

WS₂ Monolayer Based Light Emitting Devices in a Vertical p-i-n Architecture

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Supporting Information

1. Photoluminescence dependence on adjacent layers

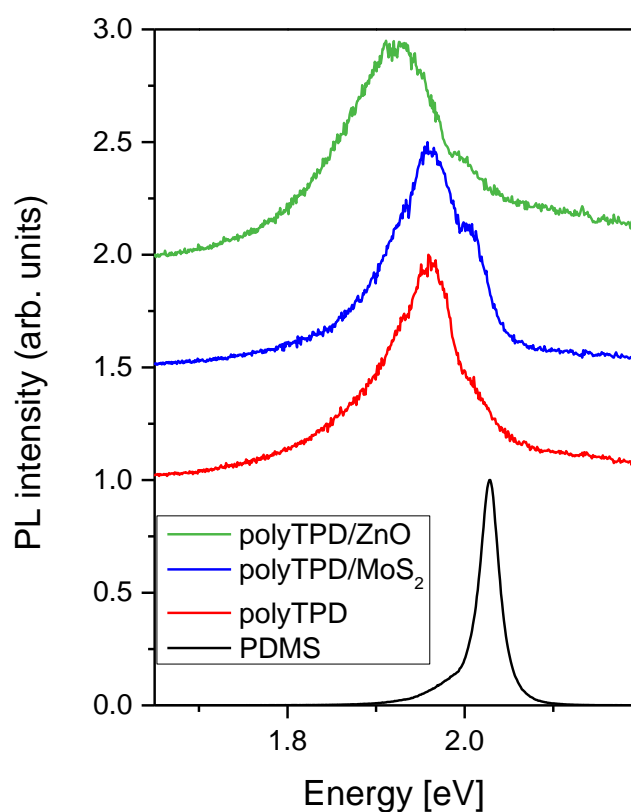


Figure S1. Change of PL emission during the process chain, showing the influence of the adjacent layers. Black: WS₂ on PDMS, Red: WS₂ on poly-TPD, Blue: WS₂ between poly-TPD and MoS₂ multilayer, Green: WS₂ between poly-TPD and ZnO-NP layer.

The spectra are normalized and vertically shifted for clarity.

The PL analysis of the active WS₂ monolayer at different stages of the process chain indicates the influence of the different surrounding on the luminescence spectrum. The exfoliated flake on PDMS (black) exhibits luminescence dominated by the A exciton with a maximum at 2.015 eV. In contrast, after WS₂ transfer onto TPD (red) we see two lines with maxima at 2 eV and 1.95 eV

that we attribute to the A-exciton, slightly shifted in energy possibly by the different dielectric surrounding, and the negatively charged X^- , which is shifted by another 50 meV towards lower energy. While the MoS₂-ESL on top of the WS₂ layer (blue) does not change the PL significantly with respect to the PL on poly-TPD, the PL in the sample with ZnO as ESL is clearly different: The PL maximum is here at 1.91 eV and the emission is significantly broadened to 130 meV. Possible reasons could be bound excitons that are enhanced by the ZnO/WS₂ interface, an enhanced contribution of trions, an energy shift due to the Quantum Confined Stark effect or a combination thereof.

2. Dependence of the EL signal on the MoS₂ thickness in the type II design.

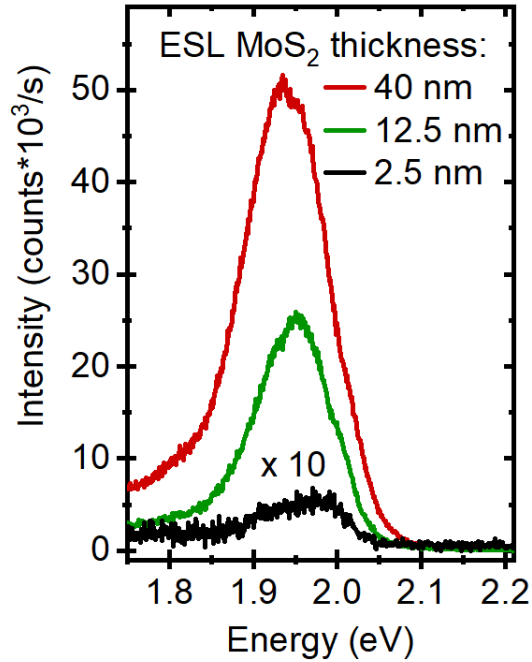


Figure S2. Electroluminescence spectra of WS₂ p-i-n devices with MoS₂ ESL layers of different thickness at 5.5 V. The intensity increases with increasing number of MoS₂ layers with a simultaneous peak shift to lower energies.

We studied the influence of the MoS₂-based electron supporting layer thickness on the electroluminescence of type II devices. As a striking result, only very weak EL is found for thin (2.5 nm) MoS₂-ESL. Obviously, the influence of the contacts is leading to a hole transfer to the contact and thus to an EL quenching [Ref. S1]. With increasing thickness, the electroluminescence becomes stronger, but with a distinct redshift. Here, we assume hole

transfer from WS₂ to the adjacent layers (poly-TPD and MoS₂), enhancing the probability for negatively charged excitons (trions) in the WS₂.

3. Estimation of internal quantum efficiency of the WS₂ monolayers

The collection efficiency of our setup was estimated by introducing a reflecting substrate (gold layer with reflectance of 80 % at a wavelength of 532 nm) and measuring the detected counts in the setup for a defined laser power.

The detected photon number per time can be written generally as

$$P_{det}/h\nu_{det} = (P_{exc}/h\nu_{exc}) \cdot IQE \cdot A \cdot \beta$$

Here $P_{det}/h\nu_{det}$ and $P_{exc}/h\nu_{exc}$ represent the photons per time detected in the setup and exciting the sample, respectively, with P_{det} (P_{exc}) as detected (exciting) power. A is the absorbance of the sample at the excitation wavelength, IQE the internal quantum efficiency in the sample and β the collection efficiency of the setup. In case of the reference with 80% reflection, $IQE \cdot A$ is set to 0.8 and β can be determined from measuring the counts in the detector and the power of the exciting laser at the position of the sample. From this, we determine a collection efficiency $\beta = 10^{-3}$. When characterizing the WS₂ monolayers by PL, the IQE can be calculated from the laser power, the detected intensity and the absorption in the monolayer that we assume to be 2% at the excitation wavelength of $\lambda=532$ nm [Ref. S2].

References

- S1 U. Bhanu, M. R. Islam, L. Tetard and S. I. Khondaker, Scientific Reports **4**, 5575 (2014).
- S2 H. Hill, A. Rigosi, C. Roquelet, A. Chernikov, T. Berkelbach, D. Reichman, M. Hybertsen, L. Brus and T. Heinz, Nano letters **15**, 2992 (2015)