

# Strains and Defects Engineered Monolayer Ni-MoS<sub>2</sub> for pH-universal Hydrogen Evolution Catalysis

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## Supporting information

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## Computational methods

Our theoretical calculations were based on the density functional theory (DFT)<sup>1</sup> as implemented in the Vienna *ab-initio* Simulation Package (VASP).<sup>2,3</sup> We used the projector augmented wave method (PAW)<sup>4,5</sup> to represent electron-ion interactions, and the Perdew-Burke-Ernzerhof (PBE) functional within the generalized gradient approximation (GGA)<sup>6</sup> was considered for exchange correlation interactions. In addition, the vdW-DF2 method was used to describe the van der Waals interaction.<sup>7-10</sup> The energy cutoff for the plane-wave basis set was chosen to be 500 eV. The structures were relaxed until the forces on each atom were less than 0.01 eV/Å and the maximum energy change was of the order of 10<sup>-6</sup> eV. To prevent interaction between two neighboring surfaces, a vacuum slab of 15 Å was employed in *z*-direction. The spin-polarization was considered in this work. The energy barriers for Heyrovsky, Tafel mechanism and water dissociation were determined using climbing image nudged elastic band (NEB) calculations.<sup>11</sup> The Gibbs free energy ( $\Delta G_H$ ) of the basal planes for both perfect and S vacancy-contained Ni-MoS<sub>2</sub> was calculated using supercells containing 4×4 and 6×6 unit cells, respectively. The *k*-points for the first Brillouin zone were sampled on a mesh grid of 4×4×1 and 2×2×1 for 4×4 and 6×6 supercells, respectively. In order to determine the dynamic stability of Ni-MoS<sub>2</sub>, we calculated the phonon spectrum based on the force constant approach using the software package Phonopy.<sup>12</sup> Finite temperature analysis of the system at 1200K was conducted through *ab-initio* molecular dynamics (AIMD). The AIMD simulations were performed using 3000 time steps with a 0.5 fs time step at 1200 K. We employed non-self-consistent G<sub>0</sub>W<sub>0</sub> correction<sup>13</sup> for MoS<sub>2</sub> to calculate the bandgap, and compare the PBE and GW bandgap energies. G<sub>0</sub>W<sub>0</sub> correction accounts for the many-body electron interactions but retains the input PBE wave functions, thus results in the more accurate bandgap. A unified *k*-point mesh 12×12×1 is adopted for the G<sub>0</sub>W<sub>0</sub>.

The adsorption energy ( $\Delta E_H$ ) is computed as<sup>14</sup>

$$\Delta E_H = E(*H) - E(*) - \frac{1}{2}E(H_2) \quad (1)$$

where  $E(*H)$  and  $E(*)$  are the total energy of a supercell with and without hydrogen adsorption, respectively, and  $E(H_2)$  is the total energy of a H<sub>2</sub> molecule.

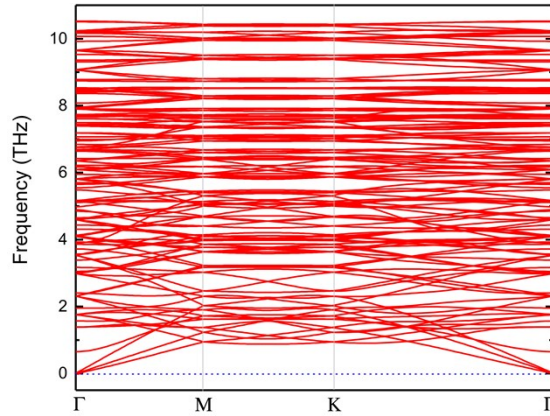
The Gibbs free energy of H ( $\Delta G_H$ ) is defined as:

$$\Delta G_H = \Delta E_H + \Delta E_{ZPE} - T\Delta S_H \quad (2)$$

where  $\Delta E_H$  is the adsorption energy,  $\Delta E_{ZPE}$  is the difference in zero-point energy, T is the temperature (300 K) and  $\Delta S_H$  is the entropy difference between H that is adsorbed and in the gas phase. We approximated the entropy of hydrogen adsorption as  $\Delta S_H \approx \frac{1}{2}(S_{H_2}^\circ)$ , where  $S_{H_2}^\circ$  is the entropy of gas phase  $H_2$  at standard conditions. Therefore, the correction factor of  $(\Delta E_{ZPE} - T\Delta S_H)$  was computed to be 0.223 eV in this study.

**Table S1.** Calculated values of lattice parameter ( $a$ ), bond length ( $d_{M-S}$ , M=Mo,Ni) for  $MoS_2$ ,  $NiS_2$ ,  $Ni-MoS_2$ , respectively, and formation energy ( $E_f$ ) of  $Ni-MoS_2$ .

	$a$ (Å)	$d_{M-S}$ (Å)	$E_f$ (eV)
$MoS_2$	3.192, 3.160 <sup>15</sup>	2.41, 2.42 <sup>15</sup>	
$NiS_2$	3.54, 3.40 <sup>16</sup>	2.58, 2.24 <sup>16</sup>	
$Ni-MoS_2$	3.351	2.418 (Mo-S) 2.279 (Ni-S) 3.857 (*Ni-S)	-1.222

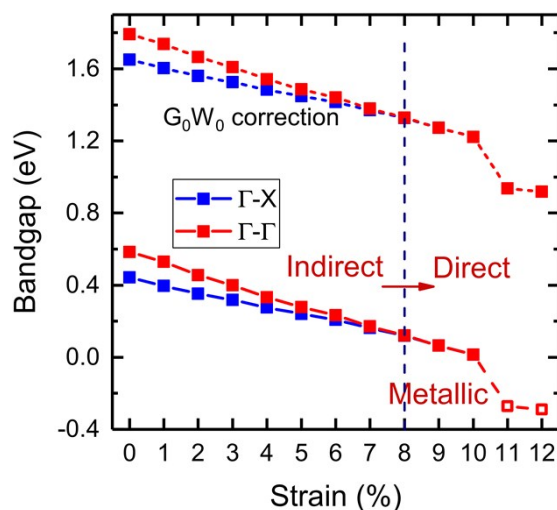


**Figure S1.** Calculated phonon-dispersion curves of  $Ni-MoS_2$  along major symmetry directions of the Brillouin zone.

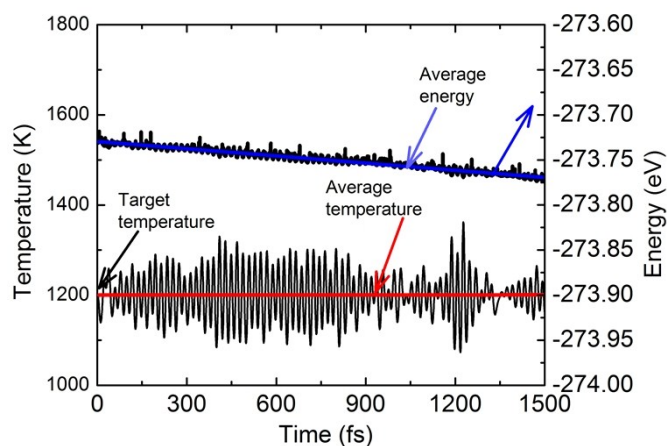
**Table S2.** Mechanical properties of  $Ni-MoS_2$ ,  $MoS_2$  and  $NiS_2$  predicted by first-principles calculations.

	Stiffness Tensor (N/m)	Young's Modulus (N/m)	Shear Modulus (N/m)	Poisson's Ratio
<b>Ni-MoS<sub>2</sub></b>	$C_{11} = 93.545$ $C_{22} = 62.575$ $C_{12} = 16.019$ $C_{66} = 16.694$	$Y_x = 89.444$ $Y_y = 59.832$	$Y_y = 16.694$	$\nu_x = 0.256$ $\nu_y = 0.171$
<b>MoS<sub>2</sub></b>	$C_{11} = 136.261$ $C_{22} = 136.260$	$Y_x = 128.178$ $Y_y = 128.177$	$Y_y = 51.754$	$\nu_x = 0.244$ $\nu_y = 0.244$

	$C_{12} = 33.188$			
	$C_{66} = 51.754$			
<b>NiS<sub>2</sub></b>	$C_{11} = 53.541$	$Y_x = 39.805$	$Y_y = 17.692$	$v_x = 0.515$
	$C_{22} = 51.852$	$Y_y = 38.549$		$v_y = 0.498$
	$C_{12} = 26.688$			
	$C_{66} = 17.692$			



**Figure S2.** Evolution of the bandgap energy as a function of applied tensile strains (%) within PBE (solid lines) and  $G_0W_0$  correction (dashed lines).



**Figure S3.** Variation of temperature and the total energy within 1500 fs during AIMD simulation around 1200 K for Ni-MoS<sub>2</sub>.

**Table S3.** Performance of MoS<sub>2</sub>, Pt(111) and Ni-MoS<sub>2</sub> catalysts for HER.

Catalyst	$\Delta G_H$ (eV)	Optimal condition ( $\Delta G_H \approx 0$ eV)	Bandgap (eV)	Water dissociation barrier (eV)
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MoS <sub>2</sub>	1.736 (present work) ~2 <sup>17</sup>	strain and S vacancy simultaneously <sup>17</sup>	1.87 (K-K) <sup>18</sup>	3.2 <sup>19</sup>
Pt(111)	~0 <sup>20</sup>		Metal	1.07 <sup>21</sup>
Ni-MoS <sub>2</sub>	0.545	11% strain or 2.5% S vacancy	0.443 (Γ-X)	1.114 (perfect); 0.866 (defective);

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