

## Supporting Information

### Boron based layered electrode materials for metal-ion batteries

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## 1. The geometrical structures of various lithium borides

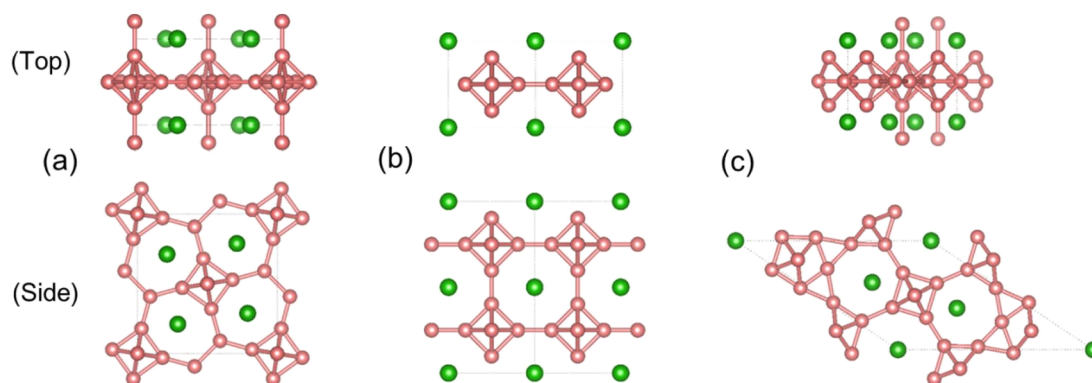


Figure S1. Top and side views of the geometric structures of lithium borides (a)  $MB_4$ ; (b)  $MB_6$ ; (c)  $M_3B_{20}$  ( $M = \text{Li, Na}$ ), respectively;

We perform a systematical investigation on various boron based layered compounds. As shown in Fig. S1, because of the structure features formed by boron octahedron and boron rings, typical metal borides including  $MB_4$ ,  $MB_6$  and  $M_3B_{20}$  ( $M = \text{Li, Na}$ ) have been considered. However, due to the dynamic instability of  $\text{LiB}_4$ ,  $\text{LiB}_6$  and  $\text{Li}_3\text{B}_{20}$  and high energy barrier of Na ion in  $\text{NaB}_4$ ,  $\text{NaB}_6$ ,  $\text{Na}_3\text{B}_{20}$ , they are restricted to the metal electrode materials.

## 2. The migration and phonon dispersion curves of LiB<sub>4</sub>

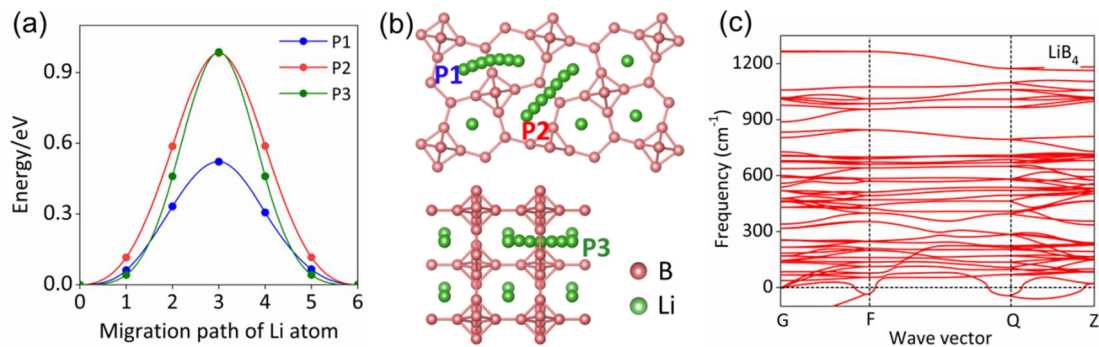


Figure S2. The diffusion energy barriers (a), three possible migration paths (b) of Li ions along the P1, P2, P3 and (c) phonon dispersion curve of LiB<sub>4</sub>.

Here, we show some results of LiB<sub>4</sub> as an example, as shown in Fig. S2 (a) and (b), with the diffusion path P1 indicated. It can be seen that just 0.52 eV energy barrier is needed to be overcome for Li ions, which is lower than that in LiAlB<sub>4</sub>. The barrier values of 0.98 eV and 0.99 eV for P2 and P3 paths were obtained, which are obviously higher than that along the P1 path. In addition, the phonon dispersion illustrates apparent imaginary modes. Thus, this compound fails in serving as electable electrode material.

### 3. Phonon dispersions of Na/MgAlB<sub>4</sub>, Na/MgGaB<sub>4</sub>

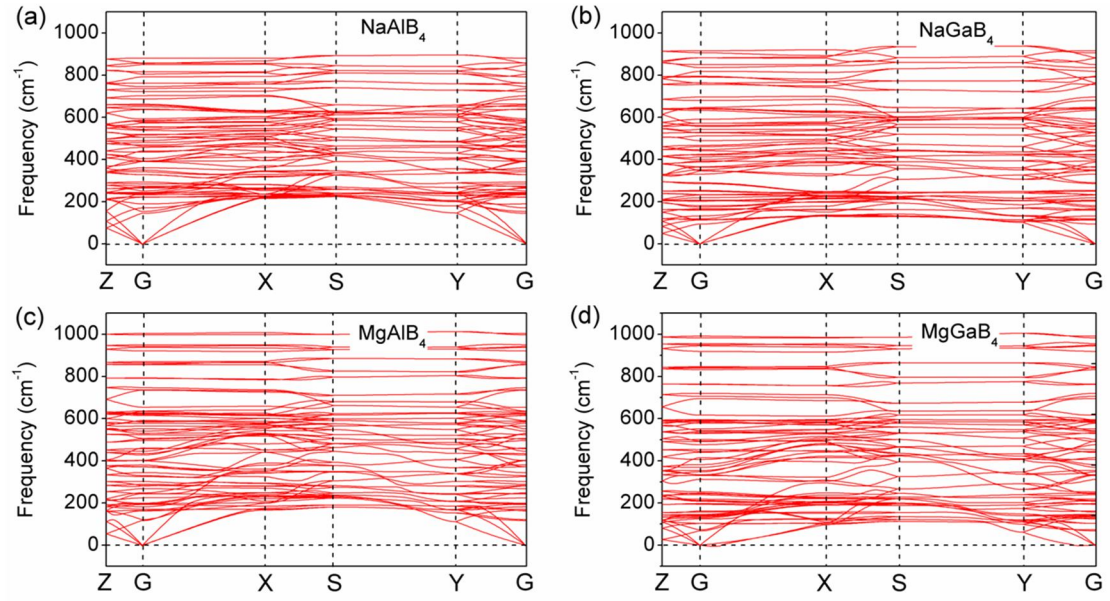


Figure S3. The calculated phonon dispersions for (a) NaAlB<sub>4</sub>, (b) NaGaB<sub>4</sub>, (c) MgAlB<sub>4</sub> and (d) MgGaB<sub>4</sub>, respectively.

In order to confirm the dynamic stability, firstly, the phonon dispersions are calculated for the structures of Na/MgAlB<sub>4</sub>, Na/MgGaB<sub>4</sub>, as presented in Fig. S3. No imaginary phonon modes are found, indicating that they all are dynamically stable.

#### 4. The electron density of states of Na/MgAlB<sub>4</sub> and Na/MgGaB<sub>4</sub>

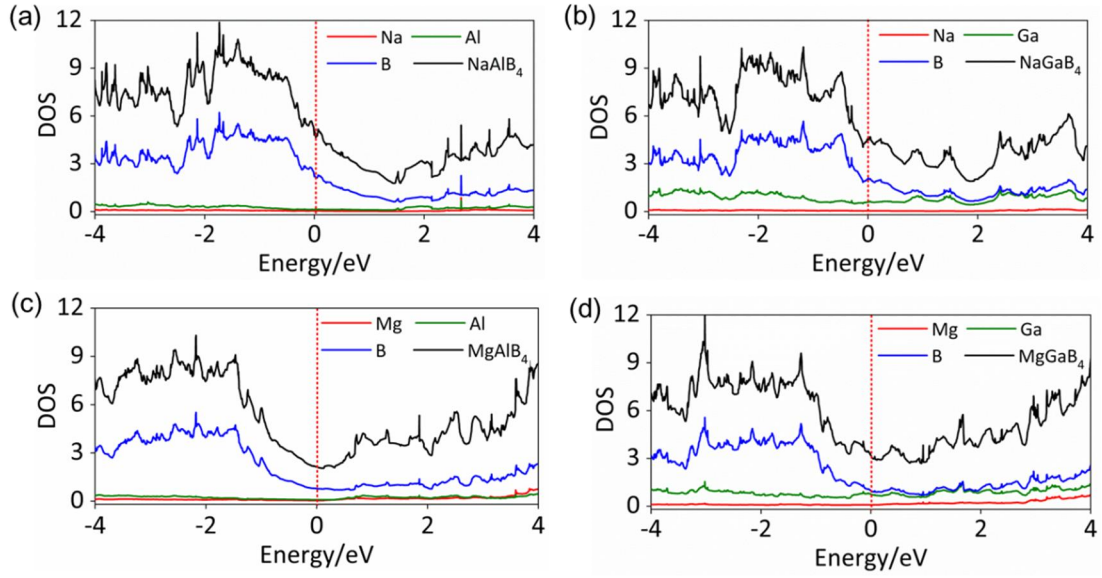


Figure S4. The total and projected electron density of states (DOS) of (a) NaAlB<sub>4</sub>, (b) NaGaB<sub>4</sub>, (c) MgAlB<sub>4</sub> and (d) MgGaB<sub>4</sub>, respectively.

As shown in Fig. S4, the calculated electron structures illustrate that they are metallic with electronic states existing on the Fermi level. The curves of DOS show that the contributions of the *s* electrons of Na and Mg atoms to the density of states especially valence states are very small and the boron electron states are prominent. Similarly, a comparison between MAIB<sub>4</sub> and MGaB<sub>4</sub> (*M* = Li, Na) also reveals that the DOS of Ga atoms are more dominant than that of Al atoms.

## 5. The electron band structures of $M_xXB_4$ ( $M = \text{Li, Na, Mg}$ ; $X = \text{Al, Ga}$ )

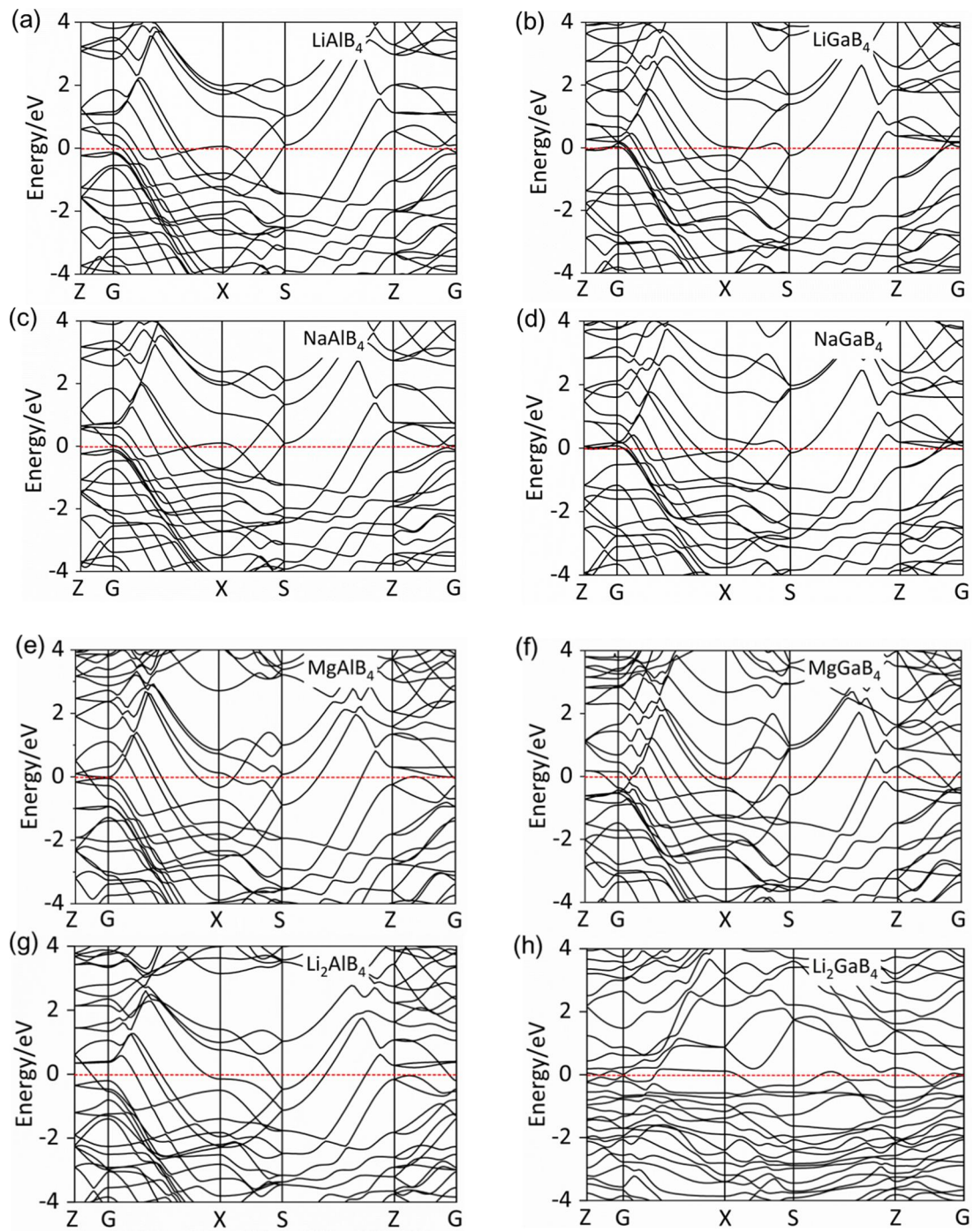


Figure S5. The electron band structures of (a)  $\text{LiAlB}_4$ , (b)  $\text{LiGaB}_4$ , (c)  $\text{NaAlB}_4$ , (d)  $\text{NaGaB}_4$ , (e)  $\text{MgAlB}_4$ , (f)  $\text{MgGaB}_4$ , (g)  $\text{Li}_2\text{AlB}_4$  and (h)  $\text{Li}_2\text{GaB}_4$ , respectively.



## 6. The volume change of $\text{Li}_x\text{AlB}_4$

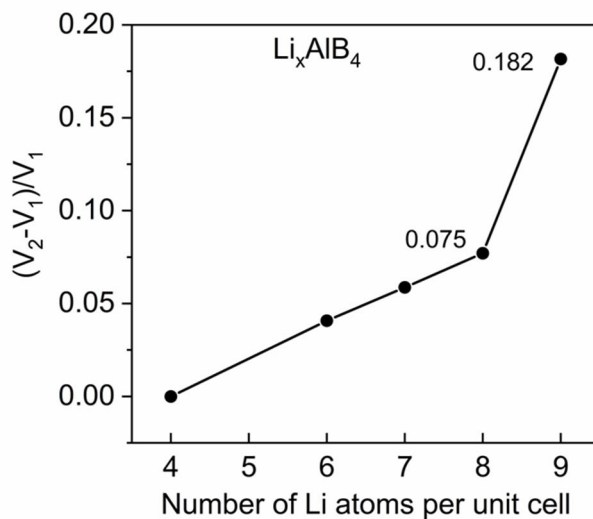


Figure S6. The volume change rate of  $\text{Li}_x\text{AlB}_4$ .

In order to further illustrate the structural stability of  $\text{Li}_x\text{AlB}_4$  during the process of lithium ions de-intercalation, the volume expansion was measured by  $\eta = (V_2 - V_1)/V_1$  ( $V_1$  is the volume of  $\text{Li}_4\text{AlB}_4$ ,  $V_2$  is the volume of  $\text{Li}_x\text{AlB}_4$  ( $x = 6, 8, 7, 9$ )). For  $\text{Li}_x\text{AlB}_4$  ( $x = 4, 6, 7, 8$ ), the structure remains intact and  $\eta$  grows linearly as shown in Fig. S6, and the value of  $\eta$  estimated as 7.5% for  $\text{Li}_8\text{AlB}_4$  goes well with the volume expansion tolerance. While  $x = 9$ , the structure is obviously distorted and  $\eta$  reaches to 18.2%, which may go beyond the lattice distortion and expansion tolerance and may be inappropriate for practical applications.



## 7. The calculated elastic constants of $M_xXB_4$ ( $M = \text{Li, Na, Mg}$ ; $X = \text{Al, Ga}$ )

	$\text{LiAlB}_4$	$\text{NaAlB}_4$	$\text{MgAlB}_4$	$\text{LiGaB}_4$	$\text{NaGaB}_4$	$\text{MgGaB}_4$	$\text{Li}_2\text{AlB}_4$	$\text{Li}_2\text{GaB}_4$
$C_{11}$	271.25	262.22	380.99	283.46	257.52	246.86	339.29	85.026
$C_{22}$	372.67	358.82	431.89	348.04	335.40	386.31	284.78	297.69
$C_{33}$	363.11	375.84	421.47	343.99	354.48	369.79	358.83	259.62
$C_{44}$	46.92	71.32	38.59	25.58	46.26	5.45	188.58	39.3
$C_{55}$	184.41	176.63	200.78	163.77	163.11	167.13	28.51	131.92
$C_{66}$	45.24	68.19	42.96	20.31	414.54	11.50	15.83	53.27
$C_{12}$	38.51	41.15	1.37	37.15	39.04	35.49	151.25	43.98
$C_{13}$	36.21	41.11	7.40	35.44	43.65	44.89	1.05	60.30
$C_{23}$	87.75	65.80	126.03	93.97	74.21	114.47	-19.54	49.82

Table S1. The calculated elastic constants  $C_{ij}$  (in GPa) for orthorhombic crystal  $M_xXB_4$  ( $M = \text{Li, Na, Mg}$ ;  $X = \text{Al, Ga}$ ).

As we all know, for a given orthorhombic crystal, there are nine independent elastic constants ( $C_{11}$ ,  $C_{22}$ ,  $C_{33}$ ,  $C_{44}$ ,  $C_{55}$ ,  $C_{66}$ ,  $C_{12}$ ,  $C_{13}$ ,  $C_{23}$ ), which should satisfy the Born stability criteria for mechanical stability as flowing:

$$C_{ii} > 0 \quad (i = 1, 2, 3, 4, 5, 6);$$

$$C_{11} + C_{22} + C_{33} + 2(C_{12} + C_{13} + C_{23}) > 0;$$

$$C_{11} + C_{22} - 2C_{12} > 0; \quad C_{11} + C_{33} - 2C_{13} > 0; \quad C_{22} + C_{33} - 2C_{23} > 0;$$

The calculated elastic constants for  $M_xXB_4$  ( $M = \text{Li, Na, Mg}$ ;  $X = \text{Al, Ga}$ ) are listed in Table S1, which shows that all of these compounds satisfy all above mechanical stability conditions. Therefore, they are mechanically stable.

## 8. The density states of phonon for LiAl/GaB<sub>4</sub> and Li<sub>2</sub>Al/GaB<sub>4</sub>

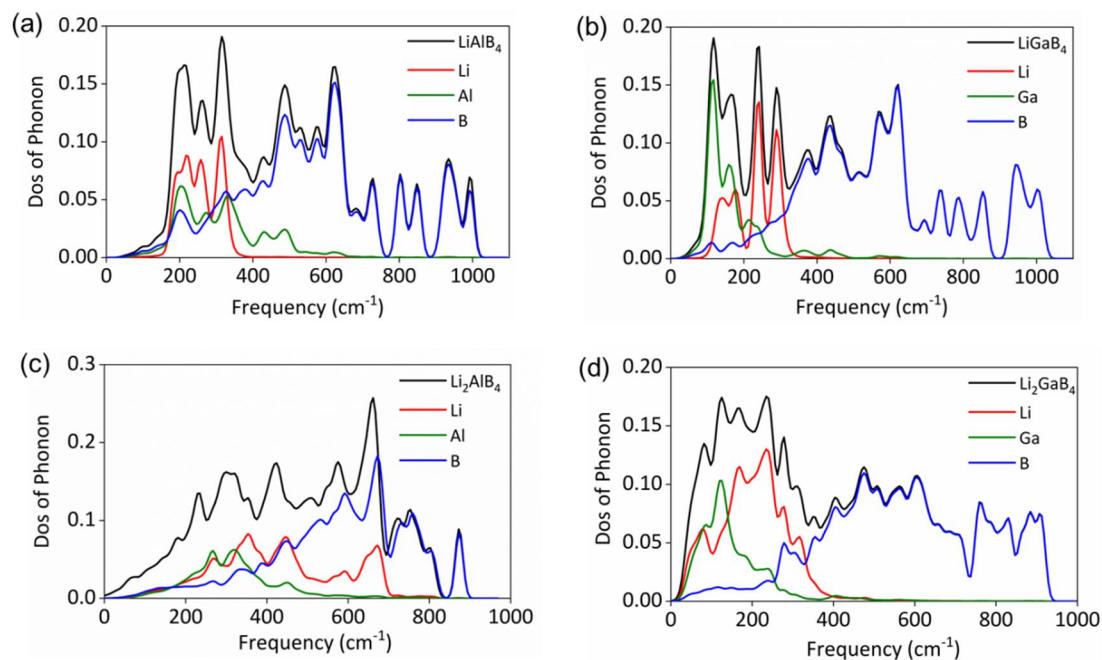


Figure S7. The total and projected phonon density of states for (a) LiAlB<sub>4</sub>, (b) LiGaB<sub>4</sub>, (c) Li<sub>2</sub>AlB<sub>4</sub> and (d) Li<sub>2</sub>GaB<sub>4</sub>, respectively.

As shown in Fig. S7, the phonon density states of Ga atoms mainly contribute in the low-frequency region and are larger than that of Al atoms. From Fig. S7 (a) (b) and (d), it is obviously seen that the primary contributions of phonon density of states in the low-frequency region come from Li and Al/Ga atoms, and the high frequency (>400 cm<sup>-1</sup>) is mainly from the boron atoms. In Li<sub>2</sub>AlB<sub>4</sub>, the distribution of DOS of phonon for Li atoms stretches across the range of 200 cm<sup>-1</sup> - 700 cm<sup>-1</sup> due to its structure.