Supporting information

for

Non-contact monitoring of extra-cellular field potentials with a multi-electrode array

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A. Representative extra-cellular field potential waveforms from non-contact measurements

Figure S1. (a-f) Representative extra-cellular field potential waveforms recorded from a multielectrode array (MEA) in the non-contact configuration from N = 10 primary cardiomyocyte culture preparations grown on glass coverslips. Spike triggered averages appear as the red solid line (averaged over the number of detected spikes, n_s , above background noise levels). Individual spike waveforms are plotted in gray. A total of 135 electrodes (from N = 10 preparations) registered spikes distinguishable above background noise: 30 percent of these spikes displayed only negative peaks, 44% displayed an initial positive component followed by negative peak and 26% registered only a positive component to the waveform. **B.** Representative extra-cellular field potenial waveforms from cells grown on the surface of a multi-electrode array



Figure S2. (a-d) Representative extra-cellular field potential waveforms produced by primary cardiomyocyte cells grown on the surface of a multi-electrode array (MEA).



C. Characterization of coverslip surface approach

Figure S3. (a) Separate measurements of the coverslip surface height h_s as a function of magnetmagnet separation distance d_m . (b) The relative angular variation between the surface ϕ_s and the coverslip ϕ_c is plotted for fixed coverslip surface heights h_s during approach. The mean angular variation is $0.17^\circ \pm 0.06^\circ$ measured over a distance of 2 mm. The variation in h_s due to angular variability during approach is $\approx 3 \ \mu m$ across the 1.1 mm distance of the recording surface of the multi-electrode array. (c-e) Optical images of the relative coverslip surface height h_s during approach.



D. Microscopy of primary cardiomyocyte cultures

Figure S4. (a) Immunofluorescence images taken of primary cardiomyocyte cultures at 10 *days in vitro* taken with a Leica SP8 resonant scanning confocal microscope using a 63x oil immersion objective. The cell membrane shown in green (WGA-AF488) and nuclei shown in blue (DAPI). The scale bar is 100 μ m. (b) Z-stack image taken along the dotted line of (a). The vertical and horizontal scale bars are 2 μ m and 100 μ m, respectively. (c) Phase contrast images of primary cardiomyocytes cultures taken at 5 days *in vitro*. The red arrows indicate spontaneously contracting cells. The scale bar is 100 μ m. (d) Phase contrast images of primary cultures taken at 14 days *in vitro*. The red arrows point to spontaneously contracting, multi-cellular populations and highlighted by the dashed line. The scale bar is 100 μ m.

E. Modeling extra-cellular field potentials

The extra-cellular voltage produced by an action potential has been previously shown by Gold *et al.*¹ to be well modeled as an isotropic volume conductor with a purely Ohmic response over a frequency range of interest between 1-3,000 Hz (capacitive effects of the extra-cellular medium are negligible). Under steady-state conditions, the electric potential V_s generated by a point current source I_s flowing through an isotropic volume conductor with electrical conductivity σ is given by

$$V_s = \frac{I_s}{4\pi\sigma \cdot r},\tag{1}$$

where r the radial distance from the point source.²

Cardiomyocyte cultures grow in confluent clusters with ion channels (current sources) spread out across their surfaces. We next approximate the cell clusters as a homogeneous distribution of current sources with current density j_0 distributed across the surface area of a planar 2d disk of radius R_0 as shown by Diagram 1. The contribution of source current, dI_s , at disc radius *a* is given by $dI_s = j_0 \cdot 2\pi a \cdot da$. The resultant field at point *P* is the linear superposition of point sources³ and is given by



Diagram 1. 2D model of a planar current source in an isotropic volume conductor.

$$V^{theory} = 2 \cdot \int_0^{R_0} \frac{dI_s}{4\pi\sigma \cdot r} = \int_0^{R_0} \frac{j_0 \cdot 2\pi a \cdot da}{4\pi\sigma \sqrt{h_s^2 + a^2}} = \frac{j_0}{\sigma} \left(\sqrt{h_s^2 + R_0^2} - |h_s| \right) + V_0, \tag{2}$$

where h_s is the height from the center of the disk. The extra-cellular voltage is multiplied by a factor of two because current flow subtends half the solid angle for a planar 2D cell culture. The constant V_0 term is added to account for a voltage drop across the recording electrode interface due to finite impedance, as well as any electrochemical voltage offsets due to grounding the liquid. The cell culture media conductivity σ has a value of 1.7 S·m⁻¹ at 37 °C.⁴ To account for the uncertainty in V_0 , we take the derivative of Eq.2 with respect to h_s and arrive at the following expression for the current density

$$j_{0} = \sigma \left(\frac{dV}{dh_{s}}\right) \cdot \left(\frac{h_{s}}{\sqrt{h_{s}^{2} + R_{0}^{2}}} - 1\right)^{-1}.$$
(3)

Note that j_0 depends only the slope (dV/dh_s), determined by the fit to experimental data, and distance h_s for a given culture size set by R_0 . This value is independent of the magnitude of the voltage signal, which can vary significantly based on the quality of the recording electrode.⁵

Figure S5a,b shows the experimentally measured extra-cellular field potential vs. time for various surface heights h_s from the recording surface of a MEA. Figure S5c shows V^{theory} as a function of height h_s from the cell-culture surface, for $R_0 = 150 \,\mu\text{m}$ and $j_0 = 1.7 \,\text{A/m}^2$. The red circles and black diamonds, shown in Figure S5d, indicate experimentally measured maximum voltage amplitudes (extracted from Figure S5a,b) measured at 6 DIV (V_6^{expt}) and 18 DIV (V_{18}^{expt}). The dotted lines are theoretical fits V_6^{theory} and V_{18}^{theory} to the experimental data using Eq.2. A two-parameter fit was used to determine j_0 and V_0 for a fixed excitable culture area set by R_0 . The excitable area was determined by spatial distribution of synchronized spiking activity measured by the MEA. Cardiomyocyte activation travels as a wave of electrical activity which propagates from cell-to-cell with an average velocity of $\approx 0.3 \,\text{m}\cdot\text{s}^{-1}$.^{6–8} Voltage signals detected at a given electrode on the MEA are the superposition of all synchronized extra-cellular field

potentials. For a fixed reference point, extra-cellular field potential signals are additive over a maximum spatial window of $\approx 300 \ \mu\text{m}$ assuming a depolarization time of $\approx 0.5 \ \text{ms}$, therefore we used an $R_0 = 150 \ \mu\text{m}$. We find $j_0 = 1.7 \ \text{A} \cdot \text{m}^{-2}$ and 2.4 $\ \text{A} \cdot \text{m}^{-2}$ for the 6 and 18 DIV respectively. These results are consistent with whole-cell patch-clamp measurements that give current densities of around $\approx 1.3 \ \text{A} \cdot \text{m}^{-2}$ (The magnitude of whole cell transmembrane currents span the range $10^{-9} \ \text{A}$ to $10^{-8} \ \text{A}$, $^{9-12}$ which flow across an area of $\approx 4 \cdot 10^{-9} \ \text{m}^{2}$.^{6,12,13}).



Figure S5. Distance dependence of extra-cellular field potentials. (a-b) Experimental non-contact measurements of extra-cellular field potentials from cardiomyocyte cultures at 6 *days in vitro* (DIV) and 18 DIV for various separation heights from the surface of a multi-electrode array. The voltage vs. time plots are the spike triggered average waveforms. (c) Theoretically modeled extra-cellular voltage as a function a distance h_s from the cell surface. (d) Experimentally measured extra-cellular peak voltage amplitude as a function of distance for two representative cardiomyocyte cultures at 6 DIV (V_6^{expt} , black diamonds) and 18 DIV (V_{18}^{expt} , red circles). The dashed lines are theoretical fits to the data.

The parameters of our model-dependent results assume homogenous spatial distribution of the cells, which is a reasonable approximation for the contiguous arrangements of cardiomyocytes whose temporal activation is synchronized via gap-junctions.⁶ Incorporating cell patterning techniques would help to define more precise spatial geometries of varying degrees of complexity, ranging from sculpting the shape of single cells¹⁴ to templating the geometry of large populations that direct information flow through out interconnected networks.^{15–17} Lastly, combining these techniques with high-channel count MEAs^{18–20} would enable the development of more precise

computational models used to describe extra-cellular fields generated by inhomogeneous spatial configurations.

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