

Supplementary Information: Anisotropic and heterogeneous dynamics in an aging colloidal gel

Avni Jain,^{*a} Florian Schulz^{bc}, Irina Lokteva,^{ac} Lara Frenzel,^{ac} Gerhard Grübel,^{ac} Felix Lehmkuhler^{ac}

Calculation of the temporal fluctuations in $C(q, t_1, t_2)$, $\chi_t(q, \tau, t_{\max})$

The figure below presents a visual aid to equation 6 in the main text.

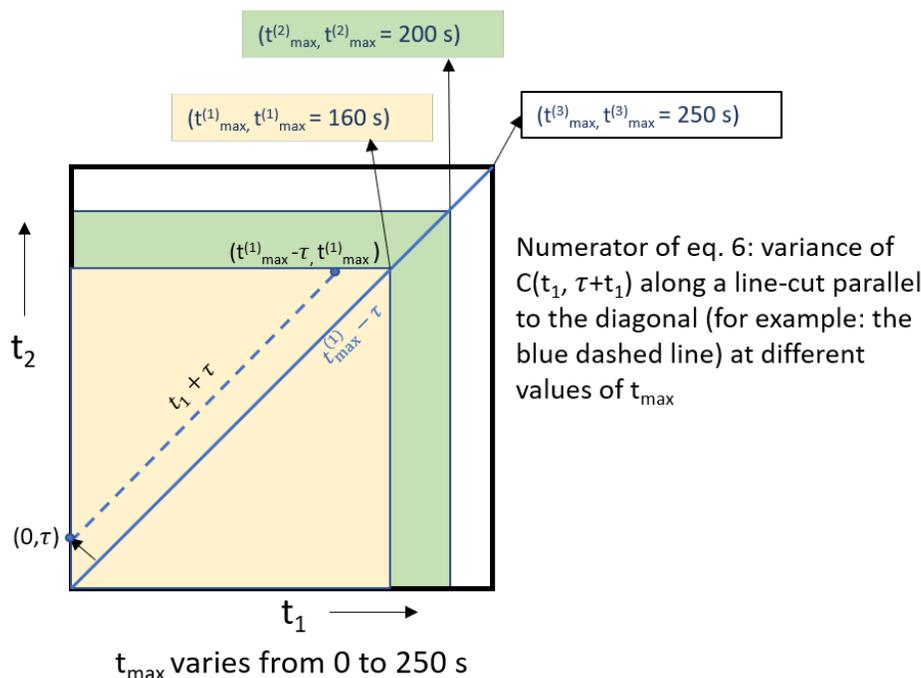


Fig. S1. 1 Calculation of the dynamical heterogeneity, χ_t for different values of t_{\max} from the two-times correlation function, $C(t_1, t_2)$ at a given value of q .

Contribution of aging on χ_t

We calculate the influence of aging dynamics on the normalized variance of fluctuations in $C(q, t_1, t_2)$, χ_t using the procedure described in¹. χ_t can be written as follows,

$$\chi_t(q, \tau, t_{\max}) = \sigma_{g_2}^2 / b^2(q), \quad (\text{S1.1})$$

where $\sigma_{g_2}^2$ is the variance of the fluctuations of the intensity autocorrelation function, g_2 , and $b^2(q)$ is a normalization factor representing the square of the short-time plateau value (or optical contrast) of the corresponding intensity autocorrelation functions. To calculate the effect of aging (i.e., increase of the characteristic relaxation time, τ_c during the experiment duration), we recall that the variance of a quantity y that depends on p variables is expressed as follows,

$$\sigma_y^2 = \sum_{i=1}^p \left[\frac{\partial y}{\partial x_i} \right]^2 \sigma_{x_i}^2 + \sum_{i \neq j} \left[\frac{\partial y}{\partial x_i} \frac{\partial y}{\partial x_j} \right] \sigma_{x_i, x_j}. \quad (\text{S1.2})$$

^a Deutsches Elektronen-Synchrotron (DESY), Notkestraße 85, 22607 Hamburg, Germany; **Email:** avni.jain@desy.de

^b Institute of Physical Chemistry, University of Hamburg, Grindelallee 117, 20146 Hamburg, Germany

^c The Hamburg Centre for Ultrafast Imaging (CUI), Luruper Chaussee 149, 22761 Hamburg, Germany

Here, $\sigma_{x_i}^2 = \langle x_i^2 \rangle - \langle x_i \rangle^2$ is the variance of x_i and $\sigma_{x_i, x_j} = \langle x_i x_j \rangle - \langle x_i \rangle \langle x_j \rangle$ is the covariance between x_i and x_j . The second sum vanishes if the distinct variables x_i 's are uncorrelated. Using equation S1.2 to calculate the contribution of the variable τ_c and assuming that the quantity τ_c is uncorrelated with any other variable, we can write,

$$\sigma_{g_2, \tau_c}^2 = \left[\frac{\partial g_2}{\partial \tau_c} \right]^2 \sigma_{\tau_c}^2. \quad (\text{S1.3})$$

Using the KWW function to describe g_2 (equation 3 in the main text), we find

$$\left[\frac{\partial g_2}{\partial \tau_c} \right]^2 = 4(g_2(q, \tau) - 1)^2 \left(\frac{\tau}{\tau_c} \right)^{2\gamma} \left(\frac{\gamma}{\tau_c} \right)^2. \quad (\text{S1.4})$$

Meanwhile, the aging dynamics in the gels has been described (equation 5 in the main text), where

$$\tau_c(q, t_w) = \frac{\alpha(t_w + \tau_0)}{q}. \quad (\text{S1.5})$$

The variance of τ_c within a certain window of measurement, t_{\max} , is simply, $\sigma_{\tau_c}^2 = \frac{1}{12} \left(\frac{\alpha t_{\max}}{q} \right)^2$. Upon combining the above expressions, equation S1.3 becomes

$$\sigma_{g_2, \tau_c}^2(q, \tau, t_{\max}) = \frac{1}{3} (g_2 - 1)^2 \left(\frac{\tau}{\tau_c} \right)^{2\gamma} \left(\frac{\gamma \alpha t_{\max}}{\tau_c q} \right)^2. \quad (\text{S1.6})$$

Using equation S1.6, for a given value of q and t_{\max} , one can evaluate the aging contribution $\sigma_{g_2, \tau_c}^2(\tau)$. Firstly, one evaluates the autocorrelation function, $g_2(\tau)$. Using the KWW-equation, one then determines the fit parameters of the characteristic relaxation time (τ_c) and stretching exponent (γ). The value of α is used from the analysis of the sample aging behaviour. Note that the value of α may change depending on the sample conditions and t_{\max} .

Characteristic time-scale of the heterogeneity vs. q , χ_t at longer waiting times

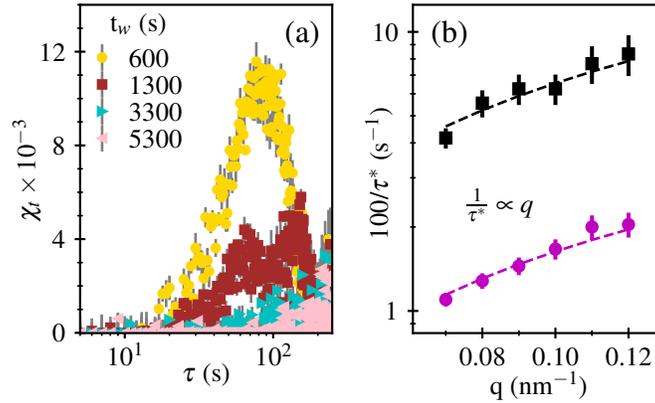


Fig. S1. 2 (a) χ_t curves averaged over all wave vectors are reported for different waiting times, t_w , at $q = 0.07 \text{ nm}^{-1}$, and (b) τ^* vs. q at the earliest waiting time for the gels at temperatures T^* (black squares) and $T^* - 2 \text{ K}$ (magenta circles).

In Fig. 2 (a), χ_t curves averaged over all wave vectors are reported for different waiting times, t_w , at $q = 0.07 \text{ nm}^{-1}$. Here, we observe at $t_w = 1300 \text{ s}$, the χ_t profile displays a double peak suggesting the presence of multiple underlying dynamical entities. At longer waiting times, the dynamical heterogeneities level off over the range of delay times, implying that at $q = 0.07 \text{ nm}^{-1}$, either the timescale of the heterogeneity is beyond the measured time window or simply, in order to resolve the dynamic susceptibility for these considerably slow dynamics, a higher value of t_{\max} (or, duration of the measurement) is required. For example, in Fig.6 of the main text, the first peaked $\chi_t(q_{\parallel})$ curve at $t_w = 600 \text{ s}$ could only be calculated for values of $t_{\max} \geq 190 \text{ s}$. Furthermore, at longer waiting times, as the system has slowed down significantly, the contribution due to aging will be trivial to the measured χ_t . For consistency, we plot the χ_t curve at $t_w = 600 \text{ s}$. The relatively high peak

amplitude is a direct outcome of the aging contribution in the vertical direction.

In Fig. 2 (b), we track the q dependency of τ^* (the delay time corresponding to the maximum value of the averaged χ_t curve) at the earliest waiting time, $t_w = 600$ s, for the gel at both temperatures. We find that $1/\tau^* \propto q$ similar to the q -scaling of the characteristic relaxation rate (in both directions). This trend has also been reported for previous experiments^{2,3}. We note that the slowing down of the relaxation time due to aging also shows a $\propto q$ behaviour, consequently, the timescale of the peak amplitude in the variance or τ^* is unaffected by the aging phenomenon.

Notes and references

- 1 A. Duri, H. Bissig, V. Trappe and L. Cipelletti, *Phys. Rev. E*, 2005, **72**, 051401.
- 2 A. Madsen, R. L. Leheny, H. Guo, M. Sprung and O. Czakkel, *New J. Phys.*, 2010, **12**, 055001.
- 3 F. Ehrburger-Dolle, I. Morfin, F. Bley, F. Livet, G. Heinrich, Y. Chushkin and M. Sutton, *Soft Matter*, 2019, **15**, 3796–3806.