## First-principles insights into tin-based two-dimensional halide perovskites for photovoltaics

Zhenyu Wang,<sup>†,‡,¶</sup> Alex M. Ganose,<sup>‡,¶,§</sup> Chunming Niu,<sup>†</sup> and David O.

 $\mathsf{Scanlon}^{*,\ddagger,\P,\S}$ 

<sup>†</sup>Xi'an Jiaotong University, Center of Nanomaterials for Renewable Energy, State Key Lab of Electrical Insulation and Power Equipment, School of Electrical Engineering, 99 Yanxiang Road, Xi'an 710054, China

<sup>‡</sup>University College London, Kathleen Lonsdale Materials Chemistry, Department of Chemistry, 20 Gordon Street, London WC1H 0AJ, UK

¶ Thomas Young Centre, University College London, Gower Street, London WC1E 6BT, UK

§Diamond Light Source Ltd., Diamond House, Harwell Science and Innovation Campus, Didcot, Oxfordshire OX11 0DE, UK

E-mail: d.scanlon@ucl.ac.uk

## Supporting Information

## Details of the spectroscopic limited maximum efficiency (SLME)

In the spectroscopic limited maximum efficiency (SLME) metric,<sup>1</sup> a photovoltaic cell is treated as an ideal diode illuminated under the incident photon flux  $I_{sun}$ , and the total current can be calculated as follows,

$$J = J_{sc} - J_0 (1 - e^{\frac{eV}{k_B T}}),$$
(1)

where V is the potential over the absorber thin film,  $k_B$  Boltzmann's constant, T the temperature, respectively.

 $J_{sc}$  is the the short-circuit current density given by,

$$J_{sc} = e \int_{0}^{\infty} a(E) I_{sun}(E) dE = e \int_{0}^{\infty} (1 - exp(-2\alpha(E)L)) I_{sun}(E) dE,$$
(2)

where  $\alpha$  and L are the absorption coefficient and thickness of the thin film, respectively.

 $J_0$  is reverse saturation current given by,

$$J_0 = f_r^{-1} J_0^r = e f_r^{-1} \int_0^\infty a(E) I_{bb}(E, T) dE = e f_r^{-1} \int_0^\infty (1 - exp(-2\alpha(E)L)) I_{bb}(E, T) dE, \quad (3)$$

$$f_r = e^{-\Delta/(k_B T)},\tag{4}$$

where  $I_{bb}$ ,  $f_r^{-1}$  and  $\Delta$  are the black-body radiation flux, fraction of the radiative recombination current and difference between the lowest direct allowed transition and the fundamental band gaps, respectively.

The maximum energy conversion efficiency, SLME, is given by,

$$SLME = \frac{[JV]_{max}}{P_{in}},\tag{5}$$

where  $[JV]_{max}$  and  $P_{in}$  are the maximum electrical output power density and the total incident solar power density, respectively.

Compounds	a (Å)	b(Å)	c(Å)	$\alpha$ (°)	$\beta$ (°)	$\gamma$ (°)
$(BA)_2 SnI_4^{rt}$	8.497	8.724	28.034	90.0	90.0	90.0
$(BA)_2SnI_4^{lt}$	8.410	8.911	26.074	90.0	90.0	90.0
$(BA)_2(MA)Sn_2I_7$	20.230	20.230	8.479	95.0	95.0	25.5
$(BA)_2(MA)_2Sn_3I_{10}$	8.551	9.085	25.307	90.5	97.7	90.9

Table S1: Lattice parameters of  $(BA)_2(MA)_{n-1}Sn_nI_{3n+1}$  unit cells using PBE-D3

Table S2: Lattice parameters of  $(BA)_2(MA)_{n-1}Sn_nI_{3n+1}$  conventional cells using PBE-D3 and PBEsol in comparison with experiment

Compounds	Functional	a (Å)	b(Å)	c(Å)	$\alpha$ (°)	$\beta$ (°)	$\gamma$ (°)
	PBE-D3	8.497	8.724	28.034	90.0	90.0	90.0
$(\mathrm{BA})_2\mathrm{SnI_4}^{\mathrm{rt}}$	PBEsol	8.519	8.691	28.411	90.0	90.0	90.0
	Ref. 2	8.591	8.814	27.644	90.0	90.0	90.0
	PBE-D3	8.410	8.911	26.074	90.0	90.0	90.0
$(BA)_2SnI_4^{lt}$	PBEsol	8.334	8.839	26.807	90.0	90.0	90.0
	Ref. 2	8.408	8.932	26.023	90.0	90.0	90.0
$(\mathrm{BA})_2(\mathrm{MA})\mathrm{Sn}_2\mathrm{I}_7$	PBE-D3	39.458	8.946	8.479	90.0	95.1	90.0
	PBEsol	40.026	8.968	8.394	90.0	92.8	90.0
	Ref. 3	39.497	8.858	8.776	90.0	90.0	90.0
$(\mathrm{BA})_2(\mathrm{MA})_2\mathrm{Sn}_3\mathrm{I}_{10}$	PBE-D3	8.551	50.184	9.085	89.4	90.9	92.0
	PBEsol	8.461	51.440	9.031	89.6	90.9	91.5
	Ref. 4	8.795	51.921	8.858	90.0	90.0	90.0

Table S3: Average in-plane and out-of-plane Born effective charges  $(Z^*)$  over Sn and I atoms. Superscripts  $\parallel$  and  $\perp$  indicate properties parallel (in-plane) and perpendicular (out-of-plane) to the 2D perovskite sheets, respectively

Compounds	Atom	$Z^*_{\parallel}$	$Z^*_\perp$
(DA) C <sub>p</sub> I rt	$\operatorname{Sn}$	4.916	2.601
$(DA)_2 SIII_4$	Ι	-2.035	-1.000
$(\mathbf{D}\mathbf{A})$ $(\mathbf{M}\mathbf{A})\mathbf{C}_{\mathbf{r}}$	$\operatorname{Sn}$	4.992	2.984
$(DA)_2(MA)SH_2I_7$	Ι	-2.000	-1.226
$(\mathbf{D}\mathbf{A})$ $(\mathbf{M}\mathbf{A})$ $\mathbf{C}_{\mathbf{n}}$ $\mathbf{I}$	$\operatorname{Sn}$	5.011	3.467
$(DA)_2(MA)_2511_31_{10}$	Ι	-2.197	-1.386



Figure S1: Crystal structures of the layered tin hybrid perovskites unit cells: (a)  $(BA)_2SnI_4^{rt}$ , (b)  $(BA)_2SnI_4^{lt}$ , (c)  $(BA)_2(MA)Sn_2I_7$ , and (d)  $(BA)_2(MA)_2Sn_3I_{10}$ 



Figure S2: Projected density of states (PDOS) and "inverted" crystal orbital Hamilton populations (-COHP) of Sn–I bonds for the layered tin hybrid perovskites: (a)  $(BA)_2SnI_4^{rt}$ , (b)  $(BA)_2SnI_4^{lt}$ , (c)  $(BA)_2(MA)Sn_2I_7$ , and (d)  $(BA)_2(MA)_2Sn_3I_{10}$ , where the energy is with respect to the Fermi level. Since the "inverted" COHP values are plotted, the positive regions (orange line) represent the bonding interactions , while the negative regions (blue line) denote antibonding interactions.



Figure S3: Charge density isosurfaces of the conduction band minimum (CBM) of the  $SnI_6^{4-}$  octahedra in (a)  $(BA)_2SnI_4^{rt}$  and (b)  $(BA)_2SnI_4^{lt}$ .

## References

- Yu, L.; Zunger, A. Identification of potential photovoltaic absorbers based on firstprinciples spectroscopic screening of materials. *Phys. Rev. Lett.* 2012, 108, 068701.
- (2) Takahashi, Y.; Obara, R.; Nakagawa, K.; Nakano, M.; Tokita, J.-y.; Inabe, T. Tunable Charge Transport in Soluble Organic–Inorganic Hybrid Semiconductors. *Chem. Mater.* 2007, 19, 6312–6316.
- (3) Stoumpos, C. C.; Mao, L.; Malliakas, C. D.; Kanatzidis, M. G. Structure-Band Gap

Relationships in Hexagonal Polytypes and Low-Dimensional Structures of Hybrid Tin Iodide Perovskites. *Inorg. Chem.* **2017**, *56*, 56–73.

(4) Mitzi, D. B.; Feild, C.; Harrison, W.; Guloy, A. Conducting tin halides with a layered organic-based perovskite structure. *Nature* **1994**, *369*, 467–469.