ast}$ | | | | 2 | $3^3\Pi$ | | |
| $1^3\Delta$ | ($1^3\Delta_u$) | 88 | $\pi_2^{nb} \rightarrow \pi$ | 6.36 | | $1^3\Delta_1$ (7) 1 | 94 | $1^3\Delta$ | 6.38 | 0.00449 |
| | | | | | | | 3 | $3^3\Pi$ | | |
| | | | | | | | 2 | $3^1\Pi$ | | |
| $1^3\Sigma^-$ | ($1^3\Sigma_u^-$) | 88 | $\pi_2^{nb} \rightarrow \pi$ | 6.53 | | $1^3\Sigma_{0+}^-$ (5) 0_+ | 96 | $1^3\Sigma^-$ | 6.53 | 0.00005 |
| | | | | | | | 4 | $3^3\Pi$ | | |
| | | | | | | $1^3\Sigma_1^-$ (8) 1 | 39 | $1^3\Sigma^-$ | 6.85 | 0.00248 |
| | | | | | | | 34 | $3^3\Sigma^+$ | | |
| | | | | | | | 21 | $2^3\Sigma^+$ | | |
| | | | | | | | 3 | $3^1\Pi$ | | |

Table continued on next page

TABLE S.V: Continue. . .

| SF state | (D_{2h}) | Weight [%] | Excitation | ΔE [eV] | f [au] | SO state/ Ω state | Weight [%] | SF state | ΔE [eV] | f [au] |
|---------------|-------------------|------------|--|-----------------|---------|------------------------------|------------|---------------|-----------------|---------|
| $3^3\Sigma^+$ | $(1^3\Sigma_g^+)$ | 41 | $\pi_1^{nb} \rightarrow \pi$ | 6.89 | | $3^3\Sigma_1^+$ (9) 1 | 60 | $3^3\Sigma^+$ | 6.94 | 0.00172 |
| | | 26 | $\sigma_1^+ \rightarrow \sigma^{++}$ | | | | 22 | $1^3\Sigma^-$ | | |
| | | 23 | $\pi_2^{nb} \rightarrow \pi$ | | | | 11 | $2^3\Sigma^+$ | | |
| | | | | | | | 3 | $3^3\Pi$ | | |
| | | | | | | | 2 | $2^3\Sigma^-$ | | |
| $3^1\Sigma^+$ | $(2^1\Sigma_g^+)$ | 57 | $\pi_1^{nb} \rightarrow \pi$ | 7.34 | 0.31756 | $3^1\Sigma_{0+}^+$ (6) 0_+ | 83 | $3^1\Sigma^+$ | 7.33 | 0.21391 |
| | | 19 | $\sigma_1^+ \rightarrow \sigma^{++}$ | | | | 11 | $2^3\Sigma^-$ | | |
| | | | | | | | 6 | $3^3\Pi$ | | |
| $2^3\Delta$ | $(1^3\Delta_g)$ | 85 | $\pi_1^{nb} \rightarrow \pi$ | 7.46 | | $2^3\Delta_1$ (10) 1 | 71 | $2^3\Delta$ | 7.30 | 0.00100 |
| | | | | | | | 15 | $4^3\Sigma^+$ | | |
| | | | | | | | 8 | $2^3\Sigma^-$ | | |
| | | | | | | | 3 | $3^3\Pi$ | | |
| $4^3\Sigma^+$ | $(2^3\Sigma_g^+)$ | 61 | $\sigma_1^+ \rightarrow \sigma^{++}$ | 7.48 | | $4^3\Sigma_1^+$ (11) 1 | 45 | $4^3\Sigma^+$ | 7.31 | 0.00023 |
| | | 25 | $\pi_1^{nb} \rightarrow \pi$ | | | | 26 | $2^3\Sigma^-$ | | |
| | | | | | | | 25 | $2^3\Delta$ | | |
| | | | | | | | 2 | $3^3\Sigma^+$ | | |
| $2^3\Sigma^-$ | $(1^3\Sigma_g^-)$ | 85 | $\pi_1^{nb} \rightarrow \pi$ | 7.60 | | $2^3\Sigma_1^-$ (8) 0_+ | 73 | $2^3\Sigma^-$ | 7.66 | 0.61485 |
| | | | | | | | 15 | $4^1\Sigma^+$ | | |
| | | | | | | | 10 | $3^1\Sigma^+$ | | |
| $3^3\Pi$ | $(2^3\Pi_u)$ | 52 | $\sigma_1^+ \rightarrow \pi$ | 7.71 | | $3^3\Pi_1$ (12) 1 | 55 | $3^3\Pi$ | 7.67 | 0.01509 |
| | | 40 | $\sigma^{+nb} \rightarrow \pi$ | | | | 33 | $3^1\Pi$ | | |
| | | | | | | | 6 | $1^3\Delta$ | | |
| | | | | | | | 4 | $2^3\Delta$ | | |
| | | | | | | $3^3\Pi_{0+}$ (7) 0_+ | 84 | $3^3\Pi$ | 7.61 | 0.03067 |
| | | | | | | | 6 | $3^1\Sigma^+$ | | |
| | | | | | | | 4 | $4^1\Sigma^+$ | | |
| | | | | | | | 4 | $1^3\Sigma^-$ | | |
| $4^1\Sigma^+$ | $(3^1\Sigma_g^+)$ | 41 | $\sigma_1^+ \rightarrow \sigma^{++}$ | 7.72 | 1.86775 | $4^1\Sigma_{0+}^+$ (9) 0_+ | 80 | $4^1\Sigma^+$ | 7.81 | 1.32056 |
| | | 25 | $\sigma^{+nb} \rightarrow \sigma^{++}$ | | | | 14 | $2^3\Sigma^-$ | | |
| | | 11 | $\pi_1^{nb} \rightarrow \pi$ | | | | 4 | $3^3\Pi$ | | |

TABLE S.VI: Symmetry-allowed vertical electronic transitions of HgBrI computed at the DKH3-SOC-MS-CASPT2(12,10)/ANO-RCC-VQZP level of theory at the ground-state equilibrium geometry, including vertical excitations energies in eV (ΔE) and oscillator strengths (f) for the spin-free (SF) and spin-orbit (SO) states (left and right parts, respectively), the corresponding SF states of $D_{\infty h}$ symmetry, the main natural orbital excitations which characterize the SF states (columns 3 and 4, only weights larger than 10% are given, for more information see Table S.XII), and the main SF states contributing to the SO states (columns 9 and 10, only weights larger than 2% are provided).

| SF state | (D_{2h}) | Weight [%] | Excitation | ΔE [eV] | f [au] | SO state/ Ω state | Weight [%] | SF state | ΔE [eV] | f [au] |
|---------------|---------------------|------------|---|-----------------|----------|------------------------------|------------|---------------|-----------------|----------|
| $1^1\Sigma^+$ | ($1^1\Sigma_g^+$) | 92 | | 0.00 | | $1^1\Sigma_{0+}^+$ (1) 0_+ | 99 | $1^1\Sigma^+$ | 0.00 | |
| $1^3\Pi$ | ($1^3\Pi_g$) | 88 | $\pi_2^{nb} \rightarrow \sigma^{+\ast}$ | 4.27 | | $1^3\Pi_1$ (1) 1 | 65 | $1^3\Pi$ | 4.10 | 0.00073 |
| | | | | | | | 34 | $1^1\Pi$ | | |
| | | | | | | $1^3\Pi_{0+}$ (2) 0_+ | 100 | $1^3\Pi$ | 4.57 | 0.00471 |
| $1^1\Pi$ | ($1^1\Pi_g$) | 86 | $\pi_2^{nb} \rightarrow \sigma^{+\ast}$ | 4.44 | 0.00151 | $1^1\Pi_1$ (2) 1 | 58 | $1^1\Pi$ | 4.60 | 0.00121 |
| | | | | | | | 29 | $1^3\Pi$ | | |
| | | | | | | | 11 | $1^3\Sigma^+$ | | |
| $2^3\Pi$ | ($1^3\Pi_u$) | 86 | $\pi_1^{nb} \rightarrow \sigma^{+\ast}$ | 4.98 | | $2^3\Pi_1$ (3) 1 | 71 | $2^3\Pi$ | 4.88 | 0.00462 |
| | | | | | | | 25 | $2^1\Pi$ | | |
| | | | | | | | 3 | $1^3\Sigma^+$ | | |
| $2^1\Pi$ | ($1^1\Pi_u$) | 85 | $\pi_1^{nb} \rightarrow \sigma^{+\ast}$ | 5.10 | 0.01888 | $2^3\Pi_{0+}$ (3) 0_+ | 100 | $2^3\Pi$ | 5.19 | 0.01382 |
| | | | | | | $2^1\Pi_1$ (4) 1 | 47 | $2^1\Pi$ | 5.06 | 0.00687 |
| | | | | | | | 39 | $1^3\Sigma^+$ | | |
| | | | | | | | 8 | $2^3\Pi$ | | |
| | | | | | | | 4 | $1^1\Pi$ | | |
| | | | | | | | 2 | $1^3\Pi$ | | |
| $1^3\Sigma^+$ | ($1^3\Sigma_u^+$) | 77 | $\sigma^{+nb} \rightarrow \sigma^{+\ast}$ | 5.15 | | $1^3\Sigma_1^+$ (5) 1 | 46 | $1^3\Sigma^+$ | 5.45 | 0.00672 |
| | | | | | | | 28 | $2^1\Pi$ | | |
| | | | | | | | 19 | $2^3\Pi$ | | |
| | | | | | | | 3 | $1^1\Pi$ | | |
| | | | | | | | 2 | $1^3\Pi$ | | |
| $2^1\Sigma^+$ | ($1^1\Sigma_u^+$) | 65 | $\pi_2^{nb} \rightarrow \pi$ | 6.04 | 0.07532 | $2^1\Sigma_{0+}^+$ (4) 0_+ | 97 | $2^1\Sigma^+$ | 6.08 | 0.05780 |
| | | 20 | $\sigma^{+nb} \rightarrow \sigma^{+\ast}$ | | | | | | | |
| $2^3\Sigma^+$ | ($2^3\Sigma_u^+$) | 72 | $\pi_2^{nb} \rightarrow \pi$ | 6.09 | | $2^3\Sigma_1^+$ (6) 1 | 64 | $2^3\Sigma^+$ | 5.77 | 0.00002 |
| | | | | | | | 34 | $1^3\Sigma^-$ | | |
| $1^3\Delta$ | ($1^3\Delta_u$) | 86 | $\pi_2^{nb} \rightarrow \pi$ | 6.24 | | $1^3\Delta_1$ (7) 1 | 94 | $1^3\Delta$ | 6.28 | 0.00264 |
| | | | | | | | 2 | $3^3\Pi$ | | |
| | | | | | | | 2 | $3^1\Pi$ | | |
| $1^3\Sigma^-$ | ($1^3\Sigma_u^-$) | 87 | $\pi_2^{nb} \rightarrow \pi$ | 6.40 | | $1^3\Sigma_{0+}^-$ (5) 0_+ | 96 | $1^3\Sigma^-$ | 6.41 | 0.00099 |
| | | | | | | | 3 | $3^3\Pi$ | | |
| | | | | | | $1^3\Sigma_1^-$ (9) 1 | 58 | $1^3\Sigma^-$ | 6.76 | 0.00315 |
| | | | | | | | 32 | $2^3\Sigma^+$ | | |
| | | | | | | | 3 | $3^3\Sigma^+$ | | |
| | | | | | | | 3 | $3^1\Pi$ | | |

Table continued on next page

TABLE S.VI: Continue...

| SF state | (D_{2h}) | Weight [%] | Excitation | ΔE [eV] | f [au] | SO state/ Ω state | Weight [%] | SF state | ΔE [eV] | f [au] | | |
|---------------|-------------------|------------|--------------------------------------|-----------------|---------|------------------------------|------------|--|-----------------|---------|----|---------------|
| $3^3\Sigma^+$ | $(1^3\Sigma_g^+)$ | 67 | $\pi_1^{nb} \rightarrow \pi$ | 6.58 | | $3^3\Sigma_1^+$ (8) 1 | 74 | $3^3\Sigma^+$ | 6.46 | 2.5E-06 | | |
| | | 12 | $\sigma_1^+ \rightarrow \sigma^{+*}$ | | | | 20 | $2^3\Sigma^-$ | | | | |
| $2^3\Delta$ | $(1^3\Delta_g)$ | 87 | $\pi_1^{nb} \rightarrow \pi$ | 6.91 | | $2^3\Delta_1$ (10) 1 | 3 | $4^3\Sigma^+$ | 6.81 | 0.00004 | | |
| | | | | | | | 76 | $2^3\Delta$ | | | | |
| | | | | | | | 9 | $3^3\Pi$ | | | | |
| $3^1\Sigma^+$ | $(2^1\Sigma_g^+)$ | 69 | $\pi_1^{nb} \rightarrow \pi$ | 7.02 | 0.13343 | $3^1\Sigma_{0+}^+$ (6) 0_+ | 6 | $3^1\Pi$ | 6.97 | 0.05845 | | |
| | | | | | | | 10 | $\sigma_1^+ \rightarrow \sigma^{+*}$ | | | 4 | $4^3\Sigma^+$ |
| | | | | | | | 26 | | | | 21 | $3^3\Pi$ |
| $4^3\Sigma^+$ | $(2^3\Sigma_g^+)$ | 69 | $\sigma_1^+ \rightarrow \sigma^{+*}$ | 7.05 | | $4^3\Sigma_1^+$ (11) 1 | 52 | $4^3\Sigma^+$ | 6.95 | 0.00027 | | |
| | | | | | | | 14 | $\pi_1^{nb} \rightarrow \pi$ | | | 18 | $2^3\Sigma^-$ |
| | | | | | | | | | | | 16 | $3^3\Sigma^+$ |
| $2^3\Sigma^-$ | $(1^3\Sigma_g^-)$ | 86 | $\pi_1^{nb} \rightarrow \pi$ | 7.09 | | $2^3\Sigma_1^-$ (7) 0_+ | 9 | $2^3\Delta$ | 7.12 | 0.09376 | | |
| | | | | | | | 54 | $2^3\Sigma^-$ | | | | |
| | | | | | | | 41 | $3^1\Sigma^+$ | | | | |
| $3^3\Pi$ | $(2^3\Pi_g)$ | 82 | $\sigma^{+nb} \rightarrow \pi$ | 7.39 | | $3^3\Pi_1$ (12) 1 | 4 | $3^3\Pi$ | 7.34 | 0.00003 | | |
| | | | | | | | 34 | $3^3\Pi$ | | | | |
| | | | | | | | 38 | $2^3\Sigma^-$ | | | | |
| | | | | | | | 16 | $4^3\Sigma^+$ | | | | |
| | | | | | | | 6 | $3^3\Sigma^+$ | | | | |
| | | | | | | | 2 | $2^3\Delta$ | | | | |
| | | | | | | | | | | | | |
| $4^1\Sigma^+$ | $(3^1\Sigma_g^+)$ | 37 | $\sigma_1^+ \rightarrow \sigma^{+*}$ | 7.61 | 2.57561 | $4^1\Sigma_{0+}^+$ (9) 0_+ | 70 | $3^3\Pi$ | 7.38 | 0.01211 | | |
| | | | | | | | 29 | $\sigma^{+nb} \rightarrow \sigma^{+*}$ | | | 19 | $2^3\Sigma^-$ |
| | | | | | | | 20 | $\pi_2^{nb} \rightarrow \pi$ | | | 6 | $3^1\Sigma^+$ |
| | | | | | | | | | | | 2 | $1^3\Sigma^-$ |
| | | | | | | | 98 | $4^1\Sigma^+$ | 7.67 | 2.53150 | | |

B. Multiconfigurational nature of the spin-free states of HgX_2 and HgXY

TABLE S.VII: Main electronic configurations of the most relevant low-lying electronic SF states in the UV-Vis spectroscopy of HgCl_2 determined with the DKH3-CASSCF(12,10)/ANO-RCC-VQZP methodology at the DKH3-CASPT2(12,10)/ANO-RCC-VQZP optimized geometry. Only configuration of weights larger than 2% are given.

| SF state | Configuration | Weight [%] | SF state | Configuration | Weight [%] |
|--|--|--|--|--|------------|
| $1^1\Sigma_g^+$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^0 \sigma_u^{+0}$ | 92.5 | $1^1\Sigma_g^+$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^0 \sigma_u^{+0}$ | 92.5 |
| | $\sigma_g^{+1} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 2.3 | | $\sigma_g^{+1} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 2.3 |
| $1^1\Sigma_u^+$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 66.4 | $2^1\Sigma_g^+$ | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 55.9 |
| | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 21.7 | | $\sigma_g^{+1} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 30.7 |
| | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 3.5 | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb0} \sigma_g^{+*2} \pi_u^0 \sigma_u^{+0}$ | 3.1 | |
| | $\sigma_g^{+1} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 3.5 | $1^1\Pi_g$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 85.1 |
| | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 2.2 | | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*2} \pi_u^0 \sigma_u^{+0}$ | 4.1 |
| $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 72.4 | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | | 4.0 | |
| $2^1\Sigma_u^+$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 24.2 | $1^1\Sigma_g^-$ | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 87.1 |
| | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 84.3 | | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 5.8 |
| $1^1\Pi_u$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*2} \pi_u^0 \sigma_u^{+0}$ | 6.1 | $1^1\Delta_g$ | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 86.8 |
| | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^0 \sigma_u^{+1}$ | 3.2 | | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 5.8 |
| $1^1\Sigma_u^-$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 87.9 | $1^3\Sigma_g^+$ | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 54.6 |
| | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 3.7 | | $\sigma_g^{+1} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 32.9 |
| $1^1\Delta_u$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 88.3 | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 4.7 | |
| | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 3.1 | $2^3\Sigma_g^+$ | $\sigma_g^{+1} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 55.5 |
| $1^3\Sigma_u^+$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 87.8 | | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 32.0 |
| | $\sigma_g^{+1} \pi_u^{nb4} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*2} \pi_u^0 \sigma_u^{+0}$ | 4.6 | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 2.9 | |
| $2^3\Sigma_u^+$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 87.7 | $1^3\Pi_g$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 87.6 |
| | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 4.5 | | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*2} \pi_u^0 \sigma_u^{+0}$ | 2.7 |
| $1^3\Pi_u$ | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*1} \pi_u^0 \sigma_u^{+0}$ | 83.0 | $1^3\Sigma_g^-$ | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 85.4 |
| | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*2} \pi_u^0 \sigma_u^{+0}$ | 4.3 | | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 6.8 |
| $1^3\Sigma_u^-$ | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^0 \sigma_u^{+1}$ | 3.9 | $1^3\Delta_g$ | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 85.8 |
| | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 87.7 | | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 6.8 |
| $1^3\Delta_u$ | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 4.5 | | | |
| | $\sigma_g^{+2} \pi_u^{nb4} \pi_g^{nb3} \sigma_u^{+nb2} \sigma_g^{+*0} \pi_u^1 \sigma_u^{+0}$ | 87.7 | | | |
| | $\sigma_g^{+2} \pi_u^{nb3} \pi_g^{nb4} \sigma_u^{+nb1} \sigma_g^{+*1} \pi_u^1 \sigma_u^{+0}$ | 4.5 | | | |

TABLE S.VIII: Main electronic configurations of the most relevant low-lying electronic SF states in the UV-Vis spectroscopy of HgBr₂ determined with the DKH3-CASSCF(12,10)/ANO-RCC-VQZP methodology at the DKH3-CASPT2(12,10)/ANO-RCC-VQZP optimized geometry. Only configuration of weights larger than 2% are given.

| SF state | Configuration | Weight [%] | SF state | Configuration | Weight [%] |
|--|--|------------|--|--|------------|
| 1 ¹ Σ _g ⁺ | σ _g ⁺² π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ⁰ σ _u ⁺⁰ | 92.8 | 1 ¹ Σ _g ⁺ | σ _g ⁺² π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ⁰ σ _u ⁺⁰ | 92.8 |
| 1 ¹ Σ _u ⁺ | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 65.1 | 2 ¹ Σ _g ⁺ | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 63.5 |
| | σ _g ⁺² π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 23.4 | | σ _g ⁺¹ π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 22.8 |
| | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 3.6 | | σ _g ⁺² π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb0} σ _g ^{+*2} π _u ⁰ σ _u ⁺⁰ | 2.7 |
| | σ _g ⁺¹ π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 2.8 | 1 ¹ Π _g | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 85.6 |
| | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 2.2 | | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*2} π _u ⁰ σ _u ⁺⁰ | 4.2 |
| 2 ¹ Σ _u ⁺ | σ _g ⁺² π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 70.5 | | σ _g ⁺² π _u ^{nb3} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 3.8 |
| | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 25.8 | 1 ¹ Σ _g ⁻ | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 87.9 |
| 1 ¹ Π _u | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 85.6 | | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 5.3 |
| | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb1} σ _g ^{+*2} π _u ⁰ σ _u ⁺⁰ | 6.1 | 1 ¹ Δ _g | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 87.6 |
| | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ⁰ σ _u ⁺¹ | 2.6 | | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 5.3 |
| 1 ¹ Σ _u ⁻ | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 88.2 | 1 ³ Σ _g ⁺ | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 68.5 |
| | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 3.6 | | σ _g ⁺¹ π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 19.1 |
| 1 ¹ Δ _u | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 88.6 | | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 5.3 |
| | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 3.1 | 2 ³ Σ _g ⁺ | σ _g ⁺¹ π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 68.2 |
| 1 ³ Σ _u ⁺ | σ _g ⁺² π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 87.7 | | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 18.8 |
| | σ _g ⁺¹ π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*2} π _u ⁰ σ _u ⁺⁰ | 4.2 | | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 3.5 |
| | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 2.5 | 1 ³ Π _g | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 88.0 |
| 2 ³ Σ _u ⁺ | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 87.1 | | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*2} π _u ⁰ σ _u ⁺⁰ | 2.6 |
| | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 4.5 | 1 ³ Σ _g ⁻ | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 86.1 |
| | σ _g ⁺² π _u ^{nb4} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 2.3 | | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 6.4 |
| 1 ³ Π _u | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*1} π _u ⁰ σ _u ⁺⁰ | 84.5 | 1 ³ Δ _g | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 86.6 |
| | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb1} σ _g ^{+*2} π _u ⁰ σ _u ⁺⁰ | 4.2 | | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 6.5 |
| | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ⁰ σ _u ⁺¹ | 3.3 | | | |
| 1 ³ Σ _u ⁻ | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 88.0 | | | |
| | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 4.5 | | | |
| 1 ³ Δ _u | σ _g ⁺² π _u ^{nb4} π _g ^{nb3} σ _u ^{+nb2} σ _g ^{+*0} π _u ¹ σ _u ⁺⁰ | 89.3 | | | |
| | σ _g ⁺² π _u ^{nb3} π _g ^{nb4} σ _u ^{+nb1} σ _g ^{+*1} π _u ¹ σ _u ⁺⁰ | 4.4 | | | |

TABLE S.IX: Main electronic configurations of the most relevant low-lying electronic SF states in the UV-Vis spectroscopy of Hg₂ determined with the DKH3-CASSCF(12,10)/ANO-RCC-VQZP methodology at the DKH3-CASPT2(12,10)/ANO-RCC-VQZP optimized geometry. Only configuration of weights larger than 2% are given.

| SF state | Configuration | Weight [%] | SF state | Configuration | Weight [%] |
|--|--|------------|--|--|------------|
| 1 ¹ Σ _g ⁺ | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ⁰ σ _u ⁺ 0 | 93.0 | 1 ¹ Σ _g ⁺ | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ⁰ σ _u ⁺ 0 | 93.0 |
| 1 ¹ Σ _u ⁺ | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 65.6 | 2 ¹ Σ _g ⁺ | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 72.6 |
| | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 24.7 | | σ _g ⁺ 1 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 14.0 |
| | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 2.9 | 1 ¹ Π _g | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 86.4 |
| 2 ¹ Σ _u ⁺ | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 68.8 | | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *2 π _u ⁰ σ _u ⁺ 0 | 3.9 |
| | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 25.8 | | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 3.6 |
| 1 ¹ Π _u | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 87.0 | 1 ¹ Σ _g ⁻ | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 89.2 |
| | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb1 σ _g ⁺ *2 π _u ⁰ σ _u ⁺ 0 | 5.7 | | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 4.3 |
| 1 ¹ Σ _u ⁻ | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 88.9 | 1 ¹ Δ _g | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 89.0 |
| | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 3.2 | | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 4.4 |
| 1 ¹ Δ _u | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 89.0 | 1 ³ Σ _g ⁺ | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 82.0 |
| | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 3.2 | | σ _g ⁺ 1 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 6.5 |
| 1 ³ Σ _u ⁺ | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 83.1 | | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 5.1 |
| | σ _g ⁺ 1 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *2 π _u ⁰ σ _u ⁺ 0 | 3.4 | 2 ³ Σ _g ⁺ | σ _g ⁺ 1 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 78.7 |
| | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 7.8 | | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 6.4 |
| 2 ³ Σ _u ⁺ | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 82.5 | | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 4.1 |
| | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 3.8 | 1 ³ Π _g | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 88.2 |
| | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 7.5 | | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *2 π _u ⁰ σ _u ⁺ 0 | 2.1 |
| 1 ³ Π _u | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *1 π _u ⁰ σ _u ⁺ 0 | 86.3 | 1 ³ Σ _g ⁻ | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 87.5 |
| | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb1 σ _g ⁺ *2 π _u ⁰ σ _u ⁺ 0 | 3.7 | | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 5.5 |
| | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ⁰ σ _u ⁺ 1 | 2.4 | 1 ³ Δ _g | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 87.9 |
| 1 ³ Σ _u ⁻ | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 89.0 | | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 5.5 |
| | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 3.9 | | | |
| 1 ³ Δ _u | σ _g ⁺ 2 π _u ^{nb} 4 π _g ^{nb} 3 σ _u ⁺ nb2 σ _g ⁺ *0 π _u ¹ σ _u ⁺ 0 | 89.1 | | | |
| | σ _g ⁺ 2 π _u ^{nb} 3 π _g ^{nb} 4 σ _u ⁺ nb1 σ _g ⁺ *1 π _u ¹ σ _u ⁺ 0 | 4.0 | | | |

TABLE S.X: Main electronic configurations of the most relevant low-lying electronic SF states in the UV-Vis spectroscopy of HgBrCl determined with the DKH3-CASSCF(12,10)/ANO-RCC-VQZP methodology at the DKH3-CASPT2(12,10)/ANO-RCC-VQZP optimized geometry. Only configuration of weights larger than 2% are given.

| SF state | Configuration | Weight [%] | SF state | Configuration | Weight [%] |
|-------------------------------|--|------------|-------------------------------|--|------------|
| 1 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^0 \sigma_2^{+0}$ | 94.6 | 1 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 85.5 |
| 2 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 62.7 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 3.9 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 23.2 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 2} \pi^0 \sigma_2^{+0}$ | 2.3 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 3.6 | 2 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 72.3 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 2.7 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 11.1 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 2.2 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 5.4 |
| 3 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 54.7 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 2.1 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 27.9 | 3 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 38.4 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb0} \sigma^{+\ast 2} \pi^0 \sigma_2^{+0}$ | 3.0 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 37.7 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 3.0 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 12.9 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 2.0 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 2.6 |
| 4 ¹ Σ ⁺ | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 50.0 | 4 ³ Σ ⁺ | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 48.6 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 26.2 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 38.2 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 4.2 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 3.0 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb0} \sigma^{+\ast 2} \pi^0 \sigma_2^{+0}$ | 3.7 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 2.1 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 3.6 | 1 ³ Π | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 85.8 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^0 \sigma_2^{+0}$ | 2.5 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 3.4 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 2.0 | 2 ³ Π | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 83.4 |
| 1 ¹ Π | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 84.3 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 3.5 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 3.4 | 3 ³ Π | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 62.5 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 2.7 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 28.6 |
| 2 ¹ Π | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 82.5 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 2.9 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 2} \pi^0 \sigma_2^{+0}$ | 4.2 | 1 ³ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 79.6 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 1} \pi^0 \sigma_2^{+0}$ | 3.3 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 10.3 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+\ast 2} \pi^0 \sigma_2^{+0}$ | 2.1 | 2 ³ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 76.1 |
| 1 ¹ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 88.6 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 10.5 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 2.7 | | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 3.8 |
| 2 ¹ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 88.5 | 1 ³ Δ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 80.1 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 3.9 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 9.9 |
| 1 ¹ Δ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 88.8 | 2 ³ Δ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 77.0 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 2.7 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 10.2 |
| 2 ¹ Δ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 0} \pi^1 \sigma_2^{+0}$ | 88.0 | | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 3.9 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+\ast 1} \pi^1 \sigma_2^{+0}$ | 3.8 | | | |

TABLE S.XI: Main electronic configurations of the most relevant low-lying electronic SF states in the UV-Vis spectroscopy of HgCII determined with the DKH3-CASSCF(12,10)/ANO-RCC-VQZP methodology at the DKH3-CASPT2(12,10)/ANO-RCC-VQZP optimized geometry. Only configuration of weights larger than 2% are given.

| SF state | Configuration | Weight [%] | SF state | Configuration | Weight [%] |
|-------------------------------|--|------------|-------------------------------|--|------------|
| 1 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^0 \sigma_2^{+0}$ | 94.5 | 1 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 68.9 |
| 2 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 59.3 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 15.6 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 20.5 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 6.1 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 7.9 | 2 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 49.3 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.4 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 22.8 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 2.7 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 17.5 |
| 3 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 57.3 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.7 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 19.1 | 3 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 40.6 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 5.6 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 26.1 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 6.2 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 23.0 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb0} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 2.2 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.5 |
| 4 ¹ Σ ⁺ | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 41.0 | 4 ³ Σ ⁺ | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 60.6 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 25.0 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 24.8 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 10.7 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 3.0 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 9.4 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.3 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb0} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 3.1 | 1 ³ Π | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 87.9 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.4 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 2.2 |
| 1 ¹ Π | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 80.4 | 2 ³ Π | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 87.4 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 7.3 | 3 ³ Π | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 51.9 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 2.6 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 39.5 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.4 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.6 |
| 2 ¹ Π | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 79.0 | 1 ³ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 87.7 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 7.3 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.6 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 2.9 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 2.6 |
| | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 2.3 | 2 ³ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 84.6 |
| | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 2.3 | | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 4.0 |
| 3 ¹ Π | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 58.6 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 2.5 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 30.8 | 1 ³ Δ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 88.2 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 6.1 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.7 |
| 1 ¹ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 86.7 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 2.1 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.7 | 2 ³ Δ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 85.4 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.5 | | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 4.1 |
| 2 ¹ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 87.0 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 2.1 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.1 | | | |
| 1 ¹ Δ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 87.0 | | | |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.7 | | | |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.5 | | | |
| 2 ¹ Δ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 87.0 | | | |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.1 | | | |

TABLE S.XII: Main electronic configurations of the most relevant low-lying electronic SF states in the UV-Vis spectroscopy of HgBrI determined with the DKH3-CASSCF(12,10)/ANO-RCC-VQZP methodology at the DKH3-CASPT2(12,10)/ANO-RCC-VQZP optimized geometry. Only configuration of weights larger than 2% are given.

| SF state | Configuration | Weight [%] | SF state | Configuration | Weight [%] |
|-------------------------------|--|------------|-------------------------------|--|------------|
| 1 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^0 \sigma_2^{+0}$ | 91.7 | 1 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 76.9 |
| 2 ² Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 65.1 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 8.2 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 19.6 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 3.9 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 4.1 | 2 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 71.8 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.2 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 3.8 |
| 3 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 69.1 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 9.4 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 10.3 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.3 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 6.4 | 3 ³ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 66.7 |
| 4 ¹ Σ ⁺ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 29.3 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 9.5 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 36.8 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 12.1 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 19.9 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.8 |
| | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 2.7 | 4 ³ Σ ⁺ | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 69.1 |
| 1 ¹ Π | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 86.0 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 13.5 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 3.7 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.9 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.8 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 3.5 |
| 2 ¹ Π | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 84.7 | | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.1 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 5.7 | 1 ³ Π | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 88.1 |
| 1 ¹ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 87.8 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 2.2 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 3.0 | 2 ³ Π | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^0 \sigma_2^{+0}$ | 85.7 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.7 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*2} \pi^0 \sigma_2^{+0}$ | 3.0 |
| 2 ¹ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 88.7 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^0 \sigma_2^{+1}$ | 2.2 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 4.4 | 3 ³ Π | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 82.3 |
| 1 ¹ Δ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 87.9 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 9.3 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.0 | | $\sigma_1^{+1} \pi_1^{nb4} \pi_2^{nb4} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.5 |
| | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 2.6 | 1 ³ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 86.8 |
| 2 ¹ Δ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 88.5 | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 3.0 |
| | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 4.4 | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.4 |
| | | | 2 ³ Σ ⁻ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 86.5 |
| | | | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 4.2 |
| | | | | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.1 |
| | | | 1 ³ Δ | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 86.4 |
| | | | | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 3.4 |
| | | | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 3.2 |
| | | | 2 ³ Δ | $\sigma_1^{+2} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*0} \pi^1 \sigma_2^{+0}$ | 86.5 |
| | | | | $\sigma_1^{+2} \pi_1^{nb4} \pi_2^{nb3} \sigma^{+nb1} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 4.1 |
| | | | | $\sigma_1^{+1} \pi_1^{nb3} \pi_2^{nb4} \sigma^{+nb2} \sigma^{+*1} \pi^1 \sigma_2^{+0}$ | 2.1 |

C. Extended data on the energetics of the excited states of HgX_2 and HgXY

TABLE S.XIII: Main vertical excitation energies (ΔE) in eV and associated oscillator strength (f) values obtained for symmetrical HgX_2 compounds from the computations of the SF states (SOC not considered) with the DKH3-MS-CASPT2(12,10)/ANO-RCC-VQZP methodology at equilibrium geometries (see computational details in Section II.A). Energetic ordering of states corresponds to the one of HgBr_2 .

| SF state | HgCl_2 | | HgBr_2 | | HgI_2 | |
|-----------------|-----------------|----------|-----------------|----------|-----------------|----------|
| | ΔE [eV] | f [au] | ΔE [eV] | f [au] | ΔE [eV] | f [au] |
| $1^1\Sigma_g^+$ | 0.00 | | 0.00 | | 0.00 | |
| $1^3\Pi_g$ | 5.15 | | 4.57 | | 4.04 | |
| $1^1\Pi_g$ | 5.23 | | 4.72 | | 4.17 | |
| $1^3\Pi_u$ | 5.85 | | 5.17 | | 4.53 | |
| $1^1\Pi_u$ | 5.92 | 0.02415 | 5.29 | 0.02444 | 4.65 | 0.01306 |
| $1^3\Sigma_u^+$ | 6.22 | | 5.5 | | 4.87 | |
| $1^1\Sigma_u^+$ | 6.96 | 0.07999 | 6.36 | 0.09646 | 5.73 | 0.08338 |
| $2^3\Sigma_u^+$ | 7.02 | | 6.45 | | 5.83 | |
| $1^3\Delta_u$ | 7.18 | | 6.59 | | 5.96 | |
| $1^1\Delta_u$ | 7.20 | | 6.62 | | 5.96 | |
| $1^1\Sigma_u^-$ | 7.25 | | 6.67 | | 6.01 | |
| $1^3\Sigma_g^+$ | 7.00 | | 6.71 | | 6.12 | |
| $1^3\Sigma_u^-$ | 7.32 | | 6.74 | | 6.09 | |
| $1^3\Delta_g$ | 7.86 | | 7.15 | | 6.41 | |
| $2^1\Sigma_g^+$ | 7.59 | | 7.18 | | 6.66 | |
| $2^3\Sigma_g^+$ | 7.86 | | 7.22 | | 6.80 | |
| $1^1\Delta_g$ | 7.94 | | 7.25 | | 6.54 | |
| $1^1\Sigma_g^-$ | 7.99 | | 7.29 | | 6.52 | |
| $1^3\Sigma_g^-$ | 8.05 | | 7.34 | | 6.57 | |
| $2^3\Pi_g$ | 8.59 | | 7.87 | | 7.03 | |
| $2^1\Pi_g$ | 8.58 | | 7.88 | | 7.09 | |
| $3^1\Sigma_g^+$ | 8.55 | | 8.08 | | 7.61 | |
| $2^1\Sigma_u^+$ | 8.61 | 2.03991 | 8.08 | 2.52704 | 7.35 | 3.18911 |

TABLE S.XIV: Main vertical excitation energies (ΔE) in eV and associated oscillator strength (f) values obtained for mixed HgXY compounds from the computations of the SF states (SOC not considered) with the DKH3-MS-CASPT2(12,10)/ANO-RCC-VQZP methodology at equilibrium geometries (see computational details in Section II.A). Energetic ordering of states corresponds to the one of HgBrCl.

| SF state | HgBrCl | | HgClI | | HgBrI | |
|---------------|-----------------|----------|-----------------|----------|-----------------|----------|
| | ΔE [eV] | f [au] | ΔE [eV] | f [au] | ΔE [eV] | f [au] |
| $1^1\Sigma^+$ | 0.00 | | 0.00 | | 0.00 | |
| $1^3\Pi$ | 4.86 | | 4.45 | | 4.27 | |
| $1^1\Pi$ | 5.15 | 0.00318 | 4.76 | 0.00274 | 4.44 | 0.00151 |
| $2^3\Pi$ | 5.70 | | 5.56 | | 4.98 | |
| $1^3\Sigma^+$ | 5.84 | | 5.48 | | 5.15 | |
| $2^1\Pi$ | 6.00 | 0.01676 | 5.82 | 0.01120 | 5.10 | 0.01888 |
| $2^3\Sigma^+$ | 6.66 | | 6.21 | | 6.09 | |
| $2^1\Sigma^+$ | 6.66 | 0.05102 | 6.28 | 0.03850 | 6.04 | 0.07532 |
| $1^3\Delta$ | 6.80 | | 6.36 | | 6.24 | |
| $1^1\Delta$ | 6.85 | | 6.44 | | 6.26 | |
| $1^1\Sigma^-$ | 6.92 | | 6.45 | | 6.32 | |
| $3^3\Sigma^+$ | 6.95 | | 6.89 | | 6.58 | |
| $1^3\Sigma^-$ | 6.99 | | 6.53 | | 6.40 | |
| $3^1\Sigma^+$ | 7.51 | 0.00683 | 7.34 | 0.31756 | 7.02 | 0.13343 |
| $2^3\Delta$ | 7.62 | | 7.46 | | 6.91 | |
| $4^3\Sigma^+$ | 7.64 | | 7.48 | | 7.05 | |
| $2^1\Delta$ | 7.72 | | 7.53 | | 6.98 | |
| $2^1\Sigma^-$ | 7.76 | | 7.54 | | 7.03 | |
| $2^3\Sigma^-$ | 7.81 | | 7.60 | | 7.09 | |
| $4^1\Sigma^+$ | 8.14 | 1.54965 | 7.72 | 1.86775 | 7.61 | 2.57561 |
| $3^3\Pi$ | 8.18 | | 7.71 | | 7.39 | |

TABLE S.XV: Main vertical excitation energies (ΔE) in eV and associated oscillator strength (f) values obtained for symmetrical HgX₂ compounds from the computations of the SO states with the DKH3-SOC-MS-CASPT2(12,10)/ANO-RCC-VQZP methodology at equilibrium geometries (see computational details in Section II.A).

| SO state | HgCl ₂ | | HgBr ₂ | | HgI ₂ | |
|--|-------------------|----------|-------------------|----------|------------------|----------|
| | ΔE [eV] | f [au] | ΔE [eV] | f [au] | ΔE [eV] | f [au] |
| (1) 0 _g ⁺ | 0.00 | | 0.00 | | 0.00 | |
| (1) 2 _g | 5.12 | | 4.43 | | 3.79 | |
| (1) 1 _g | 5.14 | | 4.49 | | 3.85 | |
| (1) 0 _g ⁻ | 5.21 | | 4.74 | | 4.39 | |
| (2) 0 _g ⁺ | 5.21 | | 4.75 | | 4.42 | |
| (2) 1 _g | 5.26 | | 4.83 | | 4.45 | |
| (1) 2 _u | 5.81 | | 5.04 | | 4.30 | |
| (1) 1 _u | 5.84 | 0.00517 | 5.09 | 0.00663 | 4.35 | 0.00440 |
| (1) 0 _u ⁻ | 5.91 | | 5.26 | | 4.53 | |
| (2) 1 _u | 5.95 | 0.01880 | 5.32 | 0.01296 | 4.57 | 0.00416 |
| (1) 0 _u ⁺ | 5.92 | 0.00001 | 5.34 | 0.00645 | 4.85 | 0.04102 |
| (2) 0 _u ⁻ | 6.24 | | 5.60 | | 5.27 | |
| (3) 1 _u | 6.24 | 0.00020 | 5.62 | 0.00474 | 5.29 | 0.00407 |
| (4) 1 _u | 6.88 | 3.2E-06 | 6.21 | 1.9E-06 | 5.48 | 0.00027 |
| (2) 2 _u | 6.94 | | 6.25 | | 5.47 | |
| (2) 0 _u ⁺ | 6.95 | 0.06352 | 6.39 | 0.09004 | 5.82 | 0.05127 |
| (3) 0 _u ⁻ | 6.95 | | 6.45 | | 5.87 | |
| (5) 1 _u | 7.02 | 0.00003 | 6.53 | 0.00053 | 6.06 | 0.00280 |
| (3) 1 _g | 7.00 | | 6.65 | | 5.87 | |
| (1) 3 _u | 7.36 | | 6.69 | | 5.96 | |
| (4) 0 _u ⁻ | 7.35 | | 6.71 | | 6.07 | |
| (2) 0 _g ⁻ | 7.01 | | 6.73 | | 6.14 | |
| (3) 0 _u ⁺ | 7.36 | 0.02594 | 6.77 | 0.00389 | 6.14 | 0.00723 |
| (2) 2 _g | 7.65 | | 6.85 | | 6.00 | |
| (3) 2 _u | 7.44 | | 6.98 | | 6.52 | |
| (6) 1 _u | 7.49 | 6.3E-07 | 7.03 | 0.00084 | 6.57 | 0.00383 |
| (4) 1 _g | 7.72 | | 7.05 | | 6.30 | |
| (5) 1 _g | 7.73 | | 7.09 | | 6.69 | |
| (3) 0 _g ⁺ | 7.59 | | 7.18 | | 6.49 | |
| (3) 0 _g ⁻ | 7.80 | | 7.19 | | 6.44 | |
| (1) 3 _g | 8.03 | | 7.24 | | 6.41 | |
| (4) 0 _g ⁻ | 8.09 | | 7.33 | | 6.81 | |
| (4) 0 _g ⁺ | 8.04 | | 7.37 | | 6.68 | |
| (3) 2 _g | 8.16 | | 7.55 | | 6.79 | |
| (6) 1 _g | 8.23 | | 7.60 | | 6.91 | |
| (5) 0 _g ⁺ | 8.41 | | 7.74 | | 7.13 | |
| (7) 1 _g | 8.41 | | 7.74 | | 7.17 | |
| (5) 0 _g ⁻ | 8.43 | | 7.75 | | 7.19 | |
| (4) 0 _u ⁺ ^a | 8.53 | 0.96067 | 8.10 | 2.40671 | 7.44 | 3.15275 |

^a Characterize with almost identical weights of the same SF states as state (5) 0_u⁺ ($\Delta E = 8.73$ eV, $f=1.07039$)

TABLE S.XVI: Main vertical excitation energies (ΔE) in eV and associated oscillator strength (f) values, obtained for asymmetrical HgXY compounds from the computations of the SO states with the DKH3-SOC-MS-CASPT2(12,10)/ANO-RCC-VQZP methodology at equilibrium geometries. Additionally, dipole moments of SO states (μ_z) in Debyes are reported (for orientation X-Hg-Y; positive end towards the second halogen Y).

| SO state | Cl-Hg-Br | | | Cl-Hg-I | | | Br-Hg-I | | |
|--------------------|-----------------|----------|-------------|-----------------|----------|-------------|-----------------|----------|-------------|
| | ΔE [eV] | f [au] | μ_z [D] | ΔE [eV] | f [au] | μ_z [D] | ΔE [eV] | f [au] | μ_z [D] |
| (1) 0 ₊ | 0.00 | | 0.65 | 0.00 | | 1.07 | 0.00 | | 0.38 |
| (1) 2 | 4.75 | | 3.84 | 4.21 | | 4.85 | 4.03 | | 4.02 |
| (1) 1 | 4.83 | 0.00053 | 3.69 | 4.33 | 0.00086 | 4.91 | 4.10 | 0.00073 | 4.23 |
| (1) 0 ₋ | 5.00 | | 2.96 | 4.68 | | 4.16 | 4.51 | | 3.44 |
| (2) 0 ₊ | 5.01 | 0.00118 | 2.90 | 4.77 | 0.01033 | 4.09 | 4.57 | 0.00471 | 2.60 |
| (2) 1 | 5.21 | 0.00267 | 3.02 | 4.85 | 0.00174 | 4.31 | 4.60 | 0.00121 | 3.94 |
| (2) 2 | 5.65 | | -3.78 | 5.52 | | -4.71 | 4.84 | | -3.79 |
| (3) 1 | 5.70 | 0.00053 | -3.51 | 5.55 | 0.00036 | -3.70 | 4.88 | 0.00462 | -3.93 |
| (2) 0 ₋ | 5.78 | | -2.85 | 5.56 | | -1.05 | 5.00 | | -2.93 |
| (3) 0 ₊ | 5.79 | 0.00073 | -2.86 | 5.66 | 0.00307 | -4.16 | 5.19 | 0.01382 | -2.45 |
| (3) 0 ₋ | 5.88 | | 1.05 | 5.69 | | -1.43 | 5.43 | | 0.32 |
| (4) 1 | 5.88 | 0.00019 | 0.91 | 5.62 | 0.00022 | 0.67 | 5.06 | 0.00687 | -3.47 |
| (5) 1 | 6.04 | 0.01579 | -2.89 | 5.87 | 0.00871 | -1.90 | 5.45 | 0.00672 | 0.18 |
| (6) 1 | 6.44 | 0.00007 | 4.69 | 5.89 | 0.00152 | 4.03 | 5.77 | 0.00002 | 5.15 |
| (3) 2 | 6.49 | | 4.59 | 5.91 | | 6.37 | 5.79 | | 5.10 |
| (4) 0 ₋ | 6.65 | | 4.33 | 6.23 | | 6.24 | 6.11 | | 5.27 |
| (4) 0 ₊ | 6.66 | 0.04221 | 2.57 | 6.31 | 0.02758 | 4.31 | 6.08 | 0.05780 | 2.49 |
| (7) 1 | 6.73 | 0.00064 | 3.94 | 6.38 | 0.00449 | 5.98 | 6.28 | 0.00264 | 4.62 |
| (5) 0 ₋ | 6.93 | | 0.51 | 6.46 | | 5.97 | 6.33 | | 4.59 |
| (1) 3 | 6.94 | | 4.37 | 6.36 | | 6.58 | 6.23 | | 5.37 |
| (8) 1 | 6.95 | 0.00003 | -1.37 | 6.85 | 0.00248 | 2.78 | 6.46 | 2.5E-06 | -3.73 |
| (6) 0 ₋ | 7.00 | | 1.95 | 6.92 | | -2.26 | 6.62 | | -3.17 |
| (5) 0 ₊ | 7.02 | 0.01252 | 4.18 | 6.53 | 0.00005 | 6.13 | 6.41 | 0.00099 | 4.60 |
| (4) 2 | 7.21 | | 3.89 | 6.87 | | 5.85 | 6.59 | | -3.79 |
| (9) 1 | 7.25 | 0.00051 | 3.75 | 6.94 | 0.00172 | 0.55 | 6.76 | 0.00315 | 4.24 |
| (5) 2 | 7.39 | | -3.73 | 7.22 | | -4.50 | 6.73 | | 4.16 |
| (10) 1 | 7.45 | 1.1E-06 | -3.17 | 7.30 | 0.00100 | -3.85 | 6.81 | 0.00004 | -3.26 |
| (11) 1 | 7.48 | 0.00024 | -3.01 | 7.31 | 0.00023 | -3.58 | 6.95 | 0.00027 | -1.60 |
| (6) 0 ₊ | 7.50 | 0.00499 | -1.97 | 7.33 | 0.21391 | -1.94 | 6.97 | 0.05845 | -1.43 |
| (7) 0 ₋ | 7.56 | | -2.64 | 7.36 | | -2.75 | 6.93 | | -1.87 |
| (2) 3 | 7.83 | | -3.81 | 7.69 | | -4.49 | 7.03 | | -4.17 |
| (7) 0 ₊ | 7.85 | 0.00259 | -3.20 | 7.61 | 0.03067 | 3.07 | 7.12 | 0.09376 | -2.45 |
| (8) 0 ₋ | 7.89 | | -3.64 | 7.63 | | 3.06 | 7.11 | | -2.96 |
| (6) 2 | 7.99 | | -2.92 | 7.81 | | -3.82 | 7.29 | | -3.11 |
| (9) 0 ₋ | 7.99 | | 2.58 | 7.72 | | -4.26 | 7.40 | | 1.29 |
| (8) 0 ₊ | 8.00 | 0.00001 | 2.61 | 7.66 | 0.61485 | -2.65 | 7.38 | 0.01211 | 1.03 |
| (12) 1 | 8.03 | 0.00601 | 1.01 | 7.67 | 0.01509 | 2.95 | 7.34 | 0.00003 | -1.98 |
| (13) 1 | 8.06 | 0.00300 | -2.01 | 7.88 | 0.00425 | -3.57 | 7.43 | 0.00294 | 0.33 |
| (9) 0 ₊ | 8.16 | 1.58677 | 2.45 | 7.81 | 1.32056 | 1.24 | 7.67 | 2.53150 | 1.12 |
| (7) 2 | 8.46 | | 2.63 | 8.09 | | 3.68 | 7.81 | | 1.59 |
| (14) 1 | 8.48 | 0.00566 | 2.02 | 8.13 | 0.01668 | 2.83 | 7.85 | 0.00125 | 1.21 |

S.III. BENCHMARK ON THE ABSORPTION CROSS SECTION SIMULATIONS

In this section, a possible problem of choosing the proper ensemble of geometries \mathbf{R}_k in the absorption cross section simulations is shortly discussed. In this small benchmark, HgBr_2 has been chosen as a representative system. Methodological parameters –active space, symmetry restrictions, number of states per IrRep, IPEA and IMAG shifts, and Gaussian broadening δ – are the same as those reported in Section II.

A. Equilibrium structures used in the generation of the Wigner distributions

The first issue is related to the information on the ground-state equilibrium structures –optimized geometries and vibrational frequencies of the normal modes– used in the process of generating the set of geometries by using the Wigner distribution. In this benchmark, the method for determining the vertical excitation energies ΔE and f has been kept unchanged – they are computed at the same level of theory, namely DKH3–SOC–MS–CASPT2(12,10)/ANO–RCC. The methodologies used to determine the gradients and Hessians for the ground-state structures are the following:

- cheap and low-level DFT with the PBE0/Def2QZVP level, which is a good candidate for the optimizations of structures of the larger mercury-based molecules¹
- DKH3–CASPT2(12,10)/ANO–RCC–V(T/Q)ZP consistent with the further computations of ΔE and f at the same level of theory
- CCSD, which is a typical reference in these type of optimization procedures, along with the Def2QZVP basis set

In the case of DFT/PBE0, CCSD, and CASPT2, the ANO–RCC–VTZP basis set is used and later the results are compared with those provided with the CASPT2 method and the larger ANO–RCC–VQZP basis set.

When comparing just the equilibrium R_{HgBr} bond lengths and vibrational frequencies (see Table S.1), it can be easily observed that the CASPT2(12,10)/ANO–RCC–VQZP methodology is actually performing less accurately than the other protocols. This problem was already highlighted in our first benchmark paper on the spectral properties of IBr and

HgBr₂²⁰ and it is related to the fact that the ground-state wavefunction of HgBr₂ has basically one-reference character and thus CASPT2 performs just as regular MP2 method. However, the final outcome of the usage of CASPT2 geometries in the $\sigma(E)$ simulations is almost the same as in the case of DFT/PBE0 or CCSD methods (see Figure S1). Firstly, the values of $\sigma(E)$ are of the same magnitude for all bands. Secondly, the shifts of the bands are negligible and much less than the uncertainty between the experimental measurements of Roxlo²¹ and Maya.²² Among them, the largest shifts (up to 0.25 eV) are noticed for the first band *a*, which still has very broad and uncertain shape. Small differences are observed for band *b*, for which the use of the ensemble of DFT/PBE0 geometries leads to the sharpest band shape. Only in the case of CCSD/Def2QZVP and CASPT2(12,10)/ANO-RCC-VQZP ensembles the side small right-side peak at ≈ 6.46 eV is noticed. For the DFT/Def2QZVP and CASPT2(12,10)/ANO-RCC-VTZP the peak is less pronounced and becomes a shoulder of band *b*. This has not been reported experimentally, due to the larger experimental broadening of this band. Moreover, a perfect agreement is observed for band *c*, for which in the case of all R_k ensembles, the same position, shape, and intensity are predicted. Lastly, there are not any other differences than those pointed out above if the ANO-RCC-VTZP or ANO-RCC-VQZP basis sets are used in the CASPT2 geometry optimization and ΔE and f computations. This is promising for future computations of these cross sections for larger mercury systems, for which smaller basis set (of a triple- ζ quality) could easily be used.

In conclusion, a good feature of the $\sigma(E)$ simulation technique based on statistical averaging is that, despite noticeable differences in the equilibrium structure properties, the final outcome of the simulation seems to be invariant on the quantum-chemical method used to generate the ensemble of Wigner geometries.

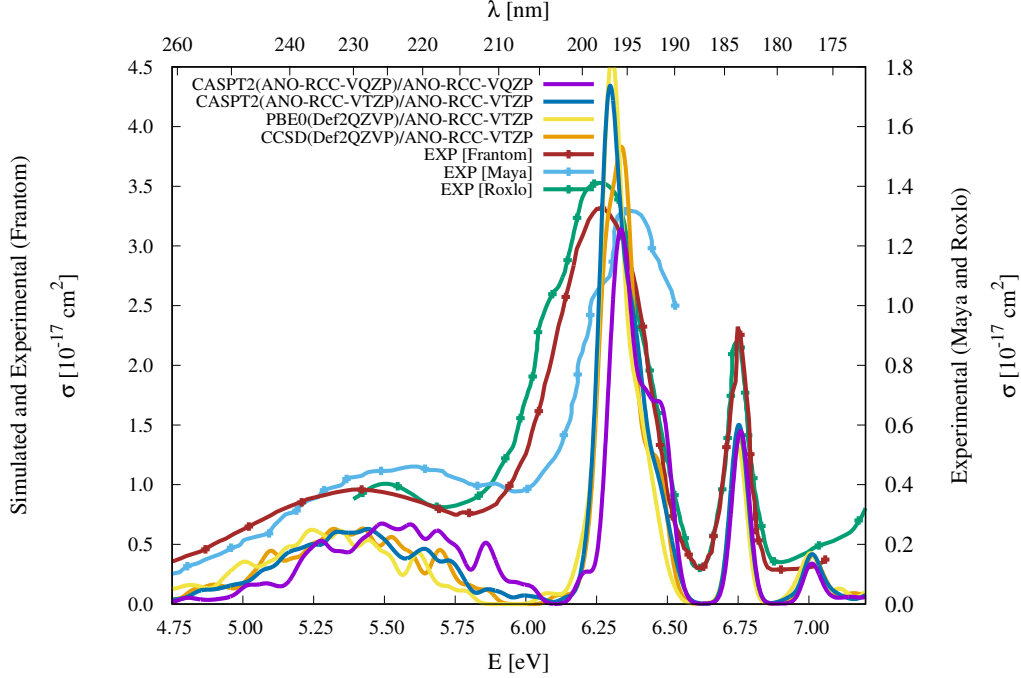


FIG. S.1: Simulated absorption spectra of HgBr_2 arising from different sampling of Wigner geometries. In the legend, the method used in the geometry optimization of the ground state and the computation of the vibrational frequencies is on the left, whereas on the right the basis set used in the DKH3–SOC–MS–CASPT2(12,10) computations of ΔE and f is given. Experimental spectra^{21–23} are included for the comparison (for better comparison, the $\sigma(E)$ values of spectrum of Frantom *et al.*²³ shown on the same axis as theoretical ones).

B. Number of sampled geometries

In this benchmark, the convergence of the obtained absorption cross sections upon increasing the number of geometries used in the averaging procedure is tested. Both equilibrium structure and the ΔE and f data are obtained at the same level of theory, DKH3–SOC–MS–CASPT2(12,10)/ANO–RCC–VQZP. The results of the simulations using $N_p=25, 50, 75,$ and 100 geometries in the ensemble are presented in Figure S.2. It can be easily noticed that the convergence of $\sigma(E)$ is reached fairly quickly, with basically no observable differences in the spectra, i.e. no shifts of bands are noticed and only minor changes in the bands curvature are observed. Moreover, they affect mostly the broad band a and the right-side peak of main band b (observed at 6.46 eV). Even with very small and insufficient number of sampled geometries $N_p=25$, the qualitative picture of the spectrum is still conserved (although this would be true only for such small triatomic system). It is worth mentioning that the statistical errors of the sampling $\delta\sigma(E)$ (not shown in Figure S2) decreases with the number

of the sampled geometries, which can be already predicted by the form of Equation 3.

In conclusion, the chosen number of sampled geometries $N_p=100$ is proved to be completely reliable.

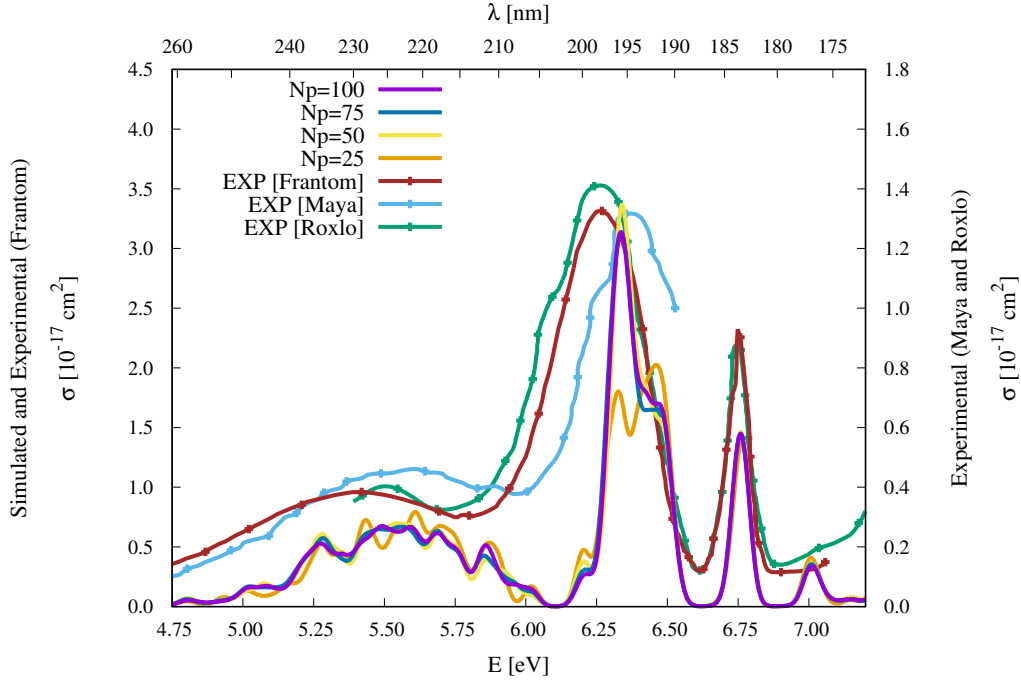


FIG. S.2: Convergence of the DKH3-SOC-MSCASPT2(12,10)/ANO-RCC-VQZP simulated spectra of HgBr_2 with different number of sampled geometries (N_p). Sampled geometries were obtained with the Wigner distribution using DKH3-CASPT2(12,10)/ANO-RCC-VQZP equilibrium structure. Experimental spectra²¹⁻²³ are included for the comparison (for better comparison, the $\sigma(E)$ values of spectrum of Frantom *et al.*²³ shown on the same axis as theoretical ones).

S.IV. SIMULATED UV–VIS ABSORPTION CROSS SECTIONS OF HGXY

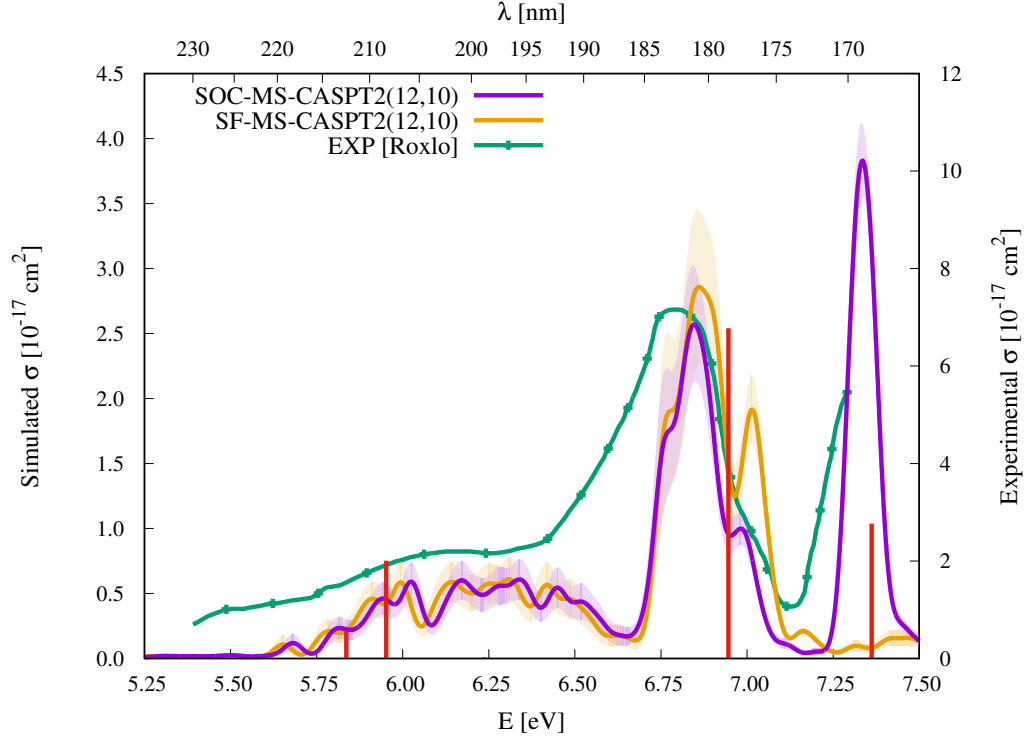


FIG. S.3: Simulated UV–Vis absorption cross sections ($\sigma(E)$) of HgCl_2 both without (SF, orange) and with (SOC, violet) taking into account the spin–orbit coupling between states. Shaded areas represent the statistical error for the sampled photon energies ($\delta\sigma(E)$). Experimental spectrum²¹ is included for the comparison. Additionally, vertical excitation energies in eV (ΔE) and associated f values between SO states for the optimized ground–state structure of HgCl_2 are graphically represented with red sticks (see the discussion in Section II).

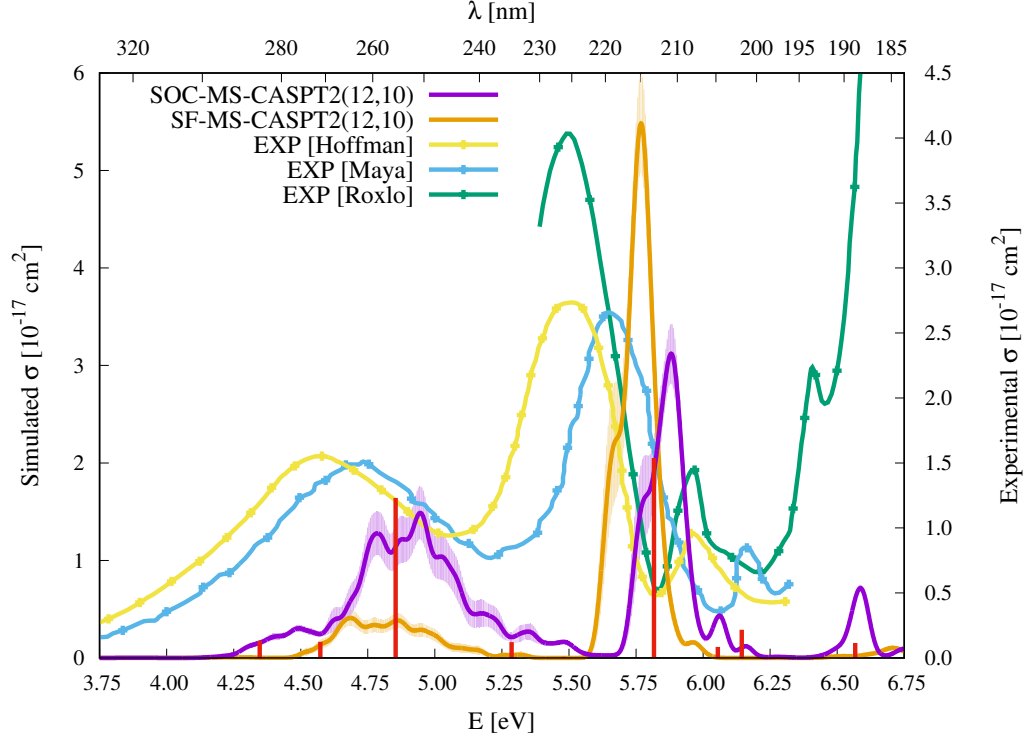


FIG. S.4: Simulated UV-Vis absorption cross sections ($\sigma(E)$) of HgI_2 both without (SF, orange) and with (SOC, violet) taking into account the spin-orbit coupling between states. Shaded areas represent the statistical error for the sampled photon energies ($\delta\sigma(E)$). Experimental spectra^{21,22,24} are included for the comparison. Additionally, vertical excitation energies in eV (ΔE) and associated f values between SO states for the optimized ground-state structure of HgI_2 are graphically represented with red sticks (see the discussion in Section II).

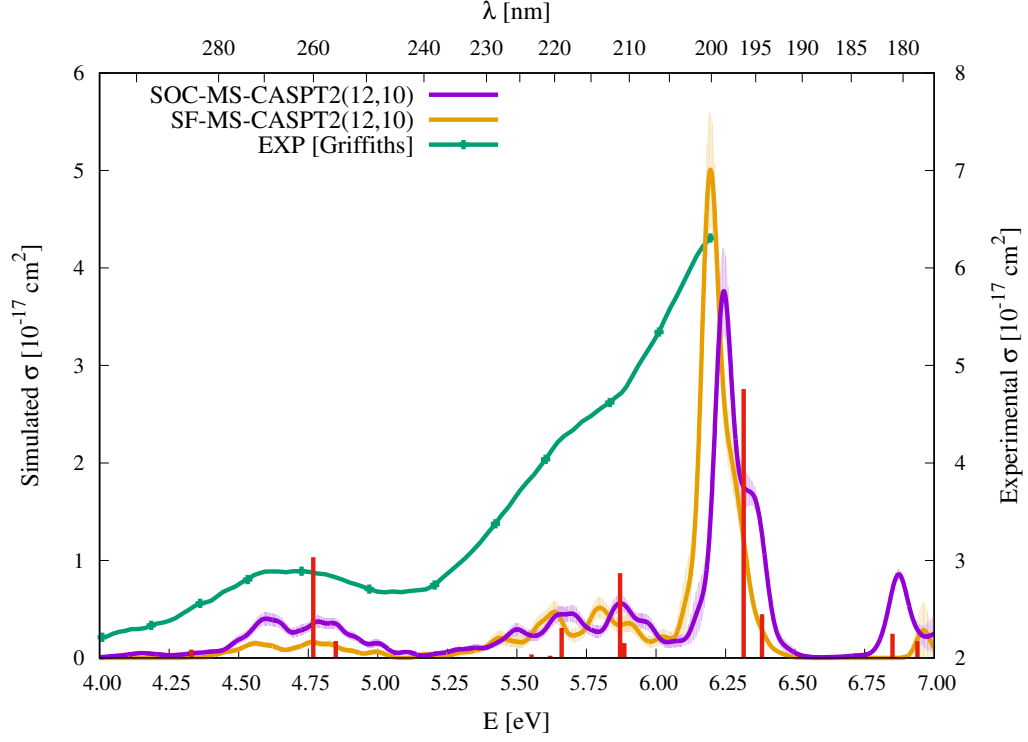


FIG. S.5: Simulated UV-Vis absorption cross sections ($\sigma(E)$) of HgClI both without (SF, orange) and with (SOC, violet) taking into account the spin-orbit coupling between states. Shaded areas represent the statistical error for the sampled photon energies ($\delta\sigma(E)$). Experimental spectrum²⁵ is included for the comparison. Additionally, vertical excitation energies in eV (ΔE) and associated f values between SO states for the optimized ground-state structure of HgClI are graphically represented with red sticks (see the discussion in Section II).

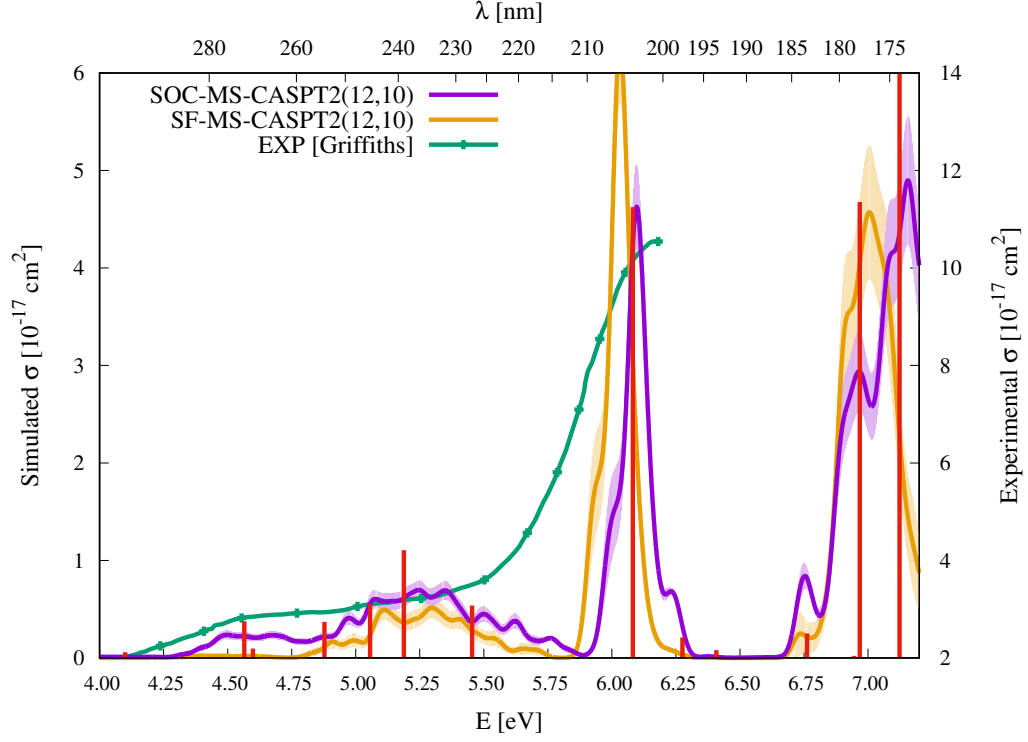


FIG. S.6: Simulated UV-Vis absorption cross sections ($\sigma(E)$) of HgBrI both without (SF, orange) and with (SOC, violet) taking into account the spin-orbit coupling between states. Shaded areas represent the statistical error for the sampled photon energies ($\delta\sigma(E)$). Experimental spectrum²⁵ is included for the comparison. Additionally, vertical excitation energies in eV (ΔE) and associated f values between SO states for the optimized ground-state structure of HgBrI are graphically represented with red sticks (see the discussion in Section II).

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