

Supporting Information for

Unveiling Lasing Mechanism in CsPbBr₃ Microspheres Cavities

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I. The setup of chemical vapor deposition method

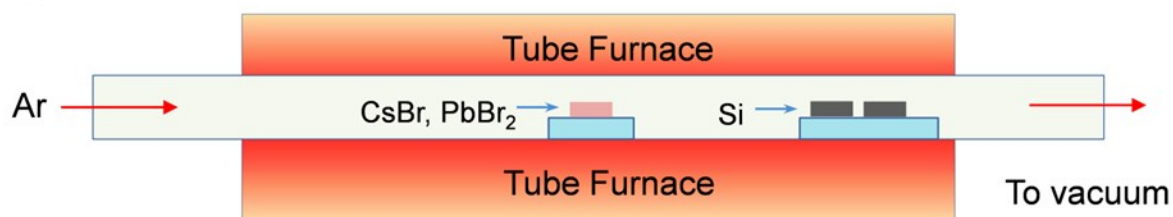


Figure S1 | Schematic of the chemical vapor deposition method setup

PbBr₂ and CsBr powders (with 1:2 stoichiometry) were placed in the heating center of a quartz tube reactor. The silicon (Si) substrate was positioned at a distance of 12 cm away from the powder sources in the downstream. Ar as the carrier gas is flowing from the upstream to the downstream. At a sample temperature of 575 °C, we observed the reliable growth of CsPbBr₃ perovskite micro-crystals with a special shape of microspheres.

II. SEM images of sphere samples with different titled angle

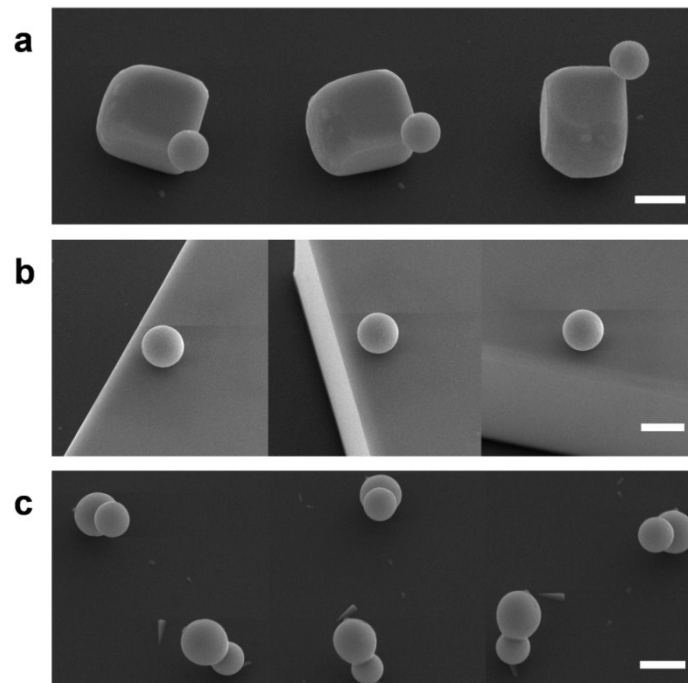


Figure S2 | SEM images of three sphere samples with different titled angle. The error bar is 1 μm .

III. Estimation of exciton binding energy

$$I(T) = R \left[1 - \exp\left(-\frac{E_b}{k_B T}\right) \right] + c,$$

The exciton binding energy can be fitted using equations,¹⁻² where $I(T)$ is the integrated PL intensity at a specific temperature T . Therefore, temperature-dependent steady-state PL spectroscopy is conducted on individual CsPbBr₃ microsphere in a backscattering configuration with an excitation laser of 405 nm. The temperature is varied from 77 K to 300 K. The integrated PL intensity versus $1/T$ of four samples are plotted in **Figure S3** (solid black dots), showing the exciton binding energy from 65 ~ 86 meV. These values are much higher than the thermal disturbance at room temperature ($KT \sim 26$ meV).

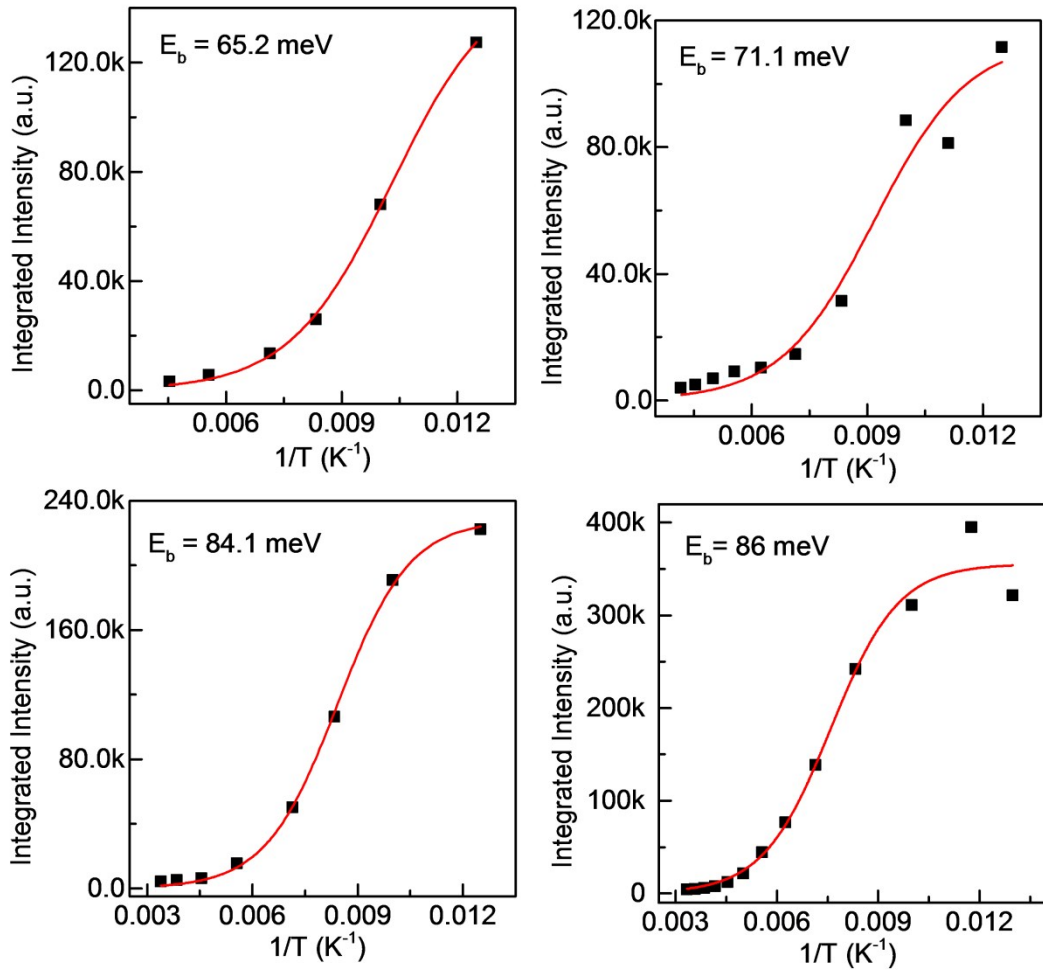


Figure S3 | The integrated PL intensity of four samples as a function of $1/T$. The experimental data are well fitted.

IV. The power dependent PL

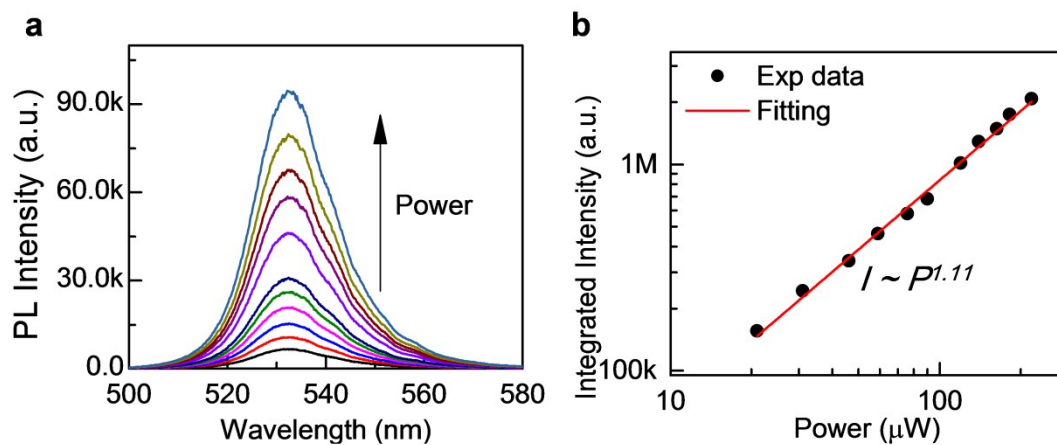


Figure S4 | (a) PL spectroscopy of CsPbBr₃ microsphere at room temperature with different laser power from 21 μW to 220 μW. (b) PL intensity as a function laser power in a log-log plot. The red line is a power-law fit $I=P^k$. Linear fitting of these plots gives a slope of 1.11 ± 0.03 , confirming exciton recombination process in these microspheres.

V. Temperature-dependent Raman spectra

To confirm the exciton-phonon coupling in sphere microcavities, Raman measurement is executed using 633 nm laser as the excitation source (**Figure S5**). As the CsPbBr₃ microsphere is heated from 103 K to 303 K, the splitting peaks near 70 cm⁻¹ of the Raman spectra are undistinguishable by degrees.

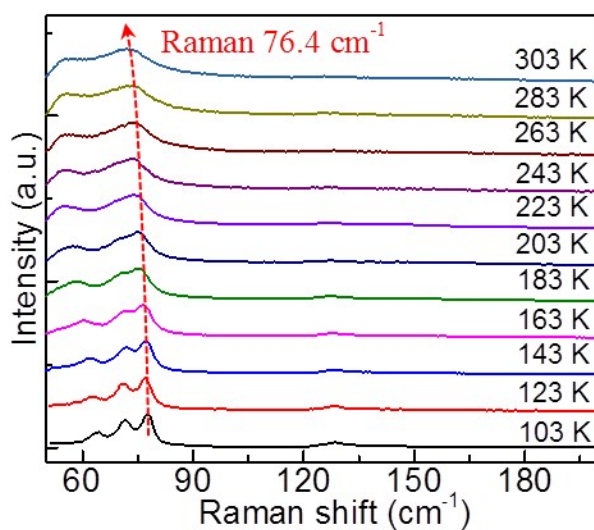


Figure S5 | Raman spectra of a typical CsPbBr₃ sphere at different temperatures from 103 to 303 K, showing the Raman shift as the temperature changes.

VI. The anharmonic model

In the anharmonic model, the variation of phonon frequency $\omega(T)$ and width $\Gamma(T)$ can be described by an anharmonic expression:

$$\omega(T) = \omega_0 + \Delta(T) \quad (\text{S1})$$

$$\Delta(T) = A \left\{ 1 + \frac{2}{e^x - 1} \right\} + B \left\{ 1 + \frac{3}{e^y - 1} + \frac{3}{(e^y - 1)^2} \right\} \quad (\text{S2})$$

for temperature dependence of wave number at 76 cm^{-1} , and

$$\Gamma(T) = C \left\{ 1 + \frac{2}{e^x - 1} \right\} + D \left\{ 1 + \frac{3}{e^y - 1} + \frac{3}{(e^y - 1)^2} \right\} \quad (\text{S3})$$

for temperature dependence of line-width at 76 cm^{-1} , where

$$x = \frac{\hbar\omega_0}{2K_B T} \quad (6)$$

$$y = \frac{\hbar\omega_0}{3K_B T} \quad (7)$$

ω_0 is the characteristic frequency of the mode; A, B, C and D are constants.

VII. Time-resolved PL spectroscopy and fitting

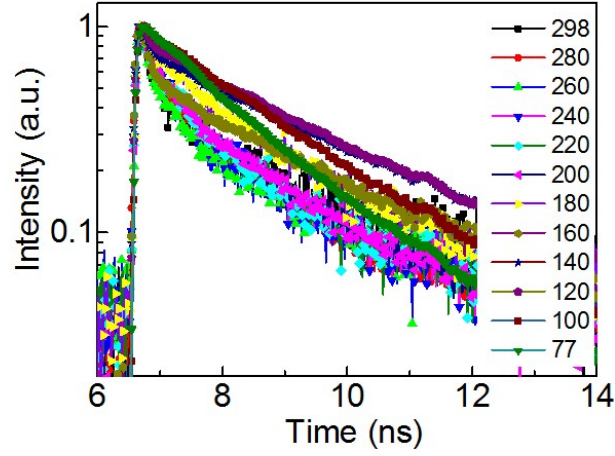


Figure S6 TRPL decay curves of CsPbBr₃ spheres with different temperatures from 77 K to 298 K monitored around 530 nm with excitation at 400 nm (76 MHz, 120 fs).

The PL decay curve in **Figure S6** can be well fitted by multi-exponential decay equation as follows:

$$I(t) = I_0 + A_1 \cdot e^{-\tau_1/t} + A_2 \cdot e^{-\tau_2/t} \quad (\text{S4})$$

$$T = \frac{A_1 \cdot \tau_1^2 + A_2 \cdot \tau_2^2}{A_1 \cdot \tau_1 + A_2 \cdot \tau_2} \quad (\text{S5})$$

where I , T , τ_i , A_i are the integrated PL intensity, average lifetime, lifetime and pre-exponential factors of CsPbBr₃ spheres at different temperature. The fitting results are shown in **Table S1**.

Table S1 The fitting results by multi-exponential decay equation

	298K	280K	260K	240K	220K	200K	180K	160K	140K	120K	100K	77K
τ_1 (ns)	2.26	1.63	1.61	1.52	1.62	1.95	1.98	2.69	2.98	2.80	2.32	2.78
τ_2 (ns)	0.23	0.18	0.21	0.23	0.25	0.23	0.31	0.14	0.15	0.12	1.12	1.32
A_1 (ns)	0.19	0.26	0.19	0.28	0.25	0.26	0.19	0.27	0.37	0.24	0.10	0.13
A_2 (ns)	0.39	0.41	0.40	0.34	0.32	0.37	0.34	0.44	0.26	0.42	0.41	0.41

VIII. The setup of optical measurement system

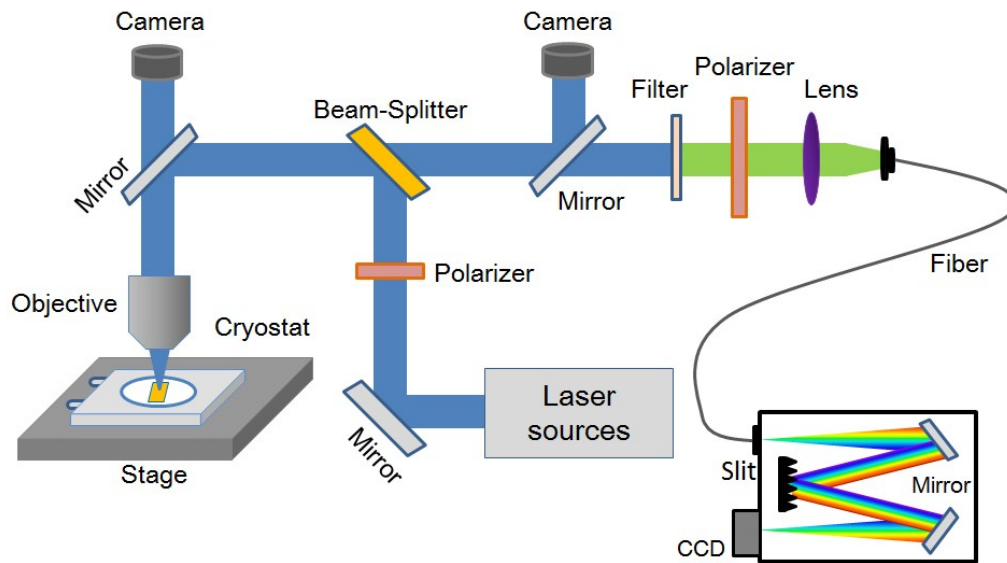


Figure S7 The setup of optical measurement system. For steady-state PL spectroscopy, a CW laser (wavelength: 405 nm) is focused onto an individual CsPbBr₃ sphere using an Olympus BX51 microscope equipped with a 100× objective (NA = 0.95). PL emission signal was collected by the same microscope objective in a backscattering configuration and analyzed by Princeton Instrument spectrometer (PI Acton, Spectra Pro 2500i) equipped with a TE-cooled charge coupled detector camera (PIXIS-400B). The PL image is recorded by a cool-snap color camera equipped on Olympus BX51 microscope. The PL lifetime measurements were conducted by a time-corrected single photon counting technique (TCSPC) with ultimate temporal resolution of ~40 ps. The excitation laser source is a frequency-doubled mode-locked Ti-sapphire oscillator laser (800 nm, repetition rate 76 MHz, pulse length 120 fs). For lasing measurements, the laser source is generated by frequency doubled from a Coherent Astrella regenerative amplifier (80 fs, 1 kHz, 800 nm) that was seeded by a Coherent Vitara-s oscillator (35 fs, 80 MHz).

References:

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- (2) Zhang, Q.; Ha, S. T.; Liu, X.; Sum, T. C.; Xiong, Q., Room-Temperature Near-Infrared High-Q Perovskite Whispering-Gallery Planar Nanolasers. *Nano Lett.* **2014**, *14*, 5995-6001.