# Supplementary Information for: Effects of symmetry breaking on the translation-rotation eigenstates of $\mathrm{H}_{2}, \mathrm{HF}$, and $\mathrm{H}_{2} \mathrm{O}$ inside the fullerene $\mathrm{C}_{60}{ }^{\dagger}$ 

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## 1 Hamiltonian Parameters

## $1.1 \quad \mathrm{H}_{2} @ \mathrm{C}_{60}$

The kinetic energy operator for $\mathrm{H}_{2} @ \mathrm{C}_{60}$ was taken to be

$$
\begin{equation*}
\hat{T}=-\frac{\nabla^{2}}{2 M}+B \hat{j}^{2} \tag{1}
\end{equation*}
$$

where $\nabla^{2}$ is the Laplacian associated with $\mathbf{R}, \hat{j}^{2}$ is the operator corresponding to the square of the rotational angular momentum of the $\mathrm{H}_{2}, M$ is the mass of the $\mathrm{H}_{2}$, and $B$ is the
rotational constant of the $\mathrm{H}_{2}$. We used $M=2.0104 \mathrm{amu}, B=58.378 \mathrm{~cm}^{-1}$ for the $v=0$ manifold, and $B=54.83 \mathrm{~cm}^{-1}$ for the $v=1$ manifold.

The $V_{M-C_{60}}$ function for both the $v=0$ and the $v=1$ manifolds was taken to be a pairwise-additive Lennard-Jones one of the form

$$
\begin{equation*}
V_{H_{2}-C_{60}}=\sum_{i=1}^{3} \sum_{k=1}^{60} 4 w_{i} \epsilon\left[\left(\frac{\sigma}{r_{i k}}\right)^{12}-\left(\frac{\sigma}{r_{i k}}\right)^{6}\right] \tag{2}
\end{equation*}
$$

where $i$ runs over three $\mathrm{H}_{2}$ sites, $k$ runs over the 60 nuclear positions of the C atoms in the central cage, and $r_{i k}$ is the distance between site $i$ and site $k$. For both manifolds the $\mathrm{H}_{2}$ site 1 was located at the center of the HH bond, and sites 2 and 3 were located at the H nuclei. For $v=0$ the HH bond distance was taken to be $0.74 \AA, w_{1}=6.7, w_{2}=w_{3}=1$, $\sigma=2.95 \AA$, and $\epsilon=3.07 \mathrm{~cm}^{-1} .{ }^{1}$ For $v=1$ the HH bond distance was taken to be 0.78132 $\AA, w_{1}=7.5, w_{2}=w_{3}=1, \sigma=2.95 \AA$, and $\epsilon=2.9886668 \mathrm{~cm}^{-1} .{ }^{2}$ The $\mathrm{C}_{60}$ geometry was taken to be that used in Felker, et al. ${ }^{3}$

As to $V_{\text {quad }}$, the $\mathrm{BF} \hat{z}$ axis was taken to be the internuclear axis, and the one nonzero BF quadrupole component for $\mathrm{H}_{2}$ was taken to be $Q_{0}^{\mathrm{BF}}=0.499$ au for both the $v=0$ and $v=1$ manifolds. This is the same value that was used in Felker, et al. ${ }^{3}$

## 1.2 $\mathrm{HF}^{( } \mathrm{C}_{60}$

The kinetic energy operator for $\mathrm{HF}_{6} \mathrm{C}_{60}$ was taken to have the same form as eqn (1) but with $M=20.006225 \mathrm{amu}$ and $B=18.523 \mathrm{~cm}^{-1}$. This value for $B$ is the cage-modified one determined by Kalugina and Roy. ${ }^{4}$

The $V_{H F-C_{60}}$ function was taken directly from Kalugina and Roy. ${ }^{4}$ It is an expansion over bipolar spherical tensors dependent on the four angles $(\Theta, \Phi, \omega)$ with $R$-dependent expansion coefficients. It does not require any input as to the $\mathrm{C}_{60}$ geometry.

The $\mathrm{BF} \hat{z}$ axis was taken by Kalugina and Roy $^{5}$ to be the internuclear axis pointing from the H nucleus to the F nucleus. As such $\vec{\mu}=\mu \hat{z}$ is antiparallel to $\hat{z}$ and $\mu$ is negative. We
take the magnitude of $\mu$ to be the screened value of -0.177 au from Krachmalnicoff, et al. ${ }^{6}$

## $1.3 \quad \mathrm{H}_{2} \mathrm{O} @ \mathrm{C}_{60}$

The kinetic energy operator for $\mathrm{H}_{2} \mathrm{O} @ \mathrm{C}_{60}$ was taken to be

$$
\begin{equation*}
\hat{T}=-\frac{\nabla^{2}}{2 M}+B_{x} \hat{j}_{x}^{2}+B_{y} \hat{j}_{y}^{2}+B_{z} \hat{j}_{z}^{2} \tag{3}
\end{equation*}
$$

where $\nabla^{2}$ is the Laplacian associated with $\mathbf{R}, \hat{j}_{x}, \hat{j}_{y}$, and $\hat{j}_{z}$ are the operators associated with the components of the rotational angular momentum of the $\mathrm{H}_{2} \mathrm{O}$ along the BF axes, which are take to be its principal inertial axes. We used $M=18.0105 \mathrm{amu}, B_{x}=27.877 \mathrm{~cm}^{-1}$, $B_{y}=9.285 \mathrm{~cm}^{-1}$, and $B_{z}=14.512 \mathrm{~cm}^{-1}$. This choice of the BF axes locates the bisector of the HOH bond angle to be along the $\mathrm{BF} \hat{z}$ axis.

The $V_{M-C_{60}}$ for $\mathrm{H}_{2} \mathrm{O} @ \mathrm{C}_{60}$ was taken from Felker and $\mathrm{Baccićc}^{7}$ and is given by

$$
\begin{equation*}
V_{\mathrm{H}_{2} \mathrm{O}-C_{60}}=\sum_{i=1}^{3} \sum_{k=1}^{60} 4 \epsilon_{i}\left[\left(\frac{\sigma_{i}}{r_{i k}}\right)^{12}-\left(\frac{\sigma_{i}}{r_{i k}}\right)^{6}\right], \tag{4}
\end{equation*}
$$

where $i$ runs over three $\mathrm{H}_{2} \mathrm{O}$ sites, $k$ runs over the 60 nuclear positions of the C atoms in the central cage, $r_{i k}$ is the distance between site $i$ and site $k, \sigma_{1}=3.372 \AA, \sigma_{2}=\sigma_{3}=2.640$ $\AA, \epsilon_{1}=36.34 \mathrm{~cm}^{-1}$, and $\epsilon_{2}=\epsilon_{3}=8.95384 \mathrm{~cm}^{-1}$. The three $\mathrm{H}_{2} \mathrm{O}$ sites are given in Table 2 of the ESI of Felker, et al. ${ }^{3}$ The $\mathrm{C}_{60}$ geometry was taken to be the same as that used for $\mathrm{H}_{2} @ \mathrm{C}_{60}$.

As to $V_{\text {quad }}$, since we take the $\mathrm{BF} \hat{z}$ axis to point from the c.m. of the water moiety toward the O nucleus along the HOH bond-angle bisector, then $\vec{\mu}=\mu \hat{z}$ is antiparallel to $\hat{z}$, and $\mu$ is negative. We used the screened dipole value, $\mu=-0.200 \mathrm{au}$, from Goh, et al. ${ }^{8}$ The BF quadrupole components of the $\mathrm{H}_{2} \mathrm{O}$ were taken to be the same as in Felker, et al.: ${ }^{3}$ $Q_{0}^{(\mathrm{BF})}=-0.09973 \mathrm{au}$ and $Q_{ \pm 2}^{(\mathrm{BF})}=1.53843 \mathrm{au}$.

## 2 Grid Parameters

As mentioned in Subsection 2.2 of the main body of the paper the TR state function, $|\psi\rangle$, employed in the Chebyshev filter diagonalization procedure was transformed to a grid representation to effect its multiplication by the potential-energy portion of $\hat{H}$. The general nature of the five-dimensional (5D) grid points used for $\mathrm{H}_{2} @ \mathrm{C}_{60}$ and for $\mathrm{HF} @ \mathrm{C}_{60}$, and the six-dimensional (6D) grid points used for $\mathrm{H}_{2} \mathrm{O} @ \mathrm{C}_{60}$ are described in Section 2.5 of Felker, et al. ${ }^{3}$ Further specifics as to the grids used in this work follow.

For $\mathrm{M}=\mathrm{H}_{2}$ we used (i) 12 Gauss-associated-Laguerre quadrature points generated as per Felker and Bačić ${ }^{9}$ with $\beta=2.9888989$ au for the $R$ coordinate, (ii) 10 Gauss-Legendre quadrature points for the $\cos \Theta$ coordinate, (iii) 18 Fourier grid points for the $\Phi$ coordinate, (iv) 10 Gauss-Legendre quadrature points for the $\cos \theta$ coordinate, and (v) 18 Fourier grid points for the $\phi$ coordinate. Here, the relevant Euler angles are $\omega=(\theta, \phi)$, where $\theta$ is the polar angle, and $\phi$ the azimuthal angle describing the orientation of the $\mathrm{BF} \hat{z}$ axis with respect to the SF axis system.

For $\mathrm{M}=\mathrm{HF}$ we used (i) 14 Gauss-associated-Laguerre quadrature points generated with $\beta=12.0$ au for the $R$ coordinate, (ii) 12 Gauss-Legendre quadrature points for the $\cos \Theta$ coordinate, (iii) 24 Fourier grid points for the $\Phi$ coordinate, (iv) 10 Gauss-Legendre quadrature points for the $\cos \theta$ coordinate, and (v) 18 Fourier grid points for the $\phi$ coordinate.

For $\mathrm{M}=\mathrm{H}_{2} \mathrm{O}$ we used (i) 12 Gauss-associated-Laguerre quadrature points generated with $\beta=24.38$ au for the $R$ coordinate, (ii) 10 Gauss-Legendre quadrature points for the $\cos \Theta$ coordinate, (iii) 18 Fourier grid points for the $\Phi$ coordinate, (iv) 10 Gauss-Legendre quadrature points for the $\cos \theta$ coordinate, (v) 18 Fourier grid points for the $\phi$ coordinate, and (vi) 18 Fourier grid points for the $\chi$ coordinate. Here, $\omega=(\phi, \theta, \chi)$ are the Euler angles, defined with the convention used in Zare, ${ }^{10}$ that specify the orientation of the BF axes of the $\mathrm{H}_{2} \mathrm{O}$ with respect to the SF axes.

## References

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