## Penta-Pd $X_2$ (X = S, Se, Te) monolayers: promising

## anisotropic thermoelectric materials

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**Abstract:** Thermoelectric materials can be used to convert waste heat into electrical energy, which is considered to be a cleaner form of energy that reduces carbon dioxide and greenhouse gas emissions. In this work, we study the thermoelectric properties of penta-PdX<sub>2</sub> (X = S, Se, Te) using first-principles calculations together with the Boltzmann transport theory. The dynamic and thermal stability of penta-PdX<sub>2</sub> is confirmed based on phonon dispersion and *ab initio* molecular dynamics (AIMD) simulations. The transport coefficients, such as thermal conductivity ( $k_l$ ), power factor (*PF*) and thermoelectric figure of merit (*ZT*), exhibit strong anisotropy in both *x* and *y* directions. The lower thermal conductivity of penta-PdX<sub>2</sub> along the same direction (*x* or *y* direction) adheres to a trend of PdS<sub>2</sub> > PdSe<sub>2</sub> > PdTe<sub>2</sub>. The anisotropy of *PF* and  $k_l$  of penta-PdX<sub>2</sub> monolayer inevitably leads to anisotropy of *ZT*.

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The largest *ZT* values of penta-PdX2 (X = S, Se) at p-type are 0.85 and 1.18 respectively, while the maximum *ZT* value of penta-PdTe2 has reached 2.42. The predicted *ZT* value in penta-PdX<sub>2</sub> are larger than the commercial TE material Bi<sub>2</sub>Te<sub>3</sub> (about 0.8) and some other transition metal di-chalcogenides, indicating that penta-PdX<sub>2</sub> (X = S, Se, Te) monolayer is a potential anisotropic thermoelectric material.

**Key words:** Pentagonal structure; First-principles investigations; Thermal transport; Thermoelectric material

**Table S1**. Calculated effective mass  $(m^*)$ , average effective mass  $(m_d)$ , elastic modulus  $(C_{2D})$ , deformation potential constant  $(E_1)$ , carrier mobility  $(\mu)$  and relaxation time  $(\tau)$  of penta-PdX<sub>2</sub> (X = S, Se, Te) monolayers at 300 K along the *x* and *y* directions.

	directions	carriers	$m^*(m_e)$	<i>m</i> <sub>d</sub> (m <sub>e</sub> )	$C_{2D} (\mathrm{eV}/\mathrm{\AA}^2)$	$E_l(eV)$	$\mu (\mathrm{cm}^2 \mathrm{V}^{-1} \mathrm{s}^{-1})$	$\tau (10^{-14} \mathrm{s})$
PdS <sub>2</sub>	x	e	0.82	0.48	3.54	-8.80	39.40	1.83
		h	0.46	0.83	3.54	-1.37	1664.15	43.61
	у	e	0.28	0.48	4.89	-9.36	137.49	2.23
		h	1.49	0.83	4.89	-1.46	631.94	53.46
PdSe <sub>2</sub>	x	e	0.33	0.20	2.36	-8.15	177.13	3.33
		h	0.36	0.25	2.36	-2.40	1485.92	30.84
	у	e	0.12	0.20	3.85	-8.85	630.21	4.61
		h	0.18	0.25	3.85	-2.67	3997.68	40.81
PdTe <sub>2</sub>	x	e	0.35	0.32	1.44	-4.01	268.55	5.38
		h	0.16	0.12	1.44	-2.84	2924.89	26.97
	у	e	0.29	0.32	3.19	-4.55	545.75	9.21
		h	0.10	0.13	3.19	-3.34	7454.75	57.88

	result	a (Å)	<i>b</i> (Å)	$d_{ ext{Pd-S}}( ext{\AA})$	$d_{\text{S-S}}(\text{\AA})$	$E_{ m gp} ({ m eV})$
PdS <sub>2</sub>	Present	5.48	5.57	2.33,2.34	2.09	1.14(PBE), 2.12(HSE)
	Exp.	5.41	5.49	-	-	-
	Wang <i>et al</i> .	5.49	5.59	2.34,2.35	2.10	1.60(HSE)
	Deng et al.	5.50	5.59	-	-	1.00(PBE)
PdSe <sub>2</sub>	Present	5.74	5.91	2.45,2.46	2.41	1.34(PBE), 2.14(HSE)
	Exp.	5.74	5.86	-	-	-
	Deng et al.	5.49	5.94	-	-	1.31(PBE)
	Qin <i>et al</i> .	5.75	5.92	-	-	1.38(PBE)
	Oyedele <i>et al</i> .	5.72	5.93	-	-	-
PdTe <sub>2</sub>	Present	6.14	6.44	2.62,2.63	2.80	1.24(PBE),1.84(HSE)
	Deng et al.	5.99	6.37	-	-	-

**Table S2**. The optimized lattice parameters (*a*, *b*), length of Pd-S ( $d_{Pd-S}$ ) and S-S ( $d_{S-S}$ ) bonds, and band gap ( $E_{gp}$ ) of the penta-PdX<sub>2</sub> (X = S, Se, Te) monolayers.

	directions	ZA	TA	LA	optical
DIC	x	13.57%	28.17%	45.03%	13.23%
PdS <sub>2</sub>	У	36.37%	12.47%	44.81%	6.35%
PdSe <sub>2</sub>	x	27.97%	33.22%	17.29%	21.52%
	У	40.30%	27.02%	20.97%	11.71%
PdTe <sub>2</sub>	x	27.74%	18.84%	35.10%	18.32%
	у	49.86%	19.13%	24.17%	6.84%

**Table S3**. Contribution rates of penta-Pd $X_2$  acoustical phonon branches (ZA, TA, LA) and optical phonons support to thermal conductance at 300 K.



**Fig. S1** Snapshot of AIMD simulations at 300 K for (a) penta-PdS<sub>2</sub>, (b) penta-PdSe<sub>2</sub>, and (c) penta-PdTe<sub>2</sub> monolayers, respectively.



Fig. S2 HSE band structures with SOC for (a) penta-PdS<sub>2</sub>, (b) penta-PdSe<sub>2</sub> and (c) penta-PdTe<sub>2</sub> monolayers, respectively.



Fig. S3 Total and partial density of states of (a) penta-PdS<sub>2</sub>, (b) penta-PdSe<sub>2</sub> and (c) penta-PdTe<sub>2</sub> monolayers.



**Fig.S4** Frequency-dependent group velocity of penta- $PdX_2$  (X = S, Se, Te) monolayers

along x (a) and y (b) direction.



**Fig. S5** The electronic thermal conductivity  $(k_e)$  of penta-PdX<sub>2</sub> (X = S, Se, Te) monolayers for both *p*-type (a) and *n*-type (b) at 300 K along the *x* (solid lines) and y (dotted lines) directions.