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Supporting Information

for

Directing self-assembly in solution towards improved cooperativity in Fe(III) complexes with

amphiphilic tridentate ligands

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1. Cyclic voltammetry studies



Fig. S1 Cyclic voltammograms of **1** in dichloromethane at different scan rates using two different concentrations 1×10^{-3} M (left) and 4.9×10^{-3} M (right). Pt was used as working electrode (2 mm diameter), Pt wire and SCE as counter and reference electrodes, respectively. NBu₄PF₆ was used as supporting electrolyte.

2. Variable temperature UV-vis spectroscopy



Fig. S2 Variable temperature UV-vis spectra of **1** in a 4.4×10^{-6} M CH₂Cl₂ solution (cooling mode).



Fig. S3 Variable temperature UV-vis spectra of **1** in a $4.4x10^{-5}$ M CH₂Cl₂ solution (heating mode).



Fig. S4 Variable temperature UV-vis spectra of **1** in a 2.2×10^{-4} M CH₂Cl₂ solution. Left: cooling mode; Right: heating mode.



Fig. S5 Extinction coefficient vs. T for **1** at low concentration ($4.4x10^{-6}$ M), medium concencentration ($4.4x10^{-5}$ M) and high concentration ($2.2x10^{-4}$ M) CH₂Cl₂ solutions both in the cooling and heating modes.

3. Mössbauer spectroscopy

The magnetic properties of **1** in the solid-state and as a frozen solution in CH_2Cl_2 were investigated by Mössbauer spectroscopy. The ⁵⁷Fe Mössbauer spectra were recorded at 78 K, in transmission mode, using a conventional constant-acceleration spectrometer and a 50 mCi ⁵⁷Co source in a Rh matrix. The measurements were performed using a liquid nitrogen flow cryostat with a temperature stability of ±0.5 K. The velocity scale was calibrated using an α -Fe foil. The spectra were fitted to Lorentzian lines using the WinNormos software program, and the isomer shifts reported are relative to metallic α -Fe at room temperature.

The Mössbauer spectra of both samples are shown in Fig. S6. For the solid sample (bottom figure) the spectrum is well fitted using two quadrupole doublets. The blue one, corresponding to around 60% of the iron nuclei, is characterized by an isomer shift (δ) value of 0.22(1) mm s⁻¹ and a quadrupole splitting (ΔE_Q) of 2.91(1) mm s⁻¹, parameters characteristic of low-spin (LS) Fe ions. The remaining 40% of the iron nuclei, represented by the green sub-spectrum, has hyperfine parameters characteristic of Fe in the high-spin (HS) state ($\delta = 0.42(4)$ mm s⁻¹ and $\Delta E_Q = 0.45(4)$ mm s⁻¹). The frozen solution spectrum (top figure) is well fitted by a single quadrupole doublet displaying hyperfine parameters ($\delta = 0.21(2)$ mm s⁻¹; $\Delta E_Q = 2.81(4)$ mm s⁻¹) very close to the ones characterizing the LS state of the Fe ions of the solid sample, although with a much larger Lorentzian line width (Table S1). This much larger quadrupole doublet, when compared to the one of the solid compound, is interpreted not only as the result of being a measure of a frozen liquid but also as an indication of an existing fraction of a HS iron component, although in lower percentage than in the solid sample. The line broadening asymmetry has been previously found in other LS Fe(III) compounds,¹⁻⁴ and this asymmetry was attributed to the relatively long paramagnetic relaxation times of the iron when compared to the ⁵⁷Fe nuclear Larmor precession time.



Fig. S6 ⁵⁷Fe Mössbauer spectra collected at 78 K of A) solid sample of 1 and B) frozen solution (glass) of 1 in CH_2Cl_2 .

Table S1 Hyperfine parameters extracted from fitting the Mössbauer spectra shown in Fig. S6: δ - isomer shift; ΔE_{Q} - quadrupole splitting; Γ - Lorentzian line width at half maximum; W_{21} : width of line 2 relative to line 1; I - relative intensity (uncertainty < 2%). Statistical errors are given in parentheses.

Sample Solid	Fe	δ (mm/s)	ΔE_{Q}	Γ (mm/s)	W_{12}	I (%)
	subspectrum		(1111/5)	(1111/5)	(1111/5)	
	blue	0.22(1)	2.91(1)	0.49(3)	0.63(4)	58.8
	green	0.42(4)	0.45(4)	1.2(4)	-	41.2
Frozen solution in CH ₂ Cl ₂	blue	0.21(2)	2.81(4)	0.89(9)	0.58(6)	100

4. Particle size determination



Fig. S7 Size distribution of particles of 1 in a 4.9×10^{-3} M CD₂Cl₂ solution.



Fig. S8 Size distribution of particles of 1 in a $4.4x10^{-5}$ M CH₂Cl₂ solution.

5. Chronoamperometric studies



Fig. S9 Chronoamperograms of **1** in dichloromethane using two different concentrations 1×10^{-3} M (left) and 4.9×10^{-3} M (right). Inset: plot of current, I, versus t^{-0.5}, data collected between 4 and 8 s. Pt was used as working electrode (2 mm diameter), Pt wire and SCE as counter and reference electrodes, respectively. NBu₄PF₆ was used as supporting electrolyte.

6. Scanning Electron Microscopy imaging



Fig. S10 SEM images of aggregates of a 4.4×10^{-5} M CD₂Cl₂ of 1; a) magnification 30000x and b) magnification 30000x.





Fig. S11 SEM images of aggregates of a 4.9×10^{-3} M CD₂Cl₂ solution of **1**: a) magnification 500x and b) magnification 5000x.

7. Molecular Dynamics

The capped head and tail groups used to calculate the atomic charges are shown in Fig. S12. The force field was derived for the high spin state of the amphiphilic complex. Thus, the geometry of the capped head group used to calculate the charges was optimized for the high spin state. This geometry depicts a bond elongation of ~10% for the 2 Fe-N_{imine} bonds and a ~12 % and ~15 % elongation for the Fe-N_{amine} bonds, whereas an elongation of ~2.5% is observed for the Fe-O bonds, relative to the low spin state.



Fig. S12 (a) Fe(III) complex head capped with a CH_3 group and (b) tail capped with a C_6H_6 -O group, used to calculate the atomic charges. The atomic charges of the head were obtained through the RESP method by constraining the Fe atomic charge to + 2.0*e* and those of the tail were obtained with AM1-BCC.

The electrostatic potential (ESP) charge of Fe was found to be +1.4 *e* and the RESP charge was < +0.1 *e*. A radius of 1.5 Å, 2.0 Å, and 2.5 Å, were used for Fe, for comparison purposes, showing no significant influence on the Merz-Kollman charge. Furthermore, similar results are found when the charges are calculated for the complex and for the capped head. Thus, because of the low value of the RESP charge on the Fe atom, and the fact that ESP charges are difficult to assess for buried atoms the charge on the Fe atom was constrained to the value of +2.0*e* in the RESP approach. This value, while somewhat arbitrary, was observed from Mulliken population analysis, reflecting a milder charge transfer, relative to the ESP charge. The charge of -1*e* of the complex was achieved, upon combination of the RESP (head) and AM1-BCC (tail) charges, by absorbing the excess negative charge (~ -0.05*e*) to the Fe atom leading to a final charge on the Fe atom of +1.95*e*.

The opbe/cc-pvtz optimized Fe-X covalent bonds (X = N, O) length and the (Fe-X-Y) and (X-Fe-X) angles (Y=C, H) were adopted. The respective force constants were defined to be 700.0 kcal.mol⁻¹.Å⁻² and 90.0 kcal.mol⁻¹.rad⁻². These values were defined based on preliminary simulations insuring geometrical stability of the octahedral metal centre. The bond and angle force constants were systematically increased until no significant geometry distortions, ending with LINCS instability and ultimately bond breaking, were observed. The dihedral angles and force constants involving the Fe atom were set to zero as these were found to be unimportant for the geometrical stability of the head group. Although the above bond and angle force constants could in principle be derived more rigorously from electronic structure calculations, their improvement should have no influence on the self-assembling of the complex, which is the main

purpose of the MD simulations reported here.

The force field for the perchlorate anions and dichloromethane were also derived based on the GAFF protocol.

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