## Supplementary Materials for

## Biomimetic multifunctional electronic hair sensor

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## Note S1. Detailed derivation of stress distribution

According to the mechanical equilibrium, when a force is applied to the hair shaft, a pressure $(P)$ and a moment $(M)$ will be generated around the hair root. As shown in Fig. 1 B , when the pressure is perpendicular to the half plane, the state of stress at any point is expressed as ${ }^{S l}$ :

$$
\left\{\begin{array}{l}
\sigma_{\rho}=-\frac{2 P \cos \theta}{\pi \rho}  \tag{S1}\\
\sigma_{\theta}=0 \\
\tau_{\rho \theta}=\tau_{\theta \rho}=0
\end{array}\right.
$$

where $\rho$ and $\theta$ are polar radius and polar angle, respectively. Upon coordinate transformation, the stress $\left(\sigma_{x}^{P}\right)$ in the x direction is expressed as:

$$
\begin{equation*}
\sigma_{x}^{P}=-\frac{2 P}{\pi} \frac{x^{3}}{\left(\mathrm{x}^{2}+\mathrm{y}^{2}\right)^{2}} \tag{S2}
\end{equation*}
$$

where x and y are abscissa and ordinate, respectively. When the half plane is under moment $(M)$, the state of stress at an arbitrary point can be expressed as:

$$
\left\{\begin{array}{l}
\sigma_{\rho}=\frac{2 M \sin 2 \theta}{\pi \rho^{2}}  \tag{S3}\\
\sigma_{\theta}=0 \\
\tau_{\rho \theta}=\tau_{\theta \rho}=-\frac{M(\cos 2 \theta+1)}{\pi \rho^{2}}
\end{array}\right.
$$

After coordinate transformation, the stress $\left(\sigma_{x}{ }^{M}\right)$ in the x direction is expressed as:

$$
\begin{equation*}
\sigma_{x}^{M}=\frac{8 M}{\pi} \frac{x^{3} y}{\left(\mathrm{x}^{2}+\mathrm{y}^{2}\right)^{3}} \tag{S4}
\end{equation*}
$$

$P$ is arbitrarily supposed to equal 0.1 N and M is assumed to equal $0.1 \mathrm{~L} \mathrm{~N} . \mathrm{mm}$. The schematic of distributed force is presented in Fig. S22. The distributed force $\left(q_{1}\right)$ at the fixed end section of the hair induced by concentrated force is expressed as ${ }^{s 2}$ :

$$
q_{1}=\frac{P}{S}=\frac{0.4}{\pi d^{2}}
$$

where S is the cross sectional area of the hair and L is the length of hair. The distributed force $\left(q_{2}\right)$ at the fixed end section of the hair induced by pure moment is expressed as ${ }^{52}$ :

$$
q_{2}=\frac{M}{I} \eta=\frac{6.4 L}{\pi d^{4}} \cdot \eta
$$

where $I$ and $\eta$ are the moment of inertia and the distance from the point $\mathrm{M}(\mathrm{x}, \mathrm{y})$ to the origin of the coordinate on the AB line, respectively. The total distributed force $q(\mathrm{y})$ at the fixed end section of hair is expressed as:

$$
\begin{equation*}
q(\mathrm{y})=q_{1}+q_{2}=\frac{0.4}{\pi d^{2}}+\frac{6.4 L}{\pi d^{4}} \cdot \eta \tag{S5}
\end{equation*}
$$

The stress $\left(\sigma_{x}\right)$ in the x direction is expressed as:

$$
\sigma_{x}=-\frac{2 P}{\pi} \frac{x^{3}}{\left(\mathrm{x}^{2}+\mathrm{y}^{2}\right)^{2}}
$$

Therefore, $\sigma_{x}$ induced by distributed force is

$$
\sigma_{x}=-\frac{2}{\pi} \int_{-r}^{r} \frac{q x^{3} d \eta}{\left(\mathrm{x}^{2}+(\mathrm{y}-\eta)^{2}\right)^{2}}
$$

where $r$ is the radius of the hair, and $r=d / 2$. According to equations (S5), $\sigma_{x}$ is expressed as:

$$
\begin{aligned}
& \sigma_{x}=-\frac{2}{\pi} \int_{-r}^{r} \frac{\left(\frac{0.4}{\pi d^{2}}+\frac{6.4 L}{\pi d^{4}} \cdot \eta\right) d \eta}{x\left(1+\left(\frac{\mathrm{y}-\eta}{x}\right)^{2}\right)^{2}} \\
& =-\frac{2}{\pi} \int_{-r}^{r} \frac{-\frac{3.2 L x}{\pi d^{4}}\left[2 \cdot\left(\frac{y}{x}-\frac{\eta}{x}\right)\right] d \eta+\left(\frac{6.4 L y}{\pi d^{4}}+\frac{0.4}{\pi d^{2}}\right) d \eta}{x\left(1+\left(\frac{\mathrm{y}-\eta}{x}\right)^{2}\right)^{2}} \\
& =\left.\frac{6.4 L x}{\pi^{2} d^{4}} \cdot \frac{1}{1+\left(\frac{\mathrm{y}-\eta}{x}\right)^{2}}\right|_{-r} ^{r}-\frac{2}{\pi} \int_{-r}^{r} \frac{\left(\frac{6.4 L y}{\pi d^{4}}+\frac{0.4}{\pi d^{2}}\right) d \eta}{x\left(1+\left(\frac{\mathrm{y}-\eta}{x}\right)^{2}\right)^{2}}
\end{aligned}
$$

Let $t=\frac{\mathrm{y}-\eta}{x}$, then $\eta=y-t x, d \eta=-x d t$.

$$
\frac{2}{\pi} \int_{-r}^{r} \frac{\left(\frac{6.4 L y}{\pi d^{4}}+\frac{0.4}{\pi d^{2}}\right) d \eta}{x\left(1+\left(\frac{\mathrm{y}-\eta}{x}\right)^{2}\right)^{2}}=\frac{2}{\pi}\left(\frac{6.4 L y}{\pi d^{4}}+\frac{0.4}{\pi d^{2}}\right) \int_{\frac{\mathrm{y}}{x} \frac{\mathrm{y}+r}{x}}^{\frac{d t}{\left(1+t^{2}\right)^{2}}}
$$

Let $t=\tan a$, thus $d t=(\sec a)^{2} \mathrm{~d} a$.

$$
\left.\int_{\frac{y-r}{x}}^{\frac{y+r}{x}} \frac{d t}{\left(1+t^{2}\right)^{2}}=\left.\frac{1}{2}(\sin a \cos a+a)\right|_{a \min } ^{a \max }=\frac{1}{2}\left[t /\left(t^{\wedge} 2+1\right)+\arctan (t)\right]\right]_{\frac{y y}{x}}^{\frac{y+r}{x}}
$$

Thus, $\sigma_{x}$ is expressed:

$$
\begin{align*}
& \left.\sigma_{x}=\left.\frac{6.4 L x}{\pi^{2} d^{4}} \cdot \frac{1}{1+\left(\frac{\mathrm{y}-\eta}{x}\right)^{2}}\right|_{-r} ^{r}-\frac{1}{\pi}\left(\frac{6.4 L y}{\pi d^{4}}+\frac{0.4}{\pi d^{2}}\right)\left[t^{\prime}\left(t^{\wedge} 2+1\right)+\arctan (t)\right]\right]_{\frac{y+y}{x}}^{\frac{y+r}{x}}  \tag{S6}\\
& \left.=\left.\frac{6.4 \gamma x}{\pi^{2} d^{3}} \cdot \frac{1}{1+\left(\frac{\mathrm{y}-\eta}{x}\right)^{2}}\right|_{-r} ^{r}-\frac{1}{\pi}\left(\frac{6.4 \gamma y}{\pi d^{3}}+\frac{0.4}{\pi d^{2}}\right)\left[t^{\prime}\left(t^{\wedge} 2+1\right)+\arctan (t)\right]\right]_{\frac{y+y}{x}}^{\frac{y+r}{x}}
\end{align*}
$$

where the aspect ratio $(\gamma)$ be defined as $L / d$. The function relation is made among the stress, diameter and aspect ratio of hair.

Therefore, $\xi_{(x=0.2, y=r)}$ can be expressed as:

$$
\begin{aligned}
& \xi_{(x=0.2, y=r)}=-\frac{\left.\frac{6.4 \gamma \times 0.2}{\pi^{2} d^{3}} \cdot \frac{1}{1+\left(\frac{\mathrm{y}-\eta}{0.2}\right)^{2}}\right|_{-r} ^{r}-\frac{1}{\pi}\left(\frac{6.4 \gamma r}{\pi d^{3}}+\frac{0.4}{\pi d^{2}}\right)\left[5 d /\left(25 d^{\wedge} 2+1\right)+\arctan (5 d)\right]}{\frac{1}{\pi}\left(\frac{0.4}{\pi d^{2}}\right)[5 d /(25 d \wedge 2+1)+\arctan (5 d)]} \\
& =1+8 \gamma-\frac{16 d \gamma}{5 d+\left(1+25 d^{2}\right) \arctan (5 d)}
\end{aligned}
$$

For a given nylon fiber, $d=0.2 \mathrm{~mm}$, we can obtain

$$
\begin{equation*}
\xi_{(x=0.2, y=r)} \approx 1+1.8 \gamma \tag{S7}
\end{equation*}
$$

## Note S2. Effect of the Young's modulus of hairs on sensitivity

The hair shaft can be considered as cantilever column under a uniaxial load. The hair shaft is deflected under compression (Fig. S23), and it will return to their initial positions after releasing the applied load. The exact differential equation of the deflection curve is then given by ${ }^{\mathrm{S} 1}$ :

$$
\left\{\begin{array}{l}
E I \frac{d \theta}{d s}=-F y  \tag{S8a-c}\\
\frac{d x}{d s}=\cos \theta \\
\frac{d y}{d s}=\sin \theta
\end{array}\right.
$$

where $\theta$ is the angle between the slope of any point of cantilever column and the $x$ axis and s is the arc length measured from the free end. The second-order differential equation is given by taking the arc length's derivative of equation (S8a)

$$
\begin{equation*}
\frac{d^{2} \theta}{d s^{2}}=-\frac{F}{E I} \sin \theta \tag{S9}
\end{equation*}
$$

Let $k^{2}=F / E I$, here for a circular cross section, the moment of inertia is $I=\pi d^{4} / 64$, the integration of $\theta$ on both sides of equation (S9) gives

$$
\begin{equation*}
\frac{1}{2}\left(\frac{d \theta}{d s}\right)^{2}=k^{2} \cos \theta+C \tag{S10}
\end{equation*}
$$

The boundary conditions are:

$$
\left\{\begin{array}{l}
\frac{d \theta}{d s}=0  \tag{S11}\\
\theta=\alpha
\end{array}\right.
$$

By applying the above boundary conditions, we can obtain

$$
d s=-\frac{d \theta}{k \sqrt{2} \sqrt{\cos \theta-\cos \alpha}}
$$

Finally, we obtain

$$
L=\int d s=\frac{1}{2 k} \int_{0}^{\alpha} \frac{d \theta}{\sqrt{\sin ^{2} \frac{\alpha}{2}-\sin ^{2} \frac{\theta}{2}}}
$$

Let $p=\sin (\alpha / 2)$, and a new variable $\varphi$ is introduced and $\sin (\theta / 2)=p \sin \varphi$, then we get:

$$
\begin{equation*}
L=\frac{1}{k} \int_{0}^{\pi / 2} \frac{d \varphi}{\sqrt{1-p^{2} \sin ^{2} \varphi}}=\frac{1}{k} K(p) \tag{S1}
\end{equation*}
$$

where $K(p)$ is complete elliptic integral of the first kind. Then, we can obtain an approximate expression:

$$
\begin{equation*}
L=\frac{1}{k} \times \frac{\pi}{2} \frac{1}{\left(1-0.97 p^{2}\right)^{1 / 4}}=\frac{\pi}{2\left(1-0.97 p^{2}\right)^{1 / 4}} \sqrt{\frac{E_{f} I}{F}} \tag{S13}
\end{equation*}
$$

where $E_{f}$ is the Young's modulus of hair. Then, we can obtain ${ }^{33, S 4}$ :

$$
\begin{equation*}
p=\sin \frac{\alpha}{2}=\sqrt{\frac{1-\left(\frac{\pi}{2 L}\right)^{4}\left(\frac{E_{f} I}{F}\right)^{2}}{0.97}} \tag{S14}
\end{equation*}
$$

When the number of hairs is $n$, then we get

$$
\begin{equation*}
p=\sin \frac{\alpha}{2}=\sqrt{\frac{1-\left(\frac{\pi}{2 L}\right)^{4}\left(\frac{n E_{f} I}{F}\right)^{2}}{0.97}} \tag{S15}
\end{equation*}
$$

Meanwhile, $K(p)=1.8541$ can be obtained at $\alpha=\pi / 2$. Considering equation (S12) and $k^{2}=F / E I$, we can obtain

$$
\begin{equation*}
F=P_{\max }=\frac{n K(p)^{2} E_{f} I}{L^{2}} \tag{S16}
\end{equation*}
$$

Meanwhile, we can obtain

$$
y=\int \sin \theta d s=-\frac{1}{\sqrt{2} k} \int_{0}^{\alpha} \frac{\sin \theta d \theta}{\sqrt{\cos \theta-\cos \alpha}}=\frac{\sqrt{2}}{k}
$$

The aspect ratio $(\gamma)$ be defined as $L / d$, where $L$ and $d$ are hair length above the skin substrate and hair diameter, respectively. Thus, we can obtain

$$
\begin{equation*}
M_{\max }=P_{\max } y=\frac{\sqrt{2} n \pi K(p) E_{f} d^{3}}{64 \gamma} \tag{S17}
\end{equation*}
$$

where $n=30$ is the number of the hairs. For the given nylon fiber, we can obtain

$$
\begin{equation*}
M_{\max }=\frac{6.5}{\gamma} \tag{S18}
\end{equation*}
$$

The aspect ratio $(\gamma)$ can be obtained from equation (S16).

$$
\begin{equation*}
\gamma=\sqrt{\frac{n \pi E_{f} K^{2}(p) d^{2}}{64 P_{\max }}} \tag{S19}
\end{equation*}
$$

For a given operating range $F=0 \sim 0.1 \mathrm{~N}$, the aspect ratio $(\gamma)$ of the hairs is estimated to equal 41.3 in terms of formula (S19). Therefore, the fiber length above the skin surface is ca. 7.4 mm .

## Note S3. Effect of the Young's modulus of human skin on sensitivity

Hooke law is expressed as ${ }^{\mathrm{s} 1}$ :

$$
\begin{equation*}
\varepsilon=\frac{\sigma}{E} \tag{S20}
\end{equation*}
$$

where $\sigma$ and $\varepsilon$ are stress and strain, respectively. $E$ is Young's modulus. For the sensing properties of human skin, the smaller the modulus, the greater the strain, and the more sensitive to external stimuli. A minimum deformation amplitudes ( $\delta_{\min }$ ) of about 40 nm can be detected by human skin. ${ }^{55}$ Therefore, the minimum strain $\left(\varepsilon_{\min }\right)$ can be expressed as

$$
\begin{equation*}
\varepsilon=\varepsilon_{\min }=\frac{\delta_{\min }}{H} \tag{S21}
\end{equation*}
$$

where $H$ is the thickness of skin and $H \approx 2 \mathrm{~mm}$. Therefore, the minimum strain $\left(\varepsilon_{\min }\right)$ can be detected is about $2.0 \times 10^{-5}$. For human fingertips, the pressure of gentle touch is about $1 \mathrm{KPa} .{ }^{56}$ Therefore, we can obtain:

$$
\begin{equation*}
\sigma=E_{m} \varepsilon_{\min }<1 K P a \tag{S22}
\end{equation*}
$$

where $E_{m}$ is the modulus of skin. Therefore, we can obtain $E_{m}<50 \mathrm{MPa}$.
At the same time, the skin needs a certain modulus to protect and buffer the human body and maintain its shape. For human skin, gravity (G) can be expressed as:

$$
G=\rho g V
$$

where $\rho$ and $g$ are density of skin and gravitational acceleration, respectively. Then, we can obtain:

$$
\sigma_{G} S=\rho g S H
$$

where $S$ and $H$ are the area and thickness of skin, respectively. $\sigma_{\mathrm{G}}$ is stress induced by gravity (see Fig. S24). Therefore, we can obtain:

$$
\sigma_{G}=\rho g H
$$

Considering equation (S20), we can get:

$$
\varepsilon_{G}=\frac{\rho g H}{E_{m}}
$$

Therefore, the deformation ( $\delta$ ) of skin under gravity can be expressed as

$$
\begin{equation*}
\delta=\int_{0}^{H} \varepsilon_{G} d H=\int_{0}^{H} \frac{\rho g H}{E_{m}} d H=\frac{\rho g H^{2}}{2 E_{m}} \tag{S23}
\end{equation*}
$$

The average strain $\left(\bar{\varepsilon}_{G}\right)$ of the skin under gravity can be expressed as

$$
\begin{equation*}
\bar{\varepsilon}_{G}=\frac{\delta}{H}=\frac{\rho g H}{2 E_{m}}<\varepsilon_{\min } \tag{S24}
\end{equation*}
$$

Therefore, the strain of skin under gravity should be less than $\varepsilon_{\text {min }}$. Then, we can obtain:

$$
\begin{equation*}
\frac{\rho g H}{2 E_{m}}<2 \times 10^{-5} \tag{S25}
\end{equation*}
$$

For the human skin, $\rho \approx 1.0 \times 10^{3} \mathrm{~kg} / \mathrm{m}^{3}, H \approx 2 \mathrm{~mm}$, and $g \approx 10 \mathrm{~m} / \mathrm{s}^{2}$. Therefore, we can obtain $E_{m}>0.5 \mathrm{MPa}$.

Meanwhile, the mean failure strain ${ }^{\mathrm{S} 7}\left(\varepsilon_{\max }\right)$ of skin is about $50 \%$ and the plantar pressure ${ }^{\mathrm{S} 8}$ obtained from barefoot is about 0.1 MPa . In order to ensure that the skin is not damaged, the strain is

$$
\begin{equation*}
\varepsilon_{\max }=\frac{\sigma}{E_{m}}<0.5 \tag{S26}
\end{equation*}
$$

Therefore, we can obtain $E_{m}>0.2 \mathrm{MPa}$. According to equations (S22), (S25) and (S26), the modulus of skin ranges from $0.5 \mathrm{MPa} \sim 50 \mathrm{MPa}$, which is similar to the experimental results for the Young modulus $(0.1 \sim 18.8 \mathrm{MPa})$ of human skin. ${ }^{33}$ Considering Skin is a highly non-linear material. In the initial loading phase, skin has
very low modulus and large deformation occurs at a relatively low applied load, which will improve the sensitivity of skin.


Fig. S1. High magnitude SEM image of one human hair.


Fig. S2. Stress-strain curve of PDMS. The Young's modulus of PDMS is estimated to be about 0.3 MPa .


Fig. S3. Stress distributions along y axis. Stresses include $\sigma_{x}{ }^{\mathrm{P}}$ (green line) induced by pressure, $\sigma_{x}{ }^{M}$ (black line) caused by moment and their sum of $\sigma_{x}{ }^{P}+\sigma_{x}{ }^{M}$ (red line). (A) $x=1,(B) x=0.5,(C) x=0.2$ and $(D) x=0.1 \mathrm{~mm}$.


Fig. S4. Dependence of stress $\left(\sigma_{\mathbf{x}}{ }^{\mathbf{P}}+\sigma_{\mathbf{x}}{ }^{\mathbf{M}}\right)$ on hair diameter. The result was obtained in terms of equation (S4) based on the aspect ratio $\gamma=41.3$ and $x=0.2 \mathrm{~mm}$.


Fig. S5. Dependence of stress $\left(\sigma_{\mathbf{x}}{ }^{\mathbf{P}}+\sigma_{\mathbf{x}}{ }^{\mathbf{M}}\right)$ on hair aspect ratio. The result was obtained from equation (S4) based on the hair diameter $d=0.2 \mathrm{~mm}$ and $x=0.2 \mathrm{~mm}$.


Fig. S6. The angle $(\alpha)$ of hair bending and the maximum force $\left(P_{\max }\right)$ as a function of Young's modulus $\left(\mathbf{E}_{\mathrm{f}}\right)$ of hairs. The results were obtained in terms of equations (S15 and S16) based on the number $(\mathrm{n}=30)$ and length $\mathrm{L}=10 \mathrm{~mm}$ of hairs, $\mathrm{I}=\pi \mathrm{d}^{4} / 64$ and the diameter $\mathrm{d}=0.2 \mathrm{~mm}$ of hairs. According to equation (S15), the angle ( $\alpha$ ) of hair bending decreases with increasing the modulus (arbitrary assumption $\mathrm{F}=0.1 \mathrm{~N}$ ), which indicates that the lower the modulus, the higher the sensitivity. According to equation (S16), the maximum force ( $\mathrm{P}_{\max }$ ) increases with increasing the Young's modulus, which indicates that the higher the modulus, the higher $\mathrm{P}_{\text {max }}$ is.


Fig. S7. Stress-strain curve of one nylon fiber as hair material. The nylon fiber is taken from a tooth brush. The Young's modulus and diameter of the nylon hair are 1.04 GPa and $180 \mu \mathrm{~m}$, respectively, which are close to those of the human hair ( $\mathrm{E} \approx 0.25 \sim 4.2$ $\mathrm{GPa}^{35}$ and $d=30 \sim 150 \mu \mathrm{~m}{ }^{36}$ ).


Fig. S8. Dependence of bending moment ( $\mathbf{N} \cdot \mathrm{mm}$ ) on diameter and Yong's modulus of hairs. The result was obtained in terms of equation (S17). When $E$ is taken as a variable, $d$ is fixed at 0.2 mm ; similarly, when $d$ is a variable, $E$ is fixed at 1.04 GPa .


Fig. S9. SEM image of the carbonized paper. Carbonized paper (CP) was prepared through carbonization of single layer tissue paper.


Fig. S10. EDX spectrum of the carbonized paper. EDX results reveal that the elements of $\mathrm{C}, \mathrm{B}, \mathrm{O}, \mathrm{N}, \mathrm{Na}, \mathrm{F}$ and Mg in the carbonized paper are $78.71,19.81,1.27$, $0.13,0.05,0.03$ and 0.01 atomic $\%$, respectively. It is obvious that carbon is the dominant element. In addition, the atomic content of $\mathrm{Mg}, \mathrm{F}, \mathrm{N}$ and Na constitutes only $0.22 \%$ of the total material, which is normal for carbonized paper obtained from biomass.


Fig. S11. The RCR response of the EH sensor with different $h$. $h$ is the distance of the piezoresistive material from the EH surface, and the EH sensor shows the maximum RCR when $\mathrm{h}=0.2 \mathrm{~mm}$.


Fig. S12. The sensitivity of the EH sensors with different $\mathbf{h}$. At $\mathrm{h}=0.2 \mathrm{~mm}$, the EH sensor shows the highest sensitivity, which is similar to human skin. ${ }^{31,32}$


Fig. S13. Theoretical and experimental verification of the maximum stress position along x-axis $(\boldsymbol{y}=\mathbf{0 . 5} \mathbf{m m})$. (A) Stress distribution of $f_{l}(x)=\sigma_{\mathrm{x}}{ }^{\mathrm{P}}+\sigma_{\mathrm{x}}{ }^{\mathrm{M}}$ at $y=0.5 \mathrm{~mm}$ and $f_{2}(x)=\sigma_{\mathrm{x}}^{\mathrm{P}}-\sigma_{\mathrm{x}}{ }^{\mathrm{M}}$ at $y=-0.5 \mathrm{~mm}$ along x axis. (B) The sensitivity of the EH sensors with different h . At $\mathrm{h}=0.5 \mathrm{~mm}$, the EH sensor shows the highest sensitivity.


Fig. S14. The RCR response of the EH sensor monitored by M2. (A) When bending is applied to M1 and (B) when cyclic loading is applied to M2.


Fig. S15. The RCR response of the $\mathbf{E H}$ sensor to a limit force of $\mathbf{1} \mathbf{m N}$. The results are similar to that $(1 \mathrm{mN})$ of the previous report. ${ }^{28}$


Fig. S16. Effect of humidity on the sensing performance of the EH sensor. When the humidity increases from $50 \%$ to $95 \%$, the relative resistance and sensitivity of the EH sensor is reduced by about $1 \%$ and $5 \%$, respectively.


Fig. S17. SEM image of plate A. The roughness is about 200 microns.


Fig. S18. The interval of grooves in plates B and C. The results are obtained in terms of the RCR curves monitored.


Fig. S19 The response and recovery time of the EH sensor to airflow. Under various air flowrates, the response and recovery time was recorded for the EH sensor.


Fig. S20 Response of the EH sensor to walking and jogging. The sensor is fixed on an arm of a man. When the man walks or jogs, the hairs will deform and cause the resistance change of the EH sensor. Thus, the motion signals can be monitored by the EH sensor.


Fig. S21. FT-IR spectrum of nylon fibers. In the frequency range below $1500 \mathrm{~cm}^{-1}$ the observed peaks have medium/weak intensities, while in the $1700-1500 \mathrm{~cm}^{-1}$ frequency range the observed peaks are strong and are assigned for amide I and II bands. The intense amide I and II bands (namely C-N stretching of amide II band at $\sim 1546$ $\mathrm{cm}^{-1}, \mathrm{C}=\mathrm{O}$ stretching vibration of amide I band at $\sim 1639 \mathrm{~cm}^{-1}$ ) are observed in the frequency range of $1700 \sim 1500 \mathrm{~cm}^{-1}$. The bands at $\sim 3034 \mathrm{~cm}^{-1}$ and $\sim 3387 \mathrm{~cm}^{-1}$ are of special interest because they are related to hydrogen bonding. N-H bending vibration of amide is at $\sim 3470 \mathrm{~cm}^{-1}$. $S 9, S 10$


Fig. S22 Half plane under distributed force. Stress distributions are analyzed in the Supplementary Materials.


Fig. S23 Schematic showing of hair shaft deformation under an applied force $\boldsymbol{F}$. When $\alpha=\pi / 2$, it reaches the maximum working force $\left(F_{\max }\right)$.


Fig. S24 Schematics of the structure of hair cells and parameters.

Table S1. Comparison of the present EH sensor with existing hair sensors.

| Materials | Force Sensitivity $\left(\mathrm{N}^{-1}\right)$ | Roughness ( $\mu \mathrm{m}$ ) | Airflow Sensitivity ( $\Omega \cdot \mathrm{s} / \mathrm{mm}$ ) | $\begin{aligned} & \text { Resolution } \\ & (\mathrm{mN}) \end{aligned}$ | Air Direction | Ref. <br> No. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Cobalt/SR | 0.034~6.9 | NA | NA | $\sim 0.1$ | NO | 28 |
| rGO/HH | 0.0049 | NA | NA | 3000 | NO | 29 |
| Alloy/SR | 0.036 | NA | NA | NA | NO | S11 |
| PD/SR | NA | $0.5 \sim 50$ | NA | NA | NO | 39 |
| CNT/SR | NA | 15 | NA | NA | NO | S12 |
| Silicon/SR | NA | 200 | NA | NA | NO | S13 |
| CNT/SR | NA | 200 | NA | NA | NO | 27 |
| rGO/SR | NA | 200 | NA | $\sim 1$ | NO | S14 |
| Graphene/s ponge | NA | 200 | NA | NA | NO | S15 |
| Gold/silico <br> n | NA | NA | $\sim 0.079$ | NA | NO | S16 |
| CB/polyimi de | NA | NA | 0.066 | NA | NO | S17 |
| IN/SR | NA | NA | 0.17 | $\sim 0.1$ | NO | 30 |
| Nylon/SR | 0.7 | 200 | 0.1 | $\sim 1$ | YES | Present work |

Note: NA indicates not available; SR indicates silicone rubber, here is PDMS; rGO indicates reduced graphene oxide; HH indicates human hair; CNT indicates carbon nanotube; IN indicates iron nanowire; PD indicates polyvinylidene difluoride; CB indicates carbon black powder.

It is clear from Table S1 that most of existing hair sensors are either single modular or are not highly sensitive. By the contrary, the present biomimetic multi-responsive EH sensor is not only highly sensible but also capable of detecting multiple signals of pressure, surface roughness and airflow rate/direction, etc. like human shin. It should be pointed out that only in the present work, human skin is fully mimicked by using nylon fibers as hairs and PDMS as human skin. The nylon fibers have similar geometrical sizes, Young's modulus and embedded depth as human hairs in human skin. PDMS as base material for encapsulation of the hair sensor has also similar

Young's modulus as that of human skin. Therefore, it is understandable that the present EH sensor is multi-responsive and highly sensitive to pressure, surface texture and airflow, etc. like human skin.

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