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Supporting Information for

Twisted Non-diffracting Beams Through All Dielectric Meta-axicon

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Figure S1. (a) Normalized propagation constant or effective refractive index as a function of diameter for some low order modes of an indexed-waveguide. (b) Working principle of conventional Axicon. A conventional axicon with angle α_0 is illuminated from the bottom with light source at incidence angle \mathbb{Z}_{i} .



Figure S2. Phase profile under circularly polarized light and numerically simulated transmission amplitude for wavelength vs diameter.

Initial verification of the proposed meta-axicons also due to limited computational resources, a small portion (18 μ m × 18 μ m) of the designed metasurfaces with a lattice constant 300 nm is numerically simulated for topological charge m = 0, 2 and 4. Full-wave electromagnetic simulations are performed with perfectly matched layer (PML) boundary conditions in all directions. Each meta-axicon with different topological charge is numerically simulated for NA = 0.1, 0.7 and 0.9. A plane wave source with the operational wavelength of 633 nm is used to impinge the light on the meta-axicons from the substrate side. Cylindrical nano-resonators of a-Si:H with 400 nm thickness are tiled on a transparent glass substrate. The diffracted light from the array of nano-resonators is collected by monitors placed along the x-y and x-z planes respectively. Numerically simulated results of the all-dielectric meta-axicons under linear and circular polarization state of the incident light for topological charge m = 0, 2 and 4 are presented in Fig. S3. For NA = 0.1, Fig. S3a-f describe the behavior of meta-axicon with different topological charges while Fig. S3g-l represents the same behavior for NA = 0.9. Fig. S3a, c and e describe the transverse intensity distribution of the diffracted light captured by the x-y monitor possessing the doughnut shape intensity pattern accompanied with its longitudinal intensity profile captured by the x-z monitor. These results show the behavior of non-diffracting BB's, for topological charge m = 0, 2 and 4 respectively under linear polarization of the source. Fig. S3b, d and f illustrate the same behavior but under circularly polarized illumination.



Figure S3. Numerically simulated results of the all-dielectric meta-axicons with different numerical apertures. (a-f) present the numerically simulated results for NA = 0.1 while (g-l) for NA = 0.9. (a), (c) and (e) represent the generation of non-diffracting twisted beams under linearly polarized plane wave for the topological charge of 0, 2 and 4. Transverse monitors placed along the x-y axis represents the intensity distribution while longitudinal monitor placed along the x-z axis show the non-diffracting behavior of the meta-axicons. (b), (d) and (f) represent the similar behavior but with the circularly polarized source. Therefore, as the topological charge increases, central defect or zero-intensity region and the helical phase fronts also increase.

A close inspection of the field distribution in Fig. S3 reveals that as topological charge increases from 0 to 4, the central defect along the direction of propagation increases and number of intertwined spirals also increases. Full-wave electromagnetic simulations demonstrate that high diffraction efficiency is observed in our meta-axicons where scattering performance of individual meta-axicon is determined by the constructive interference of the transmitted light. In conclusion, the proposed concept of polarization insensitive all-dielectric metasurfaces to generate non-diffracting twisted beams via meta-axicons is verified by having completely identical results under linear and circular polarization of the incident.



Figure S4. SEM images of the fabricated meta-axicons. (a-c) Illustrate the low magnification SEM images of the meta-a xicons with NA = 0.7 having topological charge m = 0, 2 and 4. The inset in each subfigure shows the high magnific ation version from the center of the each structure. The scale bar for low magnification images is 10 μ m while for t he inset is 2 μ m.



Figure S5. Numerically simulated results of the all-dielectric meta-axicons for numerical aperture NA=0.7. (af) present the numerically simulated results for NA = 0.7 while (g-l) represent the experimental results for generation of nondiffracting twisted beams under linearly polarized plane wave for the topological charge of 0, 2 and 4. Transverse monitors place d along the x-y axis represents the intensity distribution while longitudinal monitor placed along the x-z axis show the nondiffracting behavior of the meta-axicons.

We have also numerically verified the focusing performance of our meta-axicons, in terms of full width half maximum (FWHM), for the numerical apertures of 0.1, 0.7 and 0.9 as given in Fig. S6. The meta-axicon with the topological charge of m=0 is illuminated with circularly polarized plane wave source. The results for xy-plane and xz-plane intensity plots along with line-plots, depicting the FMWM, are represented in Fig. S6. It is found that FWHM of 3.579λ , 0.511λ , and 0.365λ can be obtained for the numerical apertures of 0.1, 0.7 and 0.9 respectively. These analyses agree well with our expected results that the higher numerical aperture meta-axicons can give highly confined spot (without requiring an additional component) leading towards a subwavelength FWHM which is impossible to achieve from the conventional bulky axicons.



Figure S6. Numerically simulated full width half maximum of the zeroth order beam for NA = 0.1, 0.7 and 0.9. (a-c) describes the intensity profile at focal plane. (d-f) describe the FWHM while (g-i) illustrate the intensity profile along the direction of propagation. It is realized that metasurface-based structures can provide subwavelength spatial resolution ultimately resulting high numerical aperture and even smaller FWHM.