Supporting Information

Elimination of Concentration Dependent Luminescence Quenching in Surface Protected Upconversion Nanoparticles

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Materials

YCl₃·6H₂O (99.99%), YbCl₃·6H₂O (99.998%), TmCl₃·6H₂O (99.99%), GdCl₃·6H₂O (99.99%), LuCl₃·6H₂O (99.99%), NH₄F (>98%), NaOH (>97%), oleic acid (OA, 90%), and 1-octadecene (ODE, 90%) were purchased from Sigma-Aldrich and used as received without further purification.

Analytical interpretation of the concentration quenching for core-shell UCNPs

Although the surface of the UCNPs is well shielded from outside, sensitizer can still cause concentration quenching due to the deleterious effects of Yb^{3+} . Here we used the mathematical modelling to simulate the processes (**Fig. S1**).

 $W_2 \sim W_5$ represents the radiative relaxation rate for different energy levels. K_{51} , K_{41} , K_{31} , denotes the cross-relaxation between different energy level. $n_1 \sim n_5$ are the energy level population for the UC system. $C_1 \sim C_4$ are energy transfer rate between Yb³⁺ and Tm³⁺ for stepwise UC processes. P_{980} is the pumping rate of Yb³⁺ ions under 980 nm excitation laser with a fixed power. τ_s denotes the intrinsic decay rate for Yb³⁺ions at excited state.

$$\frac{dn_{S2}}{dt} = P_{980}n_{S1} - \frac{n_{S2}}{\tau_{S2}} - Q - (c_1n_1 + c_2n_2 + c_3n_3 + c_4n_4)n_{S2}$$
(S1)

$$\frac{dn_2}{dt} = c_1 n_{S2} n_1 - c_2 n_{S2} n_2 - W_2 n_2 + b_{52} W_5 n_5 + 2k_{31} n_1 n_3 + k_{41} n_1 n_4$$
(S2)

$$\frac{dn_3}{dt} = c_2 n_{S2} n_2 - c_3 n_{S2} n_3 - W_3 n_3 - k_{31} n_1 n_3 + k_{41} n_1 n_4 + 2k_{51} n_1 n_5$$
(S3)

$$\frac{dn_4}{dt} = c_3 n_{S2} n_3 - c_4 n_{S2} n_4 - W_4 n_4 - k_{41} n_1 n_4 \tag{S4}$$

$$\frac{dn_5}{dt} = c_4 n_{S2} n_4 - W_5 n_5 - k_{51} n_1 n_5 \tag{S5}$$

When system reaches a steady state, all the above temporal differentiations equals to zero.

$$0 = P_{980}n_{S1} - \frac{n_{S2}}{\tau_{S2}} - Q - (c_1n_1 + c_2n_2 + c_3n_3 + c_4n_4)n_{S2}$$
(S6)

$$c_1 n_{s_2} n_1 = c_2 n_{s_2} n_2 + W_2 n_2 - b_{s_2} W_s n_5 - 2k_{31} n_1 n_3 - k_{41} n_1 n_4$$
(S7)

$$c_2 n_{S2} n_2 = c_3 n_{S2} n_3 + W_3 n_3 + k_{31} n_1 n_3 - k_{41} n_1 n_4 - 2k_{51} n_1 n_5$$
(S8)

$$c_3 n_{S2} n_3 = c_4 n_{S2} n_4 + W_4 n_4 + k_{41} n_1 n_4 \tag{S9}$$

$$c_4 n_{S2} n_4 = W_5 n_5 + k_{51} n_1 n_5 \tag{S10}$$

Substitute equation (S2) to (S5) repetitively into (S1), we obtain

$$0 = P_{980}n_{S1} - \frac{n_{S2}}{\tau_{S2}} - Q - \sum_{i=2}^{5} (i-1)W_i n_i$$
(S11)

The term $\sum_{i=2}^{5} (i-1)W_i n_i$ can be viewed as the integration of total emission spectrum. To make it clearer, we use E to substitute the term.

$$E = \sum_{i=2}^{5} (i-1)W_i n_i = P_{980} n_{S1} - \frac{n_{S2}}{\tau_{S2}} - Q$$
(S12)

In addition, $\frac{n_{S2}}{\tau_{S2}}$ denotes the Yb³⁺ that relaxed from the excited state to the ground state in Yb-Tm UC system. Since the energy level is well matched (²F_{2/5}-³H₅), the energy is mostly transferred to Tm³⁺ ions^[1]. Therefore, we can ignore $\frac{n_{S2}}{\tau_{S2}}$ since it is negligible compared to the other term, and the equation can be re-written as

$$E = \sum_{i=2}^{5} (i-1)W_i n_i = P_{980} n_{S1} - Q$$
(S13)

 $P_{980}n_{S1}$ together can be viewed as the total photons absorbed by Yb³⁺. Based on the photon absorption theory, $\frac{dN}{dx} = -\kappa N C_{Yb^{3+}} \sigma$ in which N is the total number of photons at depth x, κ is constant, $C_{Yb^{3+}}$ is the concentration of Yb³⁺, σ is the absorption cross-section. After integration, we can obtain $P_{980}n_{S1} = N_0(1 - e^{-\alpha C_{Yb^{3+}}})$, where N₀ is the total number of photons that reach on the UCNPs, α is a positive constant.

Q is the total energy distributed to the quencher. For non-radiative energy transfer, the transfer rate is proportional to the sixth power of distance^[2], namely, $k \propto \frac{1}{r^6}$. Assuming all elements ions in crystal are perfectly evenly distributed, the average distance (d) for the defects to the closest Yb³⁺ can be obtained as:

$$\frac{\frac{3}{2}\sqrt{3}a^2b}{C_{yb^{3+}}} = \frac{\pi d^3}{6}, \text{ and thus, } d = \frac{1}{\pi} \left(\frac{\frac{3}{2}\sqrt{3}a^2b}{C_{yb^{3+}}}\right)^{1/3}$$

In which, a, b is the lattice constant, thus the quenching factor Q can be related to $C_{\rm Yb3^+}$ as

$$Q \propto \left(C_{Yb^{3+}}\right)$$
, or $Q = \beta \left(C_{Yb^{3+}}\right)$

Together, the equation can be changed into

$$E = N_0 (1 - e^{-\alpha C_{Yb^{3+}}}) - \beta (C_{Yb^{3+}})$$
(S14)

Moreover, this model can be applied to other lanthanide-doped UC system, as long as the surface quenching is effectively prohibited.

Based on our model Eq. S14, we can get the first and second derivative of the E to concentration $C_{y_k^{3+}}$

$$E' = \alpha N_0 e^{-\alpha C_{\gamma b^{3+}}} - 2\beta C_{\gamma b^{3+}}$$
(S15)

$$E'' = -\alpha^2 N_0 e^{-\alpha C_{\gamma b^{3+}}} - 2\beta$$
 (S16)

Thus, we can obtain:

$$E\Big|_{C_{\gamma_{b}^{3+}=0}} = 0$$

$$E'\Big|_{C_{\gamma_{b}^{3+}=0}} > 0$$

$$E'' < 0$$

(S17)

Based on these three equations, we can conclude that the shape of the $E - C_{Yb^{3+}}$ curve is like a convex, with the maximum happened at the E' = 0. Thus, the increasing and decreasing trend of the luminescence intensity for core-shell samples can be well illustrated by using the photon-absorption and inner quenching center mechanism.

For core only UCNPs, the quenching term Q, should be revised as $Q \propto C_{\gamma b^{3+}}$, or $Q = \beta' C_{\gamma b^{3+}}$, since the energy mainly goes to the surface of the UCNPs. However, due to the difference of the power parameters between the core and core shell UCNPs, when E' = 0, different $C_{\gamma b^{3+}}$ are to be expected, thereby indicating a difference in the optimized composition for core and core-shell UCNPs.

The above analysis showed that, for Yb³⁺ sensitized co-doped core shell UCNPs, Yb³⁺ concentration cannot exceed the optimized point, otherwise, concentration quenching will happen inevitably. In the view of this, we hypothesized that moving the extra Yb³⁺ to another layer would lead to increased photon adsorption along with reducing the heavy quenching. In this case, we can introduce another term into the mathematics model:

$$E = N_0 (1 - e^{-\alpha C_{yb^{3+}}}) - \beta (C_{yb^{3+}}) + P$$
(S18)

where, *P* represents the photons injected from outside (sensitizing layer). Since varying the *P* does not affect the first two terms right side, we can simply enhance luminescence by merely increasing *P*, i.e., adding more Yb³⁺ ions to the sensitizing layer.



Figure S1. Simplified diagram of the Yb–Tm UC system and quenching mechanism.



Х	Core		Core-shell	
	Diameter(nm)	σ	Diameter(nm)	σ
20	22.72561	0.09412	29.98596	0.39393
30	23.21853	0.15663	30.06949	0.22974
45	22.83888	0.17044	29.93202	0.3439
60	23.15584	0.12458	30.17991	0.19885
70	22.77986	0.09179	30.27101	0.19726

Figure S2. Size histograms of a-e) NaYF₄: x%Yb, [(x-20)/2] %Gd, 1%Tm UCNPs with x=20, 30, 45, 60, 70 and f-j) NaYF₄: x%Yb, [(x-20)/2] %Gd, 1%Tm@NaLuF₄ UCNPs with x=20, 30, 45, 60, 70. Summary of size histograms analysis are in the table.



Figure S3. Temporal resolved decay curve showing the increasing and decay profile of NaYF₄: x%Yb, [(x-20)/2] %Gd, 1%Tm UCNPs: x=20, 30, 45, 60, 70 at peaks located at 804nm.



Figure S4. DLS measurements data showing the size increment between NaYF₄: 20%Yb, 1%Tm UCNPs and NaYF₄: 20%Yb, 1%Tm@NaLuF₄ UCNPs after NaLuF₄ shell coating.



Figure S5. EDS characterization (a) NaYF₄: 20Yb, 1Tm UCNPs and (b) NaYF₄: 20Yb, 1Tm@NaLuF₄ with peaks from lutetium highlighted in red color.



Figure S6. Temporal resolved decay curve showing the increasing and decay profile of NaYF₄: x%Yb, [(x-20)/2] %Gd, 1%Tm@NaLuF₄ UCNPs: x=20, 30, 45, 60, 70 at peaks located at 804nm.



Figure S7. Temporal resolved decay curve with duel-exponential fitted curve of ACD and ACD+ at peaks located at 804nm



Figure S8. a-c) TEM images of A, AC++ and ACD++. D) UC emission spectra of cyclohexane solutions at room temperature (25 ^oC) comprising ACD, ACD+ and ACD++.

References

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