

Electronic Supplementary Information:

Vibrationally Induced Metallization of Energetic Azide $\alpha\text{-NaN}_3$

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S1 Vibrational frequencies for $\alpha\text{-NaN}_3$

The zone-centre vibrational modes were calculated for crystalline $\alpha\text{-NaN}_3$ using both the PBE-D2 and PBE-TS schemes, with resulting frequencies shown in Table S1.1.

Table S1.1: Zone centre vibrational frequencies for $\alpha\text{-NaN}_3$ (primitive cell) based on the PBE DFT functional. Values are shown corresponding to the Grimme D2 (PBE-D2) and Tkatchenko-Scheffler (PBE-TS) dispersion corrections. LO-TO splitting is not considered.

	PBE-D2 /cm⁻¹	PBE-TS /cm⁻¹
M_1	150.974	126.970
M_2	177.998	174.320
M_3	202.210	187.008
M_4	214.510	189.240
M_5	220.220	226.954
M_6	606.277	611.254
M_7	609.716	612.105
M_8	1250.429	1252.237
M_9	1959.812	1961.848

S2 Inelastic Neutron Scattering Spectra for $\alpha\text{-NaN}_3$

At low temperature, the intensity, I , of inelastic neutron scattering increases as¹

$$I \propto (Q.u)^2$$

where Q is the momentum transfer and u is the total atomic displacement of the normal coordinate. On the TOSCA instrument, the back-scattering detectors follow a trajectory in $S(Q,\omega)$ slightly higher in Q than the forward scattering detectors.² Hence the intensity of the back-scattering detectors should exhibit larger scattering intensities. This is most notable at low wavenumber. However, for $\alpha\text{-NaN}_3$, the low wavenumber scattering intensity is higher in the forward scattering detectors, Figure S1. This likely stems

from the dominant coherent scattering (directionally dependent) of NaN_3 , and requires further investigation.

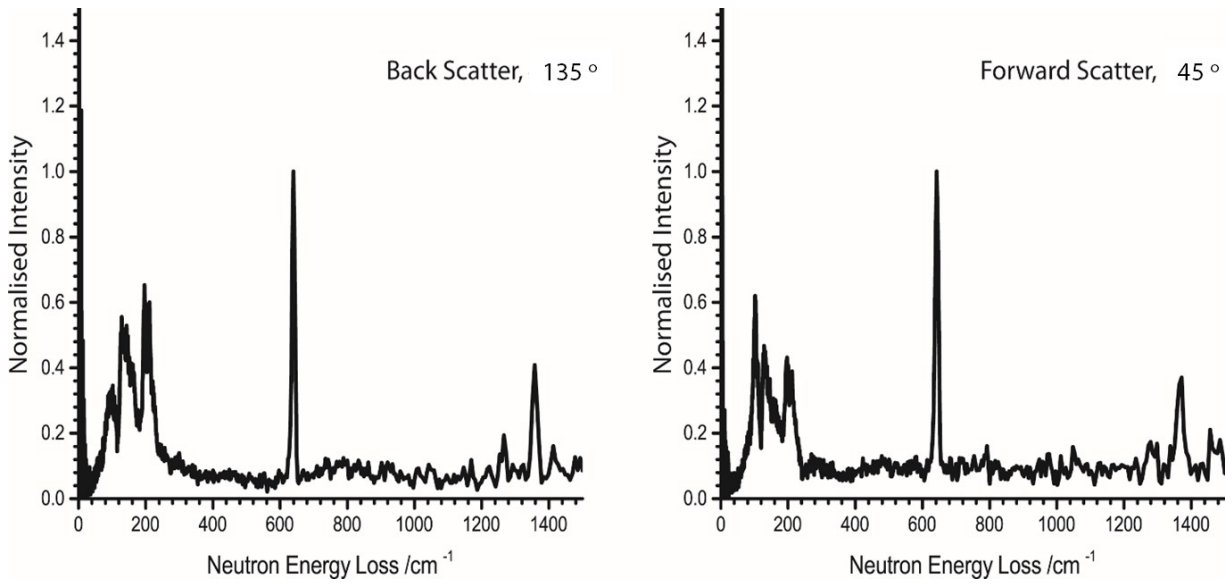


Figure S1: Inelastic neutron scattering spectra for $\alpha\text{-NaN}_3$ at *ca* 20 K. Signal from the forward and back-scattering detectors are shown separately. Intensities have been normalised to the dominant peak at *ca* 630 cm^{-1} .

S3 Eigenvector ‘walking’ for M_{16} in $\alpha\text{-NaN}_3$

The eigenvector associated with M_{16} exhibits considerable decrease in the band gap energy in $\alpha\text{-NaN}_3$. Further extension of this eigenvector results in near metallisation, Figure S2.

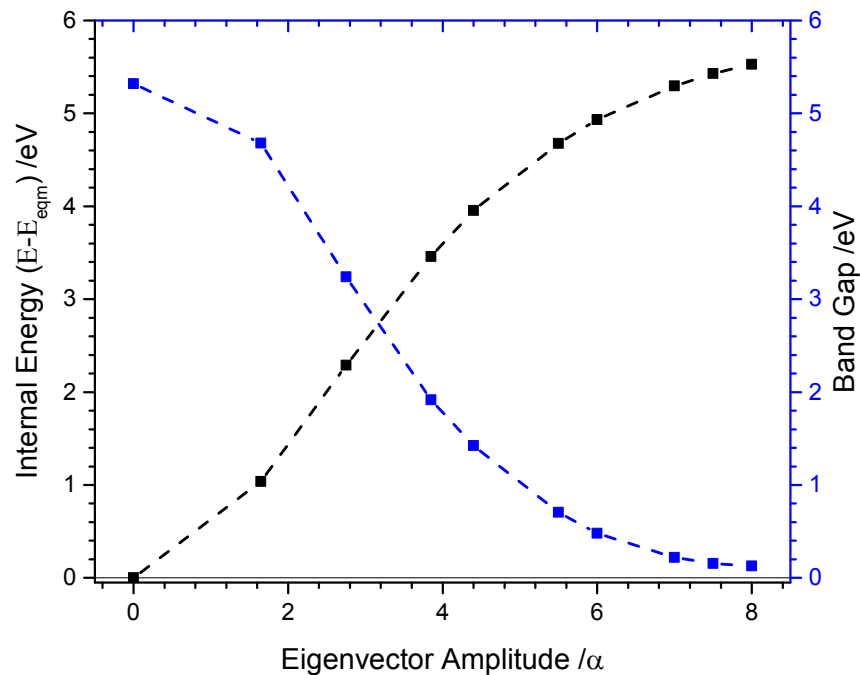


Figure S2: Band gap and internal energy along Mode 16, extended to $\alpha = 8$.

S4 Eigenvectors associated with the external modes of α -NaN₃

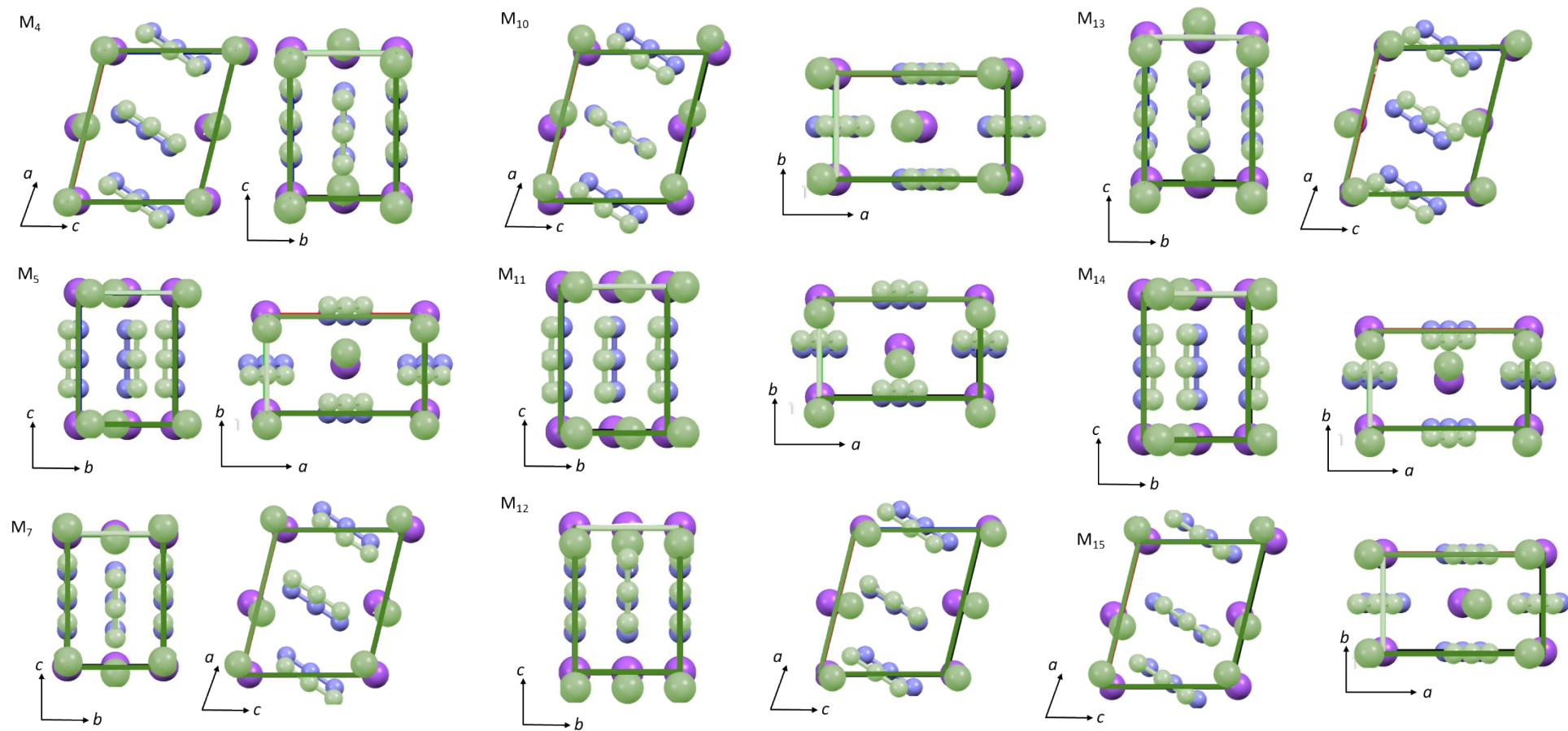


Figure S3: Eigenvectors of the external vibrational modes of α -NaN₃. Displacements are shown along the two crystallographic axes that best illustrate the distortion.

S5 Band Gap Reduction for M_{18} , M_{20} , M_{21} and M_{23}

The band gap was followed as a function of the eight internal modes of the conventional unit cell, Figure 6 and 8 in the main text, and Figure S4. Note that M_{21} is the out-of-phase symmetric stretching mode. As such, this mode could not be followed to metallisation due to compression of one of the azido anion molecules.

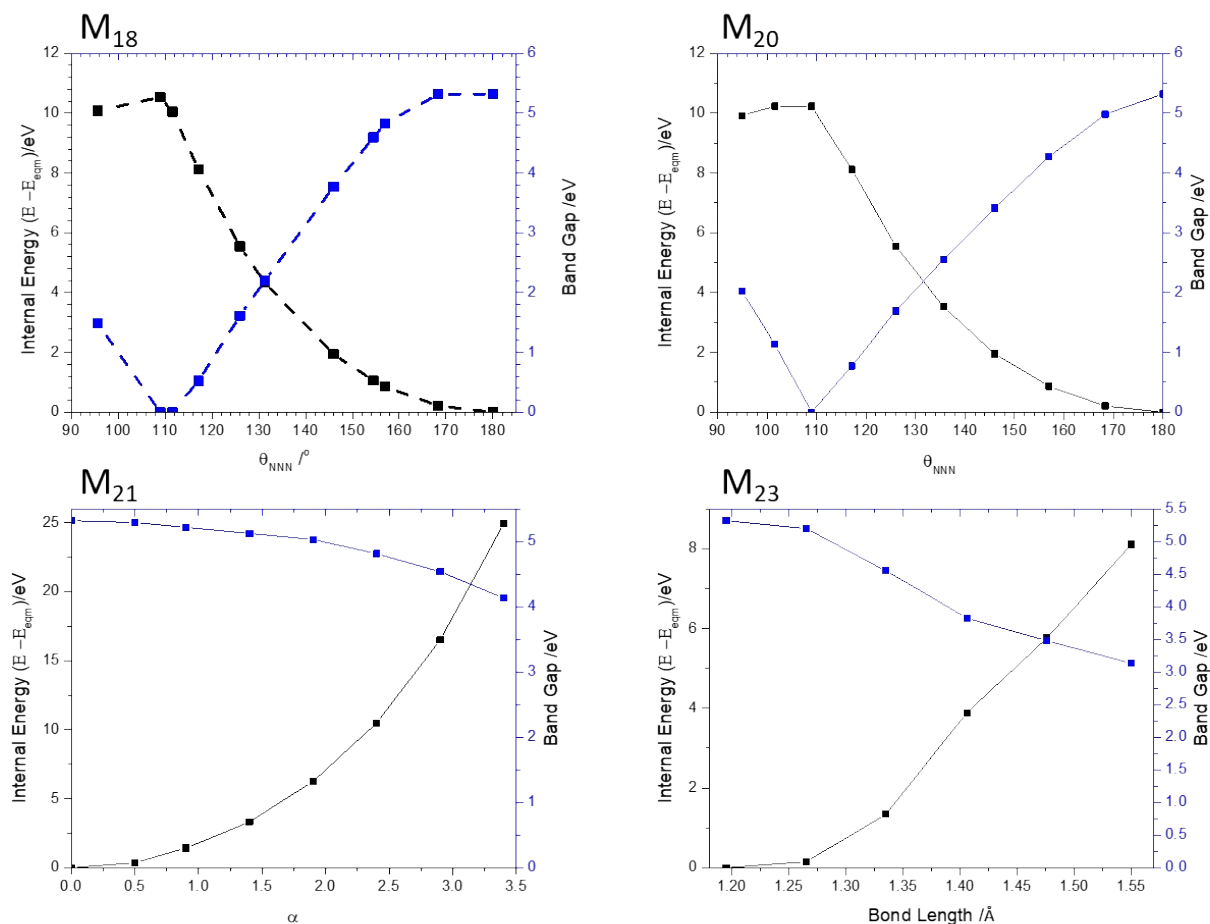


Figure S4: Band gap and internal energy as a function of the internal molecular modes of α -NaN₃.

S6 References

- 1 P. C. Mitchell, S. F. Parker, A. J. Ramirez-Cuesta and J. Tomkinson, *Vibrational Spectroscopy with Neutrons*, World Scientific, 2005.
- 2 D. Colognesi, M. Celli, F. Cilloco, R. J. Newport, S. F. Parker, V. Rossi-Albertini, F. Sacchetti, J. Tomkinson and M. Zoppi, *Appl. Phys. A Mater. Sci. Process.*, 2002, **74**, 64–66.