

Electronic Supplementary Information

Magnetic interactions and *In Vitro* Study of Biocompatible Hydrocaffeic Acid-stabilized Fe-Pt Clusters as MRI Contrast Agents

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Superparamagnetic size limit for *fcc* Fe-Pt nanoparticles

Because in the literature no information about the superparamagnetic size limit (D_{SPL}) for *fcc* Fe-Pt can be found, we used experimentally determined blocking temperatures (T_b) from the temperature-dependent field-cooled and zero-field-cooled (FC/ZFC) magnetic susceptibility curves and effective anisotropy constant (K_{eff}) values from the literature for spherical *fcc* Fe-Pt NPs with different sizes to calculate D_{SPL} . When T_b and the size of the NPs are known, then the effective anisotropy constant K_{eff} for the non-interacting monodispersed NPs can be calculated using the following equation:

$$K_{\text{eff}} = \frac{25 k_B T_b}{V}, \quad (\text{S-1})$$

where k_B is the Boltzmann constant, T_b is the blocking temperature and V is the volume of the particle. Furthermore, **Eq. (S-2)** describing the Néel relaxation time can be used for the estimation of the SPL.

$$\tau_N = \tau_0 \exp\left(\frac{K_{\text{eff}} V}{k_B T}\right), \quad (\text{S-2})$$

The superparamagnetic behaviour occurs when $k_B T \gg K_{\text{eff}} V$ is valid and **Eq. (S-3)** can be rearranged into:

$$\left(\frac{K_{eff}V}{k_bT}\right) \geq \ln\left(\frac{\tau_m}{\tau_0}\right) \quad (S-3)$$

For the materials with high K_{eff} , such as *fct* Fe-Pt, *fct* CoPt, etc., in order to maintain the thermal stability of single-domain NPs, which is important, for example, in information storage [1], this ratio should be larger than 60:

$$\left(\frac{K_{eff}V}{k_bT}\right) \geq 60 \quad (S-4)$$

The diameter of a NP at the superparamagnetic size limit (D_{SPL}) can be calculated from the following equation:

$$D_{SPL} = \sqrt[3]{\frac{60 k_b T_b}{K_{eff}}} \quad (S-5)$$

Because the value of K_{eff} is shape dependent [2], to make a more solid comparison, only data for spherical *fcc* Fe-Pt NPs were used. The calculated D_{SPL} values listed in **Table S-1** are in the size range 8.8–12.5 nm. Deviation in the results can contribute to the errors coming from the determination of T_b . A broader size distribution means a broader peak in the ZFC curve and, consequently, a less accurate determination of T_b . Moreover, the position of T_b also depends on the measurement time, which brings an additional error in the calculation. However, from the obtained data we can conclude that the D_{SPL} for spherical *fcc* Fe-Pt NPs is approximately 10 nm; therefore, in order to keep the NPs in the superparamagnetic regime, this size should not be exceeded.

Table S-1: List of T_b and K_{eff} values from the literature for the spherical *fcc* Fe-Pt NPs with different sizes that were used in the calculation of D_{SPL} .

NPs size	T_b (K)	K_{eff} (J/m ³)	D_{SPL} (nm)	Ref.
3 nm	15	$3.7 \cdot 10^5$	8.8	[3]
3 nm	15	$2.0 \cdot 10^5$	9.2	[4],[5]
4 nm	25	$2.6 \cdot 10^5$	9.9	[6]
4.6 nm	53	$3.6 \cdot 10^5$	8.8	[7]
5 nm	30	$1.6 \cdot 10^5$	11.2	[8]
5 nm	26	$1.4 \cdot 10^5$	11.4	[2]
8 nm	70	$9.0 \cdot 10^4$	11.2	[8]
9 nm	140	$1.3 \cdot 10^5$	12.5	[9], [10]

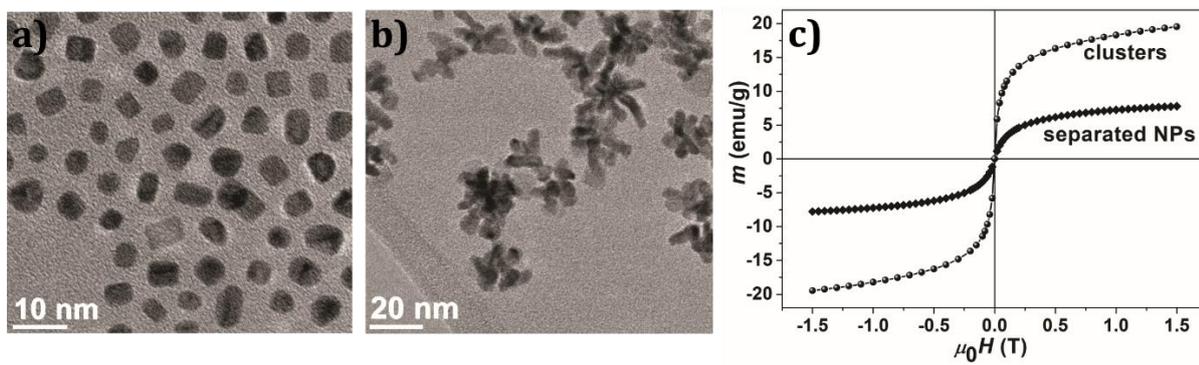


Figure S-1: TEM images of a) separated Fe-Pt NPs and b) Fe-Pt clusters and c) corresponding magnetic measurements at 300 K for both samples.

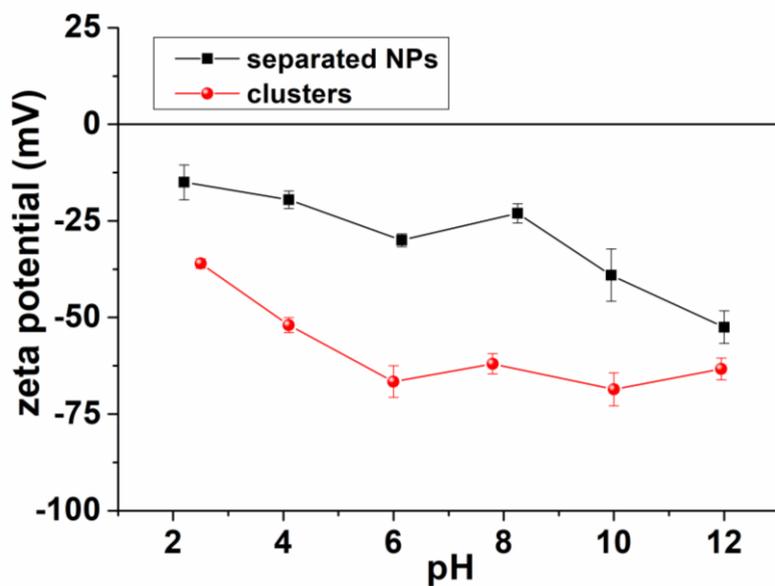


Figure S-2: Zeta-potential measurements of an aqueous suspensions of the separated Fe-Pt NPs and Fe-Pt clusters after ligand exchange with HCA.

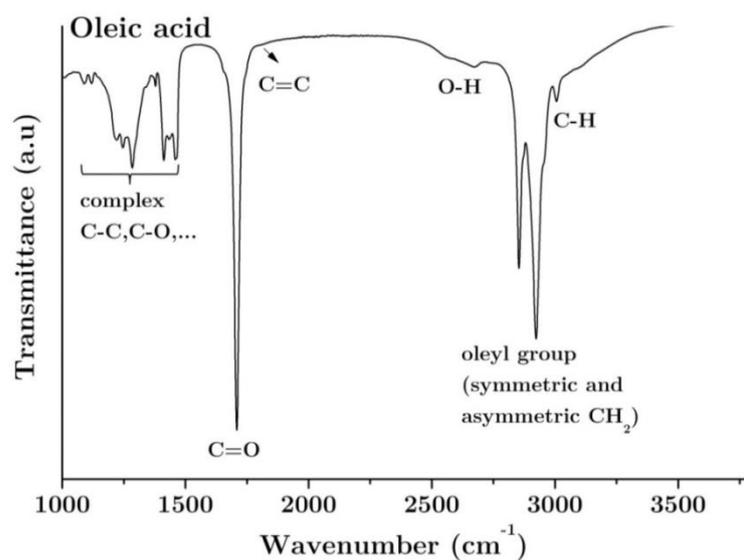


Figure S-3: FTIR spectrum of pure oleic acid (OA).

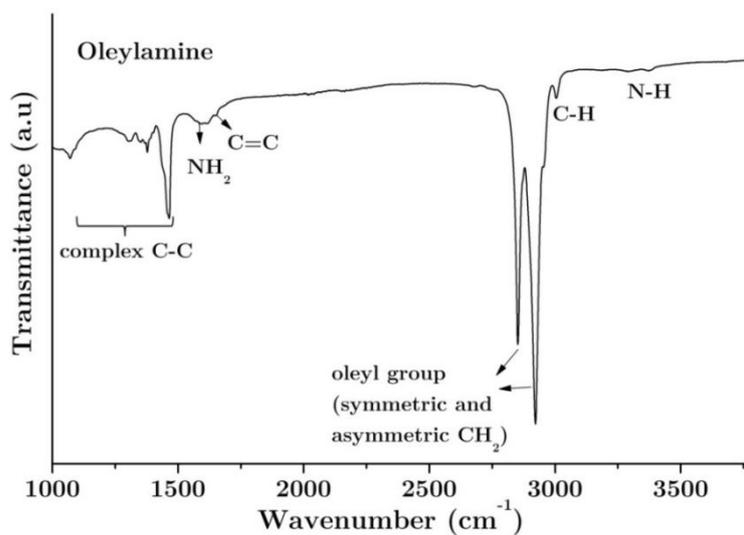


Figure S-4: FTIR spectrum of pure oleylamine (OLA).

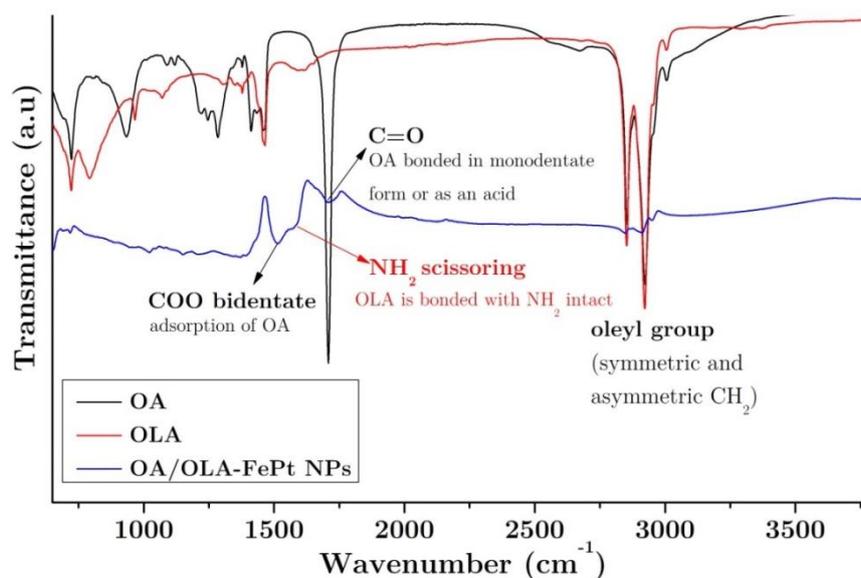


Figure S-5: FTIR spectra of pure oleic acid, pure oleylamine and oleic acid- and oleylamine-coated Fe-Pt NPs.

From the observation of $\nu(\text{COO})$ and $\nu(\text{C}=\text{O})$ vibrational modes it can be seen that OA bonds to the Fe-Pt NPs in both monodentate and bidentate forms (**Figure S-5**). More precisely, the peak at 1709 cm^{-1} corresponds to the $\nu(\text{C}=\text{O})$ stretch mode and indicates that OA is bonded on the Fe-Pt NPs, either in monodentate form or as an acid. The peak at 1512 cm^{-1} corresponds to the $\nu(\text{COO})$ mode and indicates the presence of the bidentate carboxylate bonding. The peak at 1590 cm^{-1} can be ascribed to the NH_2 scissoring, which suggests that OA is adsorbed, i.e., bonds to Fe-Pt NPs molecularly with the NH_2 group intact. This means that the OLA bonds to the Fe-Pt NPs through electron donation from the nitrogen atom of the NH_2 group [11]. Moreover, a broad peak with low intensity in the region $3000\text{--}3500 \text{ cm}^{-1}$ can be attributed to the $\nu(\text{NH})$ stretching mode, which confirms the previous argument about the NH_2 groups staying/being intact. Peaks at $2800\text{--}3000 \text{ cm}^{-1}$ correspond to the symmetric and asymmetric CH_2 stretching modes of the oleyl group; however, they give no useful information about the bonding of OA or OLA on the surface of Fe-Pt NPs.

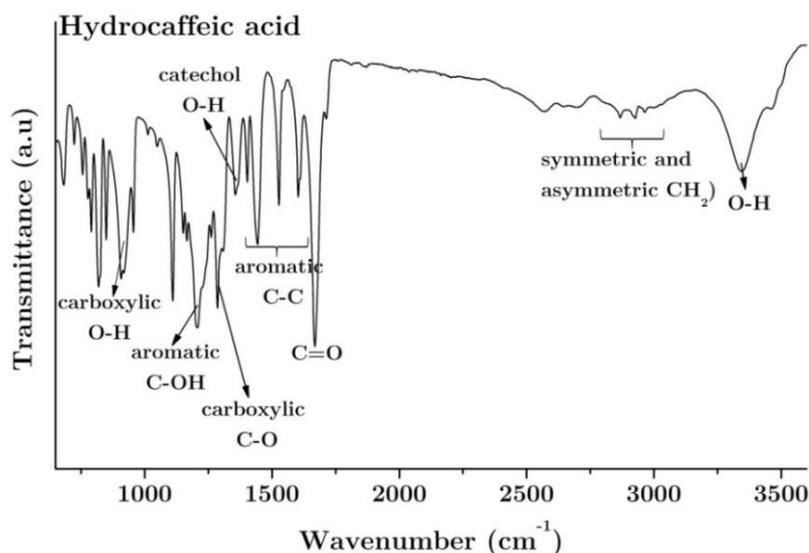


Figure S-6: FTIR spectrum of pure hydrocaffeic acid (HCA).

The FTIR spectrum of the pure HCA ligand is shown in **Figure S-5**. The broad peak centered at 3350 cm^{-1} can be attributed to the stretching vibration of the OH groups of HCA and the peak at 1360 cm^{-1} can be assigned to the bending vibrations of the catechol OH groups [11]. The peaks at $2800\text{--}3000\text{ cm}^{-1}$ correspond to the symmetric and asymmetric CH_2 stretching modes of the HCA molecule. The sharp peak at 1670 cm^{-1} corresponds to the $\nu(\text{C}=\text{O})$ stretch mode of the carboxylic group, which is found at a lower wavelength than expected for the monomeric acid ($\sim 1700\text{ cm}^{-1}$) due to the intramolecular hydrogen bonding. The peaks from $1400\text{--}1600\text{ cm}^{-1}$ appear due to the C-C stretching in the aromatic ring. The peaks at 1290 cm^{-1} and 1200 cm^{-1} correspond to the $\nu(\text{C}-\text{O})$ stretch mode of the carboxylic group and $\nu(\text{C}-\text{OH})$ of the OH catechol group, respectively [13].

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