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Electronic Supplementary Information for:

Tortuosity and mass transfer limitations in industrial hydrotreating catalysts: effect of particles shape and size distribution

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Mean Molar Flux Radius evaluation (intraparticle mass transfer limitation)

Let's consider an ideal (perfectly mixed) stirred-tank reactor loaded with a liquid phase and porous particles (external mass transfer resistance is neglected). The molar balance on the reactor gives:

$$V_{R} \cdot \varepsilon_{i} \cdot \frac{dC_{f}}{dt} = -\sum_{i} n_{T} \cdot x_{i} \cdot V_{i} \varepsilon_{p} \cdot \frac{d\overline{C}_{i,p}}{dt}$$
(SI.1.)

Where V_R is the reactor volume, ε_i the extragranular porosity, C_f the concentration in the fluid phase, n_T the total number of particles, x_i the number fraction of particles of volume V_i , ε_p the intragranular porosity and $C_{i,p}$ the average concentration in the pore volume.

Considering spherical particles:

$$\frac{d\overline{C}_{i,p}(t)}{dt} = \frac{3}{R_{pi}^{3}} \int_{0}^{R_{pi}} r^{2} \cdot \frac{\partial C_{i,p}(r,t)}{\partial t} dr$$
(SI.2.)

For a Fickian diffusion, the mass balance in the particles provides the following expression:

$$\frac{\partial C_{i,p}(r,t)}{\partial t} = \frac{1}{r^2} \cdot \frac{\partial}{\partial r} (r^2 J_i(r,t))$$
(SI.3.)

with:
$$J_i(r,t) = -D_{eff} \cdot \frac{\partial C_{i,p}(r,t)}{\partial r}$$
 (SI.4.)

Replacing equations (SI.2.) to (SI.4.) in (SI.1.) gives

$$V_{R} \cdot \varepsilon_{i} \cdot \frac{dC_{f}}{dt} = -\sum_{i=1}^{m} n_{T} \cdot x_{i} \cdot V_{i} \varepsilon_{p} \cdot \left(\frac{3}{R_{pi}^{3}} \cdot \int_{0}^{R_{i}} \frac{\partial}{\partial r} (r^{2} J_{i}(r, t)) \cdot dr \right)$$
(SI.5.)

For a spherical particle:

$$V_i^{3}/R_{pi}^{3} = 4.\pi$$
 (SI.6.)

By inserting (SI.6.) in (SI.5.) and carrying out the change of variable $ho=r/R_{pi}$, we get:

$$V_{R} \cdot \varepsilon_{i} \cdot \frac{dC_{f}}{dt} = -\sum_{i} n_{T} \cdot x_{i} \cdot \varepsilon_{p} \cdot 4 \cdot \pi \cdot R_{pi} \cdot \int_{0}^{1} \frac{\partial}{\partial \rho} (\rho^{2} J_{i}(\rho, t)) \cdot d\rho$$
(SI.7.)

The total number of particles is:

$$n_T = \frac{V_S}{V_n} \tag{SI.8.}$$

Where Vs is the total solid volume and $^{\nabla_n}$ is the average particle volume, defined by :

$$\vec{V}_n = \sum_i x_i \cdot V_i \tag{SI.9.}$$

$$V_S = (1 - \varepsilon_i).V_R \tag{SI.10.}$$

Inserting equations (SI.8.) to (SI.10.) into (SI.7.) and considering that the mean molar flux density $\mathcal{I} = \int\limits_0^1 \frac{\partial}{\partial \rho} \Big(\rho^2 J_i(\rho,t) \Big) . d\rho$ is the same for all particles finally gives:

$$\frac{dC_f}{dt} = -\frac{(1 - \varepsilon_i)}{\varepsilon_i} \cdot \varepsilon_p \cdot 3 \cdot \frac{\sum_{i=1}^m x_i \cdot R_{pi}}{\sum_{i=1}^m x_i \cdot R_{pi}^3} \cdot 7$$
(SI.11.)

For spherical particles that are homogeneous in size, the molar balance is:

$$\frac{dC_f}{dt} = -\frac{(1 - \varepsilon_i)}{\varepsilon_i} \cdot \varepsilon_p \cdot \frac{3}{R_p^2}$$
(SI.12.)

Comparing equations (SI.11.) and (SI.12.) enables to define the mean radius as:

$$R_{mmf} = R_{p} = \sqrt{\frac{\sum_{i=1}^{m} x_{i} \cdot R_{pi}^{3}}{\sum_{i=1}^{m} x_{i} \cdot R_{pi}}}$$