

Supplementary Information

Self-driven approach for local ion intercalation in vdW crystals

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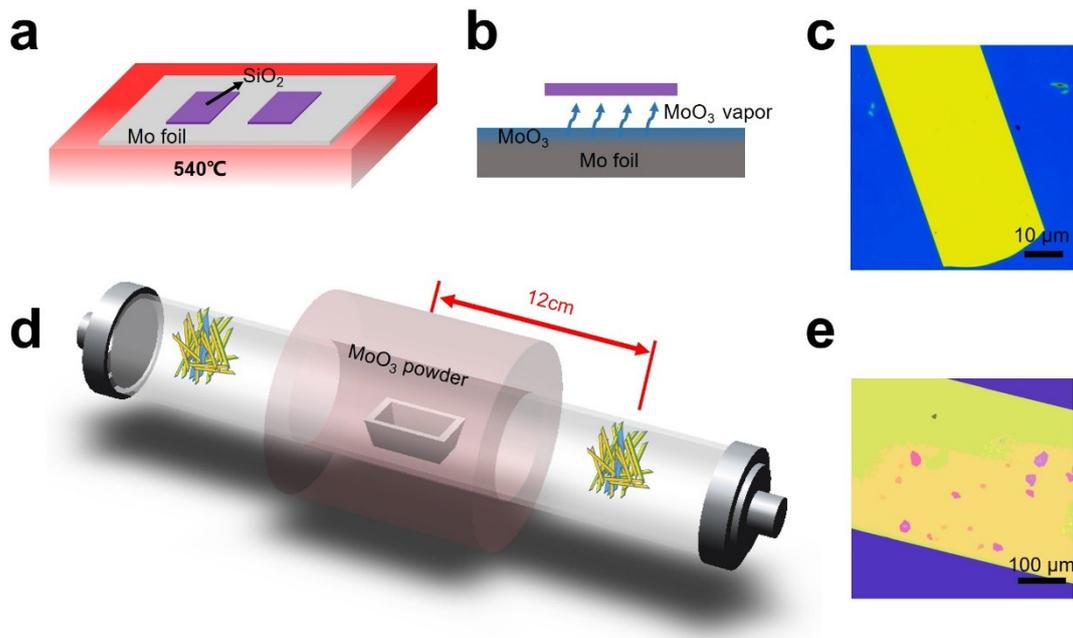


Figure S1 (a) & (b) Schematic illustration of hot-plate method to prepare MoO₃ nanosheets. (c) The optical image of the MoO₃ flakes. (d) PVD growth of MoO₃ single crystal. (e) The image of mass-produced MoO₃ single crystal

The MoO₃ in our experiment was prepared in two ways. The first was hot-plate method (Fig. S1a-c). In this method, a hot plate was heated to 540 °C and kept for a 10 minutes in atmosphere. A SiO₂/Si substrate was placed on a Mo foil with the SiO₂ facing down. It was then placed on the hot-plate for 2-5 mins. The Mo foil was oxidized and the MoO₃ nanosheet was deposited on the substrate. The sample was then taken out and cooled in air.

The second way was physical vapor deposition (PVD) preparation of MoO₃ single crystals (Fig. S1d&e). It involved the following steps: 0.3g MoO₃ powder was placed in the middle of the heating zone of the tube furnace. The tube was heated to 780 °C and held for 1 hour, and then naturally cooled down. High quality MoO₃ single crystal was deposited on the tube wall at a distance of 12 cm away from the powder source.

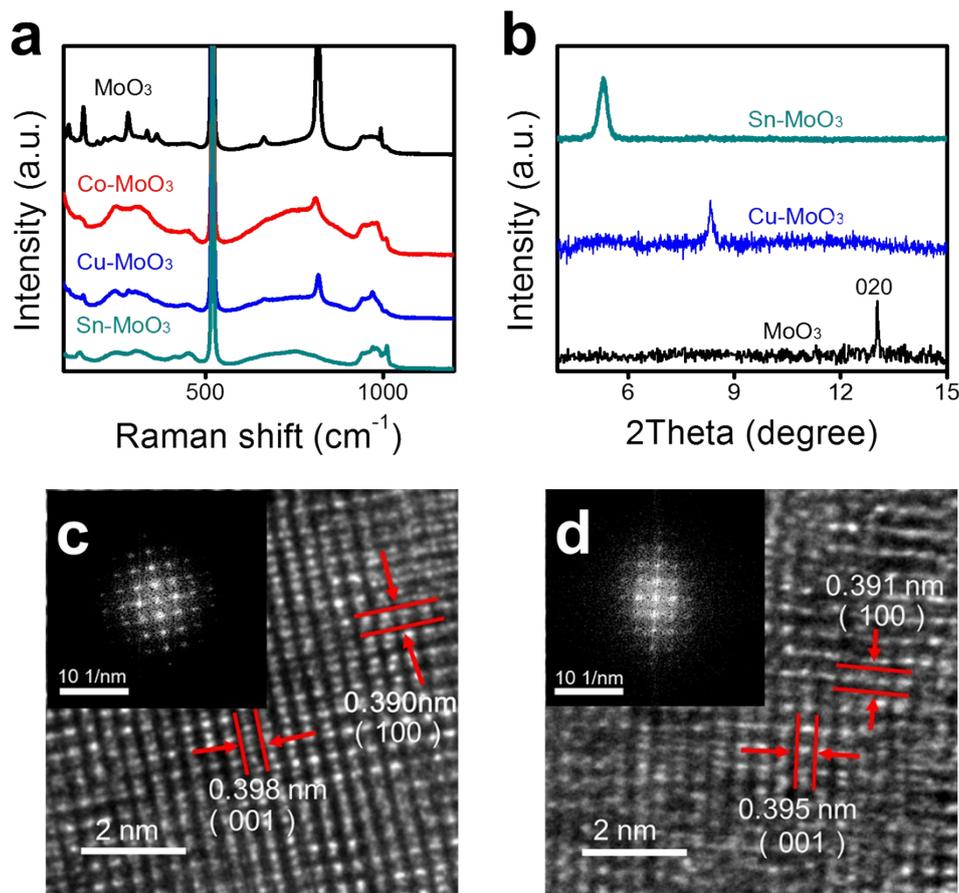


Figure S2 (a) Raman spectra of MoO₃, Co-MoO₃ Cu-MoO₃ and Sn-MoO₃. (b) XRD diffraction pattern of MoO₃, Cu-MoO₃ and Sn-MoO₃. (c) & (d) TEM diffraction images of Cu-MoO₃ and Sn-MoO₃. The insets show the corresponding SAED patterns. The intercalation concentration is 1×10^{-3} M

For Cu intercalation, 0.1 g Cu(NO₃)₂ powder (99.9% purity) was dissolved in 60 ml deionized water to obtain 10 mM Cu(NO₃)₂. For Sn intercalation, 0.15 g SnCl₄ powder (99.9% purity) was dissolved in 60 ml deionized water to obtain 10 mM SnCl₄. The solution was then diluted to desired concentration.

Intrinsic MoO₃ shows a Raman peak at 820 cm⁻¹ originated from the O₃-Mo-O₃ stretching vibration along the [100] direction, and a small peak at 993 cm⁻¹ due to the

Mo-O₁ vibration along [010] direction (Fig. S2a). After intercalation of Co, Cu or Sn ions, the intensity of the in-plane O₁-Mo-O₃ vibration significantly reduced, while the intensity of out-of-plane Mo-O₁ vibration changes slightly. It means that molecular layers are relaxed along the [010] directions so that it is slightly affected by the vdW gap intercalation. This relaxation cause the expansion along *b* axis as shown in Fig. S2b. All intercalated MoO₃ crystals show additional peaks at smaller angles. These peaks appear at different positions, which means that the expansion varies due to different interaction between the ion and the MoO₃.

The Sn and Cu intercalated sample show peaks at different positions, which means that the expansion of lattices is different. The (020) peak of MoO₃ is missing because the MoO₃ is fully intercalated. It should be noted that XRD require macro scale sample, while our method is not good at large scale intercalation. The intensity of the peaks can not be compared directly here because the amount of MoO₃ sample is small and there is great difference between samples. Both TEM images in Fig. S2c & d confirm significant lattice expansion along the [001] direction, which are consistent with that of Co-MoO₃. Disorder effect in the Sn-MoO₃ is stronger than the other two sample as shown by the blur diffraction spot in the SAED pattern. It shows that the intercalation effect depends on intercalation ion types.

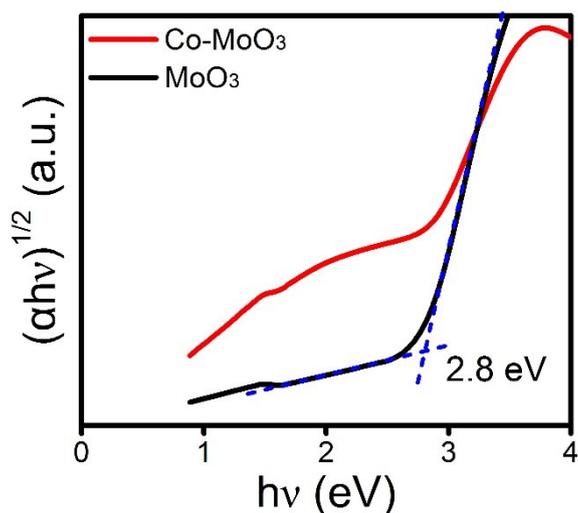


Figure S3 Optical band gap extraction of MoO₃ from Fig. 3a. The intercalated sample shows large gap states that narrow the optical gap. It's not able to extract its optical gap due to the measurement limit.

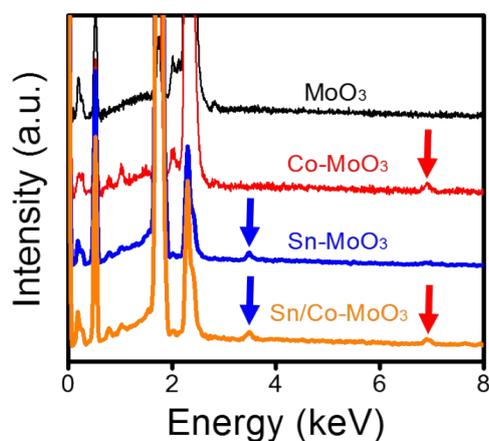


Figure S4 Energy dispersive spectroscopy (EDS) on different regions of MoO₃ nanosheets in **Fig. 4**. The red arrow indicates the peak of Co element and the blue arrow for Sn element. The atomic contents of Sn and Co are small. They can be determined roughly as Co_{0.13}MoO₃, Sn_{0.05}MoO₃, and Co_{0.8}Sn_{0.2}MoO₃. The Co content is higher than the Sn in single intercalated region.

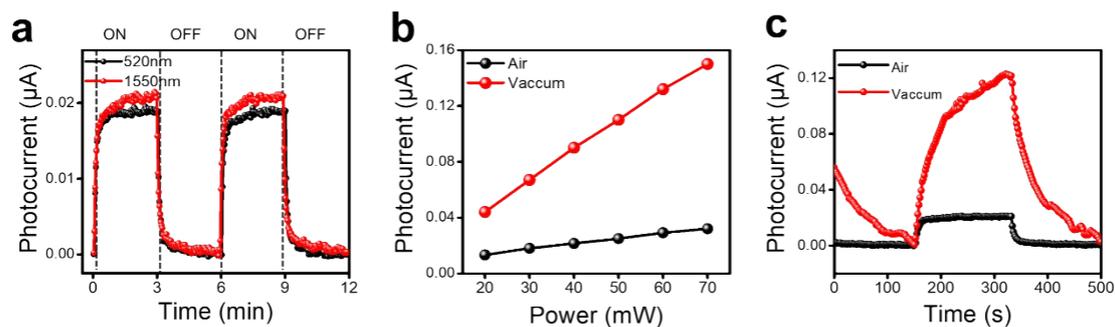


Figure S5 Broadband photoresponse of a Co-intercalated MoO₃ photoconductive device. (a) Transient photoresponse at a bias voltage of 0.1 V under 50 mW laser power in air under 520 nm and 1550 nm excitation. (b) Power dependence of photocurrent in air and vacuum under illumination of 1550 nm. (c) Photoresponse under air and vacuum condition.

Intrinsic MoO₃ shows no response to visible light, while the intercalated MoO₃ nanosheet shows broadband photoresponse to illumination wavelength of 520 nm and 1550 nm (Fig. S5a). The rise/decay response time under atmosphere is 21.8 s/27.0 s. Fig. S5b shows that the photocurrent linearly increases with increasing illumination power, giving a responsivity of 0.02 AW⁻¹ in air at 1550 nm. Photocurrent under vacuum condition is several times higher than that in air, giving a responsivity of 0.12 AW⁻¹. However, the rise/decay time is as slow as 1.82 min/1.88 min. This implies that there are localized states in the Co-intercalated MoO₃.

The difference of photo-current between vacuum and air is associated with the oxygen adsorption effect on the metal oxides. In an oxygen-rich atmosphere, oxygen molecules capture electrons from the surface of metal oxide, resulting in the formation of oxygen anion [O₂(g) + e⁻ → O₂⁻(ad)] layer on the surface. This effect causes a surface depletion layer. Under illumination, the excited electrons will be trapped by the surface layer, and leads to the decrease of photo-conductivity. In vacuum condition, the

adsorbed oxygen molecules will be desorbed, which reduces the depletion effect. Therefore, the photocurrent in vacuum is higher than that in air.

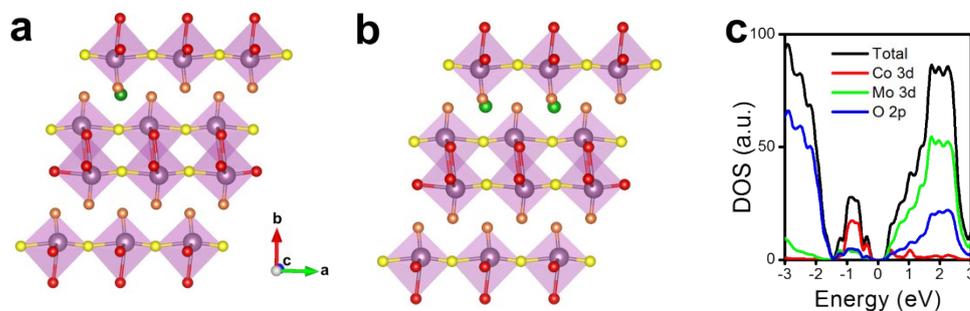


Figure S6 (a) The most stable intercalated site in MoO_3 for Sn or Co atom. (b) Two intercalated atoms in the cell. (c) The calculated DOS of intercalated MoO_3 with two Co atoms in (b).

The crystal structure of α - MoO_3 is composed of the distorted octahedral units, and each layer is connected through van der Waals interaction. There are three coordinated oxygen atoms in MoO_3 , represented by three different colors in Fig. S6 a&b. Our calculation shows that the interstitial site between two layers are the most stable intercalated site due to the large space of the van der Waals gap. The DOS was then calculated by assuming the intercalated atoms in the vdW gap.

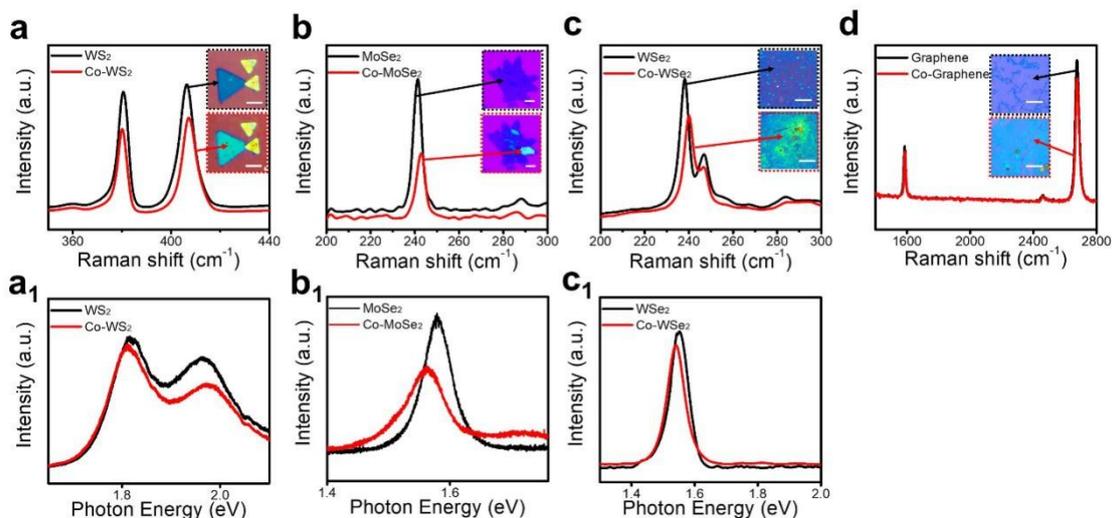


Figure S7 (a)-(d) Raman spectra of WS₂ and Co-WS₂ nanosheet, MoSe₂ and Co-MoSe₂ nanosheet, WSe₂ and Co-WSe₂ nanosheet and graphene and Co-graphene respectively. (a₁)-(c₁) Corresponding photoluminescence spectra of WS₂ and Co-WS₂ nanosheet, MoSe₂ and Co-MoSe₂ nanosheet, WSe₂ and Co-WSe₂ nanosheet in (a)-(c) respectively.

All the 2D materials are grown by chemical vapor deposition. To prepare graphene, Cu foils was annealed at 1000 °C in forming gas (20% H₂ + 80% Ar) at a pressure of 10⁻⁴ Torr. Triphenylene was introduced to the foil through electrical resistance heating at a vapor pressure of 10⁻⁵ Torr. After a 30 min reaction, the Cu foil was cooled down to room temperature at a base pressure of 10⁻⁷ Torr. The graphene was then transferred to SiO₂ surface for intercalation. To prepare MoS₂/MoSe₂/WS₂/WSe₂ Sulfur/selenium powder ($\approx 0.5 - 1$ g) and MoO₃/WO₃ powder (Alfa Aesar, 99.9995%, 5 - 50 mg) were put in different temperature zone of a homemade glass container with a long, narrow output tube (OD ≈ 3 mm). 300 nm - thick SiO₂/Si substrates (1 cm \times 1 cm) were placed downstream relative to the Ar gas (500 sccm) flow direction. The sulfur was evaporated by a heating belt at 120–150 °C, and 250 °C for selenium. The amount of

sulfur/selenium could be well controlled/reduced with the help of the narrow output tube. The growth temperature was raised to $\approx 800 - 900$ ° C to further enable the reaction between MoOx/Wox flakes with S/Se for crystalline sample for 0.5 - 3 h. Finally, the system was naturally cooled down to room temperature. More detail growth parameters of graphene, MoS₂, WS₂, MoSe₂ and WSe₂ and can be found in previous publications in Ref.[1-4].

For cobalt ion intercalation in other 2D materials, 0.1 g CoSO₄ powder (99.9 % purity) was dissolved in 60 ml deionized water to obtain 10 mM CoSO₄ solution. The solution was then diluted to desired concentration. Zn powder (99.9 % purity) with diameter of around 10 μ m was dispersed on the surface of WS₂, MoSe₂, WSe₂ and graphene nanosheet. A drop of CoSO₄ solution was dropped onto the nanosheet. Intercalation was observed under an optical microscope. After intercalation, the nanosheets showed distinct color change relative to the intrinsic ones. The change of Raman spectra and photoluminescence spectra was shown in Fig. S7.

It's clear that first, the color of the 2D materials change after intercalation. The reduction of Raman intensity are close to what has been observed in MoS₂. It implies that intercalation has altered the interlayer coupling. The reduction of PL intensity gives the support.

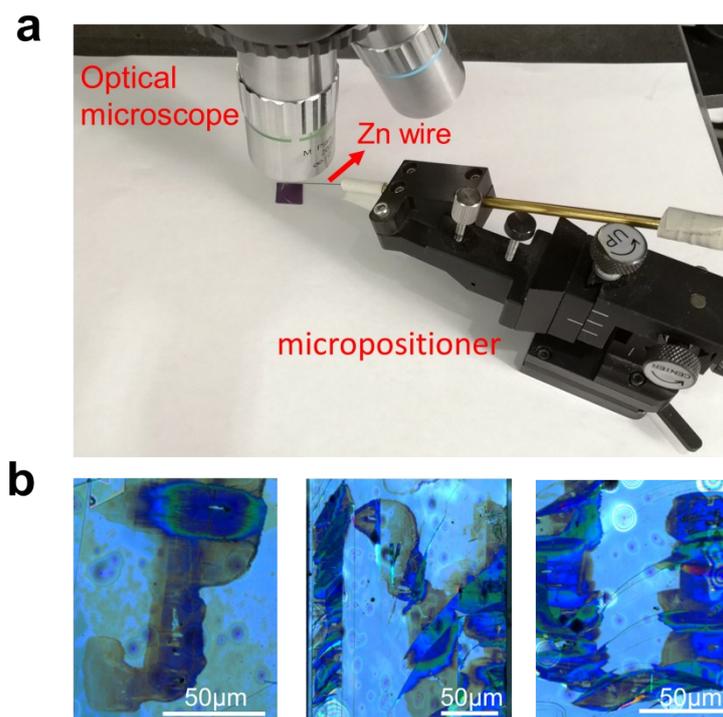


Figure S8 (a) Apparatus to realize local intercalation. (b) Optical image of “JNU” by Co intercalation on a single crystal MoO₃ using a Zn wire.

The Zn wire is fixed on a micropositioner with accuracy of 5 μm (Fig. S8a). The wire is brought by the micropositioner into contact with the surface of the MoO₃ under an optical microscope. The intercalation process can be monitored under the optical microscope.

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