Supporting Information for

Revealing cooperative Li-ion migration in Li_{1+x}Al_xTi_{2-x}(PO₄)₃ solid state electrolyte with high Al doping

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Computational details

To evaluate the stability of Al-dopants in LATP, we calculate the defect formation energy of Aldopants in $Li_{1+x}Al_xTi_{2-x}(PO_4)_3$ (LATP) using a supercell with six LATP units by the following equation:

$$\Delta E_{\rm f} = E(\text{defected}) + \mathbf{n} \cdot \boldsymbol{\mu}_{\rm Al} - E(\text{bulk}), \tag{1}$$

where E(defected) and E(bulk) represent the total energies of the defected and the pristine LATP supercells, respectively, μ_{Al} is the chemical potential of Al, and n is the number of Al removed from the pristine LATP when defects are generated. For Al chemical potential, we define it by the following equation:

$$\mu_{Al} = [E(Al_2O_3) - 3 \cdot E(O_2)/2]/2, \qquad (2)$$

where $E(Al_2O_3)$ and $E(O_2)$ are the total energies of bulk Al_2O_3 and gas phase O_2 molecular. After computing work, the formation energies of Al vacancies are 4.13, 2.98, and 3.10 eV in LATP-0.16, LATP-0.33 and LAPT-0.50, respectively. The more positive the defect formation energy is, the more unfavorably thermodynamical reaction. Moreover, we calculate the decomposition energy with one combination, ΔE_f , which is defined as

$$\Delta E_{\rm f} (\rm LATP) = E(\rm LTP) - E(\rm LiAlPO_5) - E(\rm LATP) + E(\rm TiP_2O_7). \tag{3}$$

where, the E(LTP), $E(LiAIPO_5)$, E(LATP), and $E(TiP_2O_7)$ is the total energies of $LiTi_2(PO_4)_3$, LiAIPO₅, LATP, and TiP₂O₇. After calculation, the decomposition energies for LATP-0.16, LATP-0.33 and LATP-0.50 are 262.41, 389.93, and 642.34 meV/atom, respectively. Again, these positive values suggest an unfavorably thermodynamical reaction. We note that the extent of Al defect formation and phase decomposition for LATP-0.33 is even a little higher than that for LATP-0.50. Now many experiments have successfully demonstrated the synthesis of LATP-0.50. Thus, the stability of Al-dopant in LATP-0.50 is less problematic in practice because the thermodynamically unfavorable reaction and associated kinetic barrier.

The time-evolution of spatial correlations between atoms is described by the van Hove or spacetime correlation functions. This function can be decomposed into self:

$$G_{s}(r,t) = \frac{1}{4\pi r^{2}N} \langle \sum_{i=1}^{N} \delta[r - |r_{i}(t_{0}) - r_{i}(t + t_{0})|] \rangle_{t_{0}}$$
(4)

and distinct:

$$G_{\rm d}(r,t) = \frac{1}{4\pi r^2 N} \langle \sum_{i\neq j}^N \delta[r - |r_i(t_0) - r_j(t+t_0)|] \rangle_{t_0}$$
(5)

parts, in which the atom whose position is referred to at time *t* is either the same or a different atom from that whose position was specified at time zero. The defined $G(\mathbf{r}, t)$ gives the probability that, at the specified time *t*, an atom will be located a distance *r* away from the location occupied by an atom at an earlier time zero. This method has been widely used to study the diffusion mechanism in other ISEs, such as $\mathrm{Li}_{4\pm x}\mathrm{Si}_{1-x}\mathrm{X}_{x}\mathrm{O4}^{[1]}$ and garnet-type $\mathrm{Li}_{7-x}\mathrm{La}_{3}\mathrm{Zr}_{2-x}\mathrm{M}_{x}\mathrm{O}_{12}^{[2]}$.

To examine the effective migration of Li-ion, we calculate the site displacement function (SDF) of Li-ion as a function of simulation time. The SDF equation is defined d(t) = |r(t) - r(0)|, i.e, the distance between the position of the Li-ion at time t (r(t)) and the initial position of the Li-ion atom at time 0)).



Fig. S1. Relative energies of 11 possible structures considering different AlO₆ arrangement.



Fig. S2. Site displacement function (SDF) plots for one Li-ion trajectory in LATP0.50 taken from 60 ps NVT AIMD simulations at 600 K.



Fig. S3. The MSD of O and P atoms in LTP, LATP0.16 and LATP0.50 taken from 30 ps NVT AIMD simulations at 500 K.

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